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Synthesis, structure and magnetic properties of tetrakis- μ -carboxylato-bis(dodecylnicotinato)dicopper(II) complexes; crystal and molecular structure of the decyl carboxylate derivative

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Dodecylnicotinate bis-adducts of binuclear copper carboxylates, of the general formula $\text{Cu}_2(\text{O}_2\text{CC}_{n-1}\text{H}_{2n-1})_4(\text{C}_5\text{H}_4\text{N}-\text{COOC}_{12}\text{H}_{25})_2$, were synthesized for $n = 10, 12, 14, 16, 18$ and 20 , and their crystal structure, thermal behavior and magnetic properties studied. The molecular structure of the decyl derivative has been determined from single-crystal X-ray diffraction data. The dimer is centrosymmetric with the Cu^{II} ions in a square-pyramidal coordination with four O-alkyl O atoms [average $d(\text{Cu}-\text{O})$ 1.960 (6) Å] in the basal plane and the nicotine N atom at apical positions [$d(\text{Cu}-\text{N})$ 2.183 (3) Å]. The copper ions, 2.615 (1) Å apart, are bridged by four O-alkyl carboxylate groups. Both the $n = 20$ and $n = 18$ homologues exhibit lamellar phases, which can be related to the supramolecular arrangement found in the $n = 10$ derivative. The magnetic behavior of the decyl and octadecyl dimers was studied in the 2–300 K temperature range. They exhibit a strong intramolecular antiferromagnetic interaction ($\text{Cu}-\text{Cu}$ superexchange coupling constant $J = -347 \text{ cm}^{-1}$ for the decyl derivative), which can be attributed to a large overlap of the metal 3d orbitals and the oxygen lone pair orbitals of the linking carboxylate groups.

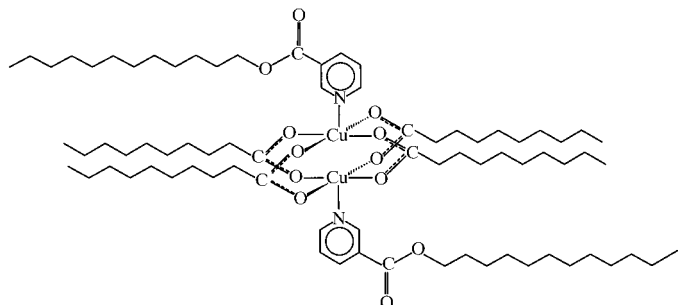
1. Introduction

Investigations on dicopper carboxylates exhibiting the lantern structure have been directed for a variety of interests (Mehrotra & Bohra, 1983; Kato & Muto, 1988). The structural and magnetic properties of the acetate derivative in the past were considered to be very important in the understanding of magnetic exchange in dinuclear compounds, relevant topics in inorganic and physical chemistry, and in the development of appropriate models to describe the magnetic coupling between metal centers (Bleaney & Bowers, 1952; Kahn, 1993).

More recent studies on long-chain derivatives showed their potential application as advanced materials, as they exhibit thermotropic columnar mesophases (Abied *et al.*, 1987; Ibn-Elhaj *et al.*, 1992; Marchon *et al.*, 1992), and also as biologically active substances and antifungal agents, for which bis(nicotinamide derivatives) adducts are examples (Mehrotra & Bohra, 1983; Kozlevcar *et al.*, 1996).

All the copper carboxylate liquid crystals reported thus far are polymeric, the axial position of each Cu^{2+} ion being occupied, in the $\text{Cu}_2(\text{O}_2\text{CR})_4$ series, by an O atom of a neighbor dimer. To break this polymeric structure into distinct dimers we decided to 'block' the axial positions of the long-chain dicopper carboxylates by forming bis(pyridine derivative) adducts (like in the nicotinamide antifungals). As in the equatorial carboxylate, the axial ligand we used also bears a long chain.

We report here the synthesis, spectroscopic and structural characterization, phase transitions and magnetic properties of tetrakis- μ -carboxylato-bis(dodecylnicotinato)-dicopper(II) complexes $\text{Cu}_2(\text{O}_2\text{CC}_{n-1}\text{H}_{2n-1})_4(\text{C}_5\text{H}_4\text{NCOOC}_{12}\text{H}_{25})_2$ for $n = 10, 12, 14, 16, 18$ and 20 , including the crystal and molecular structure of the decyl derivative.



2. Experimental

2.1. Synthesis of dodecyl nicotinate ($\text{C}_5\text{H}_4\text{NCOOC}_{12}\text{H}_{25}$)

Nicotinic acid (2.53 g, 20.5 mmol) was refluxed with a large excess of SOCl_2 for 4 h under an N_2 atmosphere in an anhydrous toluene and DMF (10:1) mixture. Excess thionyl chloride was removed after reaction by vacuum distillation. Dodecanol (14 ml, 62 mmol) dissolved in toluene was added to the nicotiny chloride and the mixture was stirred for 4 h at 353 K under an N_2 atmosphere. The reaction mixture was then poured into water and NaOH 10% was added to alkalinity. The product was extracted with diethyl ether, the organic fraction washed with water and the solvents removed in vacuum. The oily product was purified by silica gel chromatography in cyclohexane: diethyl ether gradient. Yield: 83%. $^1\text{H NMR}$: δ 9.23 (d, 1H, pyridine ring), 8.77 (dd, 1H, pyridine ring), 8.29 (m, 1H, pyridine ring), 7.39 (m, 1H, pyridine ring), 4.35 (t, 2H, CH_2), 1.78 (q, 2H, CH_2), 1.26 (m, 18H, CH_2), 0.88 (t, 3H, CH_3).

2.2. Synthesis of $\text{Cu}_2(\text{O}_2\text{CC}_{n-1}\text{H}_{2n-1})_4(\text{C}_5\text{H}_4\text{NCOOC}_{12}\text{H}_{25})_2$ compounds

Dodecyl nicotinate (2 mmol) was dissolved in 10 ml of boiling *n*-heptane. The corresponding copper carboxylate (0.5 mmol in 10 ml of boiling *n*-heptane), synthesized as described by Ibn-Elhaj *et al.* (1992), was added dropwise. The mixture was refluxed for 30 min and then cooled to 255 K overnight. A green solid was recovered by filtration at low temperature. The solid was dried in vacuum over silica. The compound with $n = 10$ crystallized as green platelet crystals, from which single-crystal samples were selected for crystallographic studies.

Anal.: exp. (calc.) for $n = 10$: C 65.57 (65.43), N 2.08 (2.01), H 9.39 (9.68), yield 49%; for $n = 12$: C 66.65 (66.94), N 1.85 (1.85), H 9.91 (10.03), yield 54%; for $n = 14$: C 67.99 (68.23), N 1.79 (1.73), H 9.19 (10.33), yield 84%; for $n = 16$: C 69.59 (69.36), N 1.61 (1.61), H 10.68 (10.59), yield 91%; for $n = 18$: C

69.59 (70.35), N 1.49 (1.51), H 10.76 (10.82), yield 70%; for $n = 20$: C 71.12 (71.23), N 1.49 (1.43), H 10.94 (11.03), yield 75%.

2.3. Physicochemical methods

IR spectroscopy was carried out on a Nicolet FTIR 510P with KBr pellets and UV-vis measurements on a HP8453 spectrometer. $^1\text{H NMR}$ spectra were run on a Bruker AC200 in CDCl_3 . Differential scanning calorimetry (DSC) was performed on a Shimadzu DSC-50 calorimeter. For optical microscopy observations, a Leitz DMRX microscope equipped with a Leitz 1350 heating stage device was used. Elemental analyses were performed on a Carlo Erba CHNS-O EA1108 instrument.

Magnetic measurements between 2 K and room temperature were performed on a Quantum Design MPMS-XL susceptometer in a 5000 G magnetic field. The powdered samples, typically 5–15 mg, were placed in a cellophane bag. Raw magnetic data were corrected for the sample holder contribution and for the sample diamagnetism, which was evaluated using Pascal's constants (Selwood, 1956).

Variable-temperature powder X-ray diffraction (XRD) patterns were recorded using an INEL CPS-120 curved position-sensitive detector and $\text{Cu K}\alpha$ radiation from an INEL X-ray generator; room-temperature data were collected on Siemens D5000 equipment.

2.4. Single-crystal diffraction data and structure solution and refinement of $\text{Cu}_2(\text{O}_2\text{CC}_9\text{H}_{19})_4(\text{C}_5\text{H}_4\text{NCOOC}_{12}\text{H}_{25})_2$

Crystal data, data collection procedure, structure determination methods and refinement results are summarized in Table 1.¹ As expected for a complex with extended organic ligands, the samples turned out to diffract poorly. Only about 59% of the X-ray intensity data up to 0.841 Å resolution were above two standard deviations of measurement errors.

Several H atoms were detected at approximate locations in a difference Fourier map. They, however, were positioned stereochemically and refined with the riding model. During the refinement, the CH_3 H atoms were allowed to rotate as a rigid group around the C– CH_3 bond so as to maximize the sum of the electron density at the three calculated hydrogen positions.

3. Results and discussion

3.1. Synthesis and characterization

The formation of the expected complexes was first studied in solution. Experiments carried out with $\text{Cu}_2(\text{O}_2\text{CC}_9\text{H}_{19})_4$ dissolved in toluene showed a progressive conversion of the 678 nm absorption band of the starting complex into a new band centered at *ca* 702 nm, as the dodecyl nicotinate

¹ Supplementary data for this paper, including listings of atomic anisotropic displacement parameters, full bond distances and angles, H-atom positions, calculated and observed structure factor amplitudes for the $n = 10$ compound are available from the IUCr electronic archives (Reference: CA0004). Services for accessing these data are described at the back of the journal. Powder XRD patterns for $n = 12, 14, 16, 18$ and 20 derivatives are available from the authors.

Table 1

Experimental details.

Crystal data	
Chemical formula	[Cu(C ₁₀ H ₁₉ O ₂) ₄ (C ₁₈ H ₂₉ NO ₂) ₂]
Chemical formula weight	1394.93
Cell setting	Triclinic
Space group	<i>P</i> $\bar{1}$
<i>a</i> (Å)	10.231 (3)
<i>b</i> (Å)	12.878 (4)
<i>c</i> (Å)	16.804 (2)
α (°)	92.39 (2)
β (°)	104.31 (2)
γ (°)	106.17 (3)
<i>V</i> (Å ³)	2045.8 (9)
<i>Z</i>	1
<i>D_x</i> (Mg m ⁻³)	1.132
<i>D_m</i> (Mg m ⁻³)	1.12
Density measured by	Immersion
Radiation type	Mo <i>K</i> α
Wavelength (Å)	0.71073
<i>F</i> (000)	758
No. of reflections for cell parameters	23
θ range (°)	9.8–15.71
μ (mm ⁻¹)	0.574
Temperature (K)	293 (2)
Crystal form	Plate
Crystal size (mm)	0.44 × 0.40 × 0.16
Crystal colour	Blue–green
Data collection	
Diffractometer	Enraf–Nonius CAD-4
Data collection method	ω – 2θ scans
Absorption correction	Numerical (Busing & Levy, 1957)
<i>T_{min}</i>	0.793
<i>T_{max}</i>	0.916
No. of measured reflections	7627
No. of independent reflections	6369
No. of observed reflections	4529
Criterion for observed reflections	<i>I</i> > 2 σ (<i>I</i>)
<i>R_{int}</i>	0.036
θ_{max} (°)	24.97
Range of <i>h, k, l</i>	–12 → <i>h</i> → 11 –15 → <i>k</i> → 15 0 → <i>l</i> → 19
No. of standard reflections	1
Frequency of standard reflections	Every 30 min
Intensity decay (%)	1.5
Refinement	
Refinement on	<i>F</i> ²
<i>R</i> [<i>F</i> ² > 2 σ (<i>F</i> ²)]	0.0504
<i>wR</i> (<i>F</i> ²)	0.1295
<i>S</i>	1.086
No. of reflections used in refinement	6369
No. of parameters used	423
H-atom treatment	H-atom parameters constrained
Weighting scheme	$w = 1/[\sigma^2(F_o^2) + (0.0878P)^2 + 0.2591P]$, where $P = (F_o^2 + 2F_c^2)/3$
(Δ/σ) _{max}	–0.006
$\Delta\rho_{max}$ (e Å ⁻³)	0.393
$\Delta\rho_{min}$ (e Å ⁻³)	–0.288
Extinction method	None
Source of atomic scattering factors	<i>International Tables for Crystallography</i> (1992, Vol. C, Tables 4.2.6.8 and 6.1.1.4)
Computer programs	
Data collection	Enraf–Nonius CAD-4
Cell refinement	Enraf–Nonius CAD-4
Data reduction	<i>SDP</i> (B. A. Frenz & Associates, 1983)

Table 1 (continued)

Structure solution	<i>SHELXS86</i> (Sheldrick, 1990)
Structure refinement	<i>SHELXL93</i> (Sheldrick, 1993)
Preparation of material for publication	<i>SHELXL93</i> (Sheldrick, 1993)

concentration increased. This shift is consistent with a change in the coordination number around Cu²⁺ from four to five. For example, the pyridine adduct band shifts down in energy to 713 nm owing to its strong interaction with copper. Job's method (Reilley & Sawyer, 1961) clearly indicated a 1:2 complex-to-nicotinate ligand ratio, as expected for a bis-adduct, but also showed that the stability constant for this Cu₂(O₂CR)₄(nicotinate)₂ compound is not very high. To overcome this we performed the synthesis of the compounds in excess dodecyl nicotinate ligand; their high solubility in non-polar solvents indicated that a low-temperature crystallization process was necessary.

The characterization of the solid product was completed by the crystal structure determination of the decyl derivative by XRD methods (see below) and IR spectroscopy. The main IR bands are, in addition to the ν_{CH_3} , ν_{CH_2} stretching vibrations characteristic of long-chain organic compounds, the ν_{CO_2} vibration of the ester at 1725–1740 cm⁻¹, the ν_{py} symmetric stretching mode of the pyridine ring (1591 cm⁻¹ in the free ligand) shifted by coordination to 1616–1618 cm⁻¹, as well as the ν_{CO_2} symmetric and asymmetric stretching modes of the bridging carboxylates at 1424–1437 cm⁻¹ (depending on the carboxylate chain length) and 1599–1588 cm⁻¹, respectively. These values agree with those found for bis-nicotinamide copper carboxylates (Kozlevcar *et al.*, 1996).

3.2. Crystallographic structural results and discussion for the *n* = 10 derivative

Atomic fractional coordinates and equivalent isotropic displacement parameters are given in Table 2. Selected bond distances and angles are listed in Table 3. Fig. 1 is an *ORTEP* (Johnson, 1965) drawing of the dimeric complex.

The Cu₂(O₂C(CH₂)₈CH₃)₄(C₅H₄NCO₂(CH₂)₁₁CH₃)₂ complex is located on a crystallographic inversion center. The pair of Cu^{II} ions in a dimer are bridged through the carboxylate groups of four CH₃(CH₂)₈CO₂⁻ ligands. Each Cu atom is in an elongated octahedral environment, coordinated equatorially to four carboxylate O atoms of two independent CH₃(CH₂)₈CO₂⁻ moieties [Cu–O bond distances vary from 1.950 (3) to 1.966 (2) Å] and axially by the pyridine nitrogen of a C₅H₄NCO₂(CH₂)₁₁CH₃ molecule [*d*(Cu–N) = 2.183 (3) Å] and by the other inversion-related Cu atom located at 2.615 (1) Å. The local molecular symmetry is close to *D*_{4h}. The nicotinate ring plane is approximately coplanar to one of the C–O–Cu–O–C bonds, but slightly bent by 11° with respect to the Cu–Cu direction.

The Cu—Cu distance is the second shortest Cu—Cu distance found for $\text{Cu}_2(\text{O}_2\text{C}_n\text{H}_{2n-1})_4\text{L}_2$ compounds. This can be related to the moderate basicity of the axial pyridinic ligand (Melník *et al.*, 1985; Melník, 1982); in fact, the only compound which exhibits a shorter Cu—Cu distance is the pyrazine derivative, the least basic axial ligand ($\text{p}K_a = 0.65$) for which the bis-adduct has been structurally characterized. The Cu—O distances are also slightly shorter than the corresponding values reported previously in the range 1.96–1.98 Å (Melník *et al.*, 1985). The observed Cu—O_{eq}/Cu—N_{ax} distances relationship is as expected on the basis of the local electro-neutrality principle (Kato & Muto, 1988; Pauling, 1960) for such a weak basic ligand.

One of the two independent alkyl chains of the equatorial carboxylates is completely elongated, in a fully extended zigzag *trans* arrangement, whereas the other aligns parallel to the first one after a *gauche* conformation at the C(12)—C(13) bond. The alkyl chain of the nicotinate ligand is also in an extended conformation, after turning at the C(37) atom, rotated from the nicotinate plane to align with the carboxylate chains (see Fig. 1). Both the equatorial ligands and the nicotinate alkyl chain are straight and tilt slightly with respect to the plane perpendicular to the N—Cu—Cu—N axis. Three crystallographically non-equivalent alkyl chains converge in a plane: that belonging to the elongated carboxylate is positioned between the alkyl chain of the other non-equivalent carboxylate from the same dimer and a third one belonging to a nicotinate of a neighboring molecule. These alkyl chains are completely interdigitated (see Fig. 2), leading to the formation of two different layers along the crystal: one is defined by the polar copper carboxylate cores and the second, non-polar one contains the alkyl chains. This type of polar–nonpolar layer arrangement is typically found in mesomorphic systems.

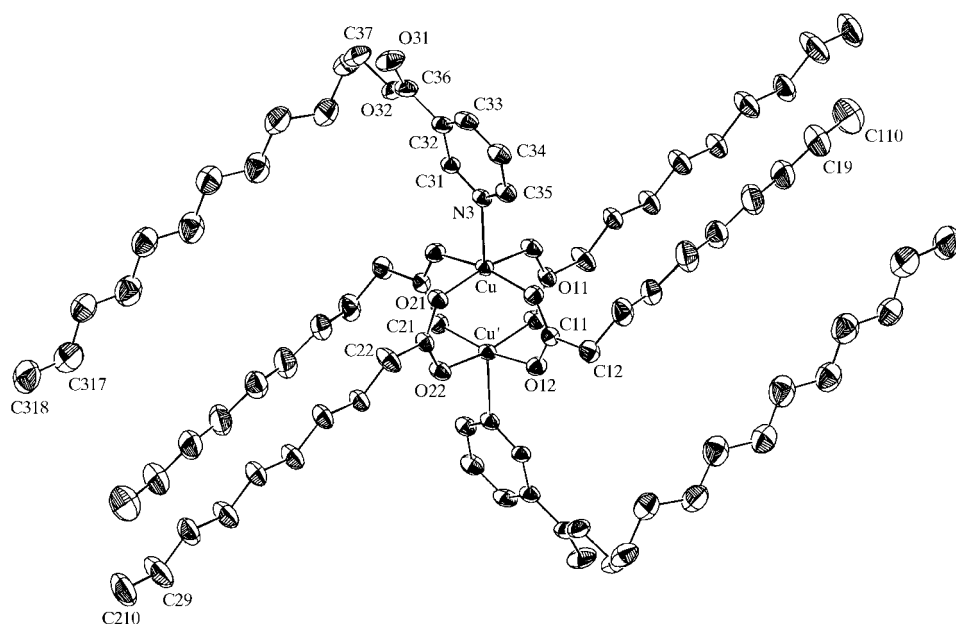


Figure 1
ORTEP (Johnson, 1965) drawing of $\text{Cu}_2(\text{O}_2\text{CC}_9\text{H}_{19})_4(\text{C}_5\text{H}_4\text{NCOOC}_{12}\text{H}_{25})_2$ showing the numbering scheme of the non-H atoms.

3.3. Phase transitions and thermal behavior of the $\text{Cu}_2(\text{O}_2\text{CC}_{n-1}\text{H}_{2n-1})_4(\text{C}_5\text{H}_4\text{NCO}_2\text{C}_{12}\text{H}_{25})_2$ series

The compound in which $n = 20$ exhibits a lamellar phase from room temperature up to 377 K, where it melts to the isotropic liquid. The interlamellar distance, d_{100} , is 49.4 Å at 298 K, but it sharply changes to 49.1 Å at 325 K and then to 52.8 Å at 339 K. This contraction–expansion sequence correlates with phase transitions detected in DSC measurements. Table 4 contains the thermodynamic parameters of these transitions. The experimental interlamellar distance is in good agreement with the molecular length in an elongated conformation, estimated as 52 Å by extrapolation from the solved structure for the $n = 10$ derivative. The $n = 18$ homologue showed a transition to a lamellar phase at 333 K ($d_{100} = 46.7$ Å, consistent with an estimated length of 47 Å). During the phase transitions of both compounds some differences in the high diffraction angle region have been observed, revealing changes in the crystalline order at the aliphatic chain packing level. The powder XRD patterns of the $n = 12, 14$ and 16 derivatives are much more complicated, precluding any straightforward assignment. As mentioned above, the crystal structure of the decyl derivative can also be interpreted as a layered arrangement of polar cores on the *ab* plane separated by aliphatic chains, but now the interlamellar separation is about half that expected for a segregated bilayer. This is due to the intercalation of the alkyl chains belonging to neighboring dimers related to each other by a unit cell translation along the *c* axis.

Interestingly, the differences in the supramolecular features between long-chain and short-chain derivatives can be correlated with molecular characteristics. Indeed, the IR spectra showed that the position of the ν_{CO_2} symmetric stretching band depends on the chain length: it appears at 1599 cm^{-1} for the $n = 10, 12, 14$ and 16 compounds, and at 1588 cm^{-1} for the $n = 18$ and 20 derivatives. The difference between the ν_{CO_2} asymmetric and symmetric stretching mode wavenumbers decreases from 176 cm^{-1} for $n = 10$ to 150 cm^{-1} for $n = 20$, thus reflecting a more symmetric coordination (Deacon & Phillips, 1980) of the equatorial ligands as the chain length increases.

No liquid crystal phases have been found for these compounds, despite the presence of the long alkyl chains and the rich polymorphism they exhibited before melting to the isotropic liquid. All the studied compounds decompose between 493 and 513 K. The thermal stability is higher than those of nicotinamide adducts of

copper carboxylates (Kozlevcar *et al.*, 1996).

3.4. Magnetic behavior

Fig. 3 shows the molar magnetic susceptibility (χ) of the $\text{Cu}_2(\text{O}_2\text{CC}_9\text{H}_{19})_4(\text{C}_5\text{H}_4\text{NCOOC}_{12}\text{H}_{25})_2$ dimer measured as a function of temperature. As for the other lantern-structured copper carboxylates, these data correspond to a system with two antiferromagnetically coupled $S = \frac{1}{2}$ centers (the increase of χ at low temperature is due to a very small amount of a paramagnetic impurity, very likely being a monomeric Cu^{2+} species). The octadecyl derivative, $\text{Cu}_2(\text{O}_2\text{CC}_{17}\text{H}_{35})_4(\text{C}_5\text{H}_4\text{NCOOC}_{12}\text{H}_{25})_2$, exhibits essentially the same χ versus T variation.

The magnetic behavior of the copper unpaired electrons ($S_1 = S_2 = 1/2$) in the dimeric complex can be described by the following Hamiltonian (Bleaney & Bowers, 1952)

$$H = (V_1 + V_2) - JS_1 \cdot S_2 + \gamma(\mathbf{L}_1 \cdot \mathbf{S}_1 + \mathbf{L}_2 \cdot \mathbf{S}_2) + \beta\mathbf{H}(\mathbf{L}_1 + 2\mathbf{S}_1 + \mathbf{L}_2 + 2\mathbf{S}_2), \quad (1)$$

where the first term describes ligand field effects, the second is the exchange coupling between the electrons, and the third and fourth contributions are the spin-orbit and Zeeman interactions. The crystal field effects lead to a mainly metal $d(x^2 - y^2)$ ground state orbital with its lobes along the Cu–O bonds, energetically well separated from the electronic excited states. The exchange interaction produces the largest (first-order) perturbation to the spin fourfold-degenerated ground state by splitting it into diamagnetic singlet (total spin $S = 0$) and paramagnetic triplet ($S = 1$) levels, energy separated in $|J|$. Perturbation theory up to second-order including exchange spin-orbit and Zeeman interactions shows that the triplet state can be described by the effective spin ($S = 1$) Hamiltonian

$$H_{\text{eff}} = \mathbf{S} \cdot \mathbf{D} \cdot \mathbf{S} + \beta\mathbf{S} \cdot \mathbf{g} \cdot \mathbf{H}, \quad (2)$$

where \mathbf{D} and \mathbf{g} are the fine and gyromagnetic tensors, respectively. The fine interaction splits the triplet state in the absence of an external magnetic field into a singlet and a doublet. It arises from the combined effect of ligand field, exchange and spin-orbit interactions and vanishes for uncoupled ($J = 0$) electrons. The symmetric \mathbf{D} and \mathbf{g} tensors share the same principal (x, y, z) axes, where the spin Hamiltonian adopts the form

$$H_{\text{eff}} = DS_z^2 + E(S_x^2 - S_y^2) + \beta(g_x H_x S_x + g_y H_y S_y + g_z H_z S_z). \quad (3)$$

Single-crystal EPR measurements in the closely related system of $\text{Cu}_2(\text{CH}_3\text{COO})_4(\text{H}_2\text{O})_2$ shows that the dimer exhibits approximate axial symmetry and therefore $E \simeq 0$, $g_x \simeq g_y \simeq g_{\perp}$ and $g_z = g_{\parallel}$ (Bleaney & Bowers, 1952). Assuming the same holds for our system, (3) further simplifies to

$$H_{\text{eff}} = DS_z^2 + \beta[g_{\perp}(H_x S_x + H_y S_y) + g_{\parallel} H_z S_z]. \quad (4)$$

The fine and Zeeman interactions (a fraction of 1 cm^{-1}) are much smaller than the exchange interaction (in the present case, a few hundredths of a cm^{-1}). This interaction, in turn, is much less than the energy separation between electronic

Table 2

Fractional atomic coordinates and equivalent isotropic displacement parameters (\AA^2).

	$U_{\text{eq}} = (1/3)\Sigma_i \Sigma_j U^{ij} a_i' a_j'$			U_{eq}
	x	y	z	
Cu	0.10395 (4)	0.45719 (3)	0.01842 (2)	0.0564 (2)
O11	0.2103 (2)	0.6037 (2)	0.0749 (2)	0.0755 (7)
O12	0.0340 (3)	0.6753 (2)	0.0450 (2)	0.0738 (7)
C11	0.1594 (4)	0.6801 (3)	0.0795 (2)	0.0635 (9)
C12	0.2546 (5)	0.7861 (3)	0.1286 (3)	0.0900 (13)
C13	0.3620 (6)	0.7753 (4)	0.2037 (3)	0.119 (2)
C14	0.3002 (5)	0.7241 (4)	0.2708 (3)	0.1032 (15)
C15	0.4061 (8)	0.7233 (5)	0.3483 (4)	0.149 (3)
C16	0.3615 (7)	0.6762 (5)	0.4168 (4)	0.125 (2)
C17	0.4698 (9)	0.6783 (6)	0.4930 (4)	0.156 (3)
C18	0.4390 (8)	0.6327 (6)	0.5634 (4)	0.150 (2)
C19	0.5486 (11)	0.6364 (7)	0.6368 (5)	0.193 (4)
C110	0.5294 (12)	0.5972 (8)	0.7081 (6)	0.243 (5)
O21	0.1358 (2)	0.5074 (2)	−0.08604 (14)	0.0690 (6)
O22	−0.0400 (3)	0.5803 (2)	−0.1173 (2)	0.0776 (7)
C21	0.0622 (4)	0.5559 (3)	−0.1319 (2)	0.0652 (9)
C22	0.1020 (5)	0.5869 (5)	−0.2092 (3)	0.107 (2)
C23	−0.0062 (4)	0.6012 (4)	−0.2789 (2)	0.0834 (12)
C24	0.0357 (5)	0.6294 (4)	−0.3555 (2)	0.0963 (14)
C25	−0.0698 (5)	0.6432 (5)	−0.4265 (3)	0.111 (2)
C26	−0.0331 (6)	0.6702 (4)	−0.5030 (3)	0.1043 (15)
C27	−0.1412 (7)	0.6847 (6)	−0.5740 (3)	0.136 (2)
C28	−0.1105 (7)	0.7088 (5)	−0.6513 (3)	0.132 (2)
C29	−0.2204 (9)	0.7217 (8)	−0.7215 (4)	0.195 (4)
C210	−0.1972 (10)	0.7449 (7)	−0.7960 (4)	0.206 (4)
N3	0.2843 (3)	0.3929 (3)	0.0471 (2)	0.0620 (7)
C31	0.2706 (4)	0.2926 (3)	0.0172 (2)	0.0650 (9)
C32	0.3890 (4)	0.2567 (3)	0.0207 (2)	0.0731 (10)
C33	0.5196 (4)	0.3283 (4)	0.0548 (3)	0.0857 (13)
C34	0.5326 (4)	0.4298 (4)	0.0875 (3)	0.0880 (13)
C35	0.4130 (4)	0.4599 (4)	0.0821 (2)	0.0749 (11)
C36	0.3746 (5)	0.1470 (4)	−0.0151 (3)	0.0896 (13)
O31	0.4702 (4)	0.1206 (3)	−0.0311 (3)	0.1292 (13)
O32	0.2456 (4)	0.0796 (2)	−0.0285 (2)	0.0933 (9)
C37	0.2233 (7)	−0.0308 (4)	−0.0645 (4)	0.123 (2)
C38	0.0733 (8)	−0.0924 (4)	−0.0752 (4)	0.132 (2)
C39	−0.0347 (7)	−0.0551 (5)	−0.1366 (4)	0.128 (2)
C310	−0.0369 (8)	−0.0738 (5)	−0.2247 (5)	0.144 (2)
C311	−0.1377 (7)	−0.0354 (6)	−0.2844 (5)	0.149 (2)
C312	−0.1397 (8)	−0.0449 (6)	−0.3740 (5)	0.145 (2)
C313	−0.2330 (8)	0.0001 (6)	−0.4301 (5)	0.157 (3)
C314	−0.2346 (8)	−0.0035 (5)	−0.5189 (5)	0.144 (2)
C315	−0.3294 (8)	0.0446 (6)	−0.5730 (5)	0.161 (3)
C316	−0.3324 (8)	0.0469 (5)	−0.6608 (5)	0.139 (2)
C317	−0.4337 (11)	0.0948 (8)	−0.7127 (6)	0.201 (4)
C318	−0.4441 (11)	0.1012 (7)	−0.7969 (6)	0.198 (4)

ground and excited states. This warrants the description of the powder molar magnetic susceptibility by the Van Vleck expression, which to second-order approximation leads to the Bleaney–Bowers equation

$$\chi(T) = \left[\frac{2Ng^2\beta^2}{kT(3 + e^{-J/kT})} + \text{TIP} \right] (1 - \rho) + \frac{Ng_p^2\beta^2}{2kT} \rho. \quad (5)$$

The TIP contribution to χ corresponds to temperature-independent paramagnetic effects and the last term in (5) accounts for the presence of small amounts of paramagnetic impurities in a proportion ρ . g^2 is an isotropic average of the squared \mathbf{g} -tensor, the g_p value can be taken to be equal to 2 and the other parameters have their usual meaning.

Table 3
Selected geometric parameters (Å, °).

Cu—O12 ⁱ	1.950 (3)	O11—C11	1.242 (4)
Cu—O11	1.960 (3)	O12—C11	1.254 (4)
Cu—O22 ⁱ	1.962 (2)	C11—C12	1.504 (5)
Cu—O21	1.966 (2)	O21—C21	1.252 (4)
Cu—N3	2.183 (3)	O22—C21	1.247 (4)
Cu—Cu ⁱ	2.6146 (10)	C21—C22	1.493 (5)
N3—C31	1.323 (4)	C33—C34	1.356 (6)
N3—C35	1.329 (4)	C34—C35	1.366 (5)
C31—C32	1.401 (5)	C36—O31	1.206 (5)
C32—C33	1.362 (6)	C36—O32	1.321 (5)
C32—C36	1.466 (6)		
O12 ⁱ —Cu—O22 ⁱ	90.31 (12)	O11—Cu—Cu ⁱ	82.90 (8)
O11—Cu—O22 ⁱ	89.42 (12)	O22 ⁱ —Cu—Cu ⁱ	85.18 (7)
O12 ⁱ —Cu—O21	88.59 (11)	O21—Cu—Cu ⁱ	83.16 (7)
O11—Cu—O21	89.27 (11)	N3—Cu—Cu ⁱ	177.08 (8)
O22 ⁱ —Cu—O21	168.33 (9)	C11—O11—Cu	125.0 (2)
O12 ⁱ —Cu—N3	96.76 (11)	C11—O12—Cu ⁱ	122.5 (2)
O11—Cu—N3	95.15 (11)	O11—C11—O12	124.4 (3)
O22 ⁱ —Cu—N3	97.00 (10)	O11—C11—C12	118.1 (3)
O21—Cu—N3	94.66 (10)	O12—C11—C12	117.5 (3)
O12 ⁱ —Cu—Cu ⁱ	85.15 (8)	C21—O21—Cu	124.4 (2)
C21—O22—Cu ⁱ	122.3 (2)	C31—C32—C36	121.4 (4)
O22—C21—O21	124.9 (3)	C34—C33—C32	120.0 (3)
O22—C21—C22	118.4 (3)	C33—C34—C35	118.7 (4)
O21—C21—C22	116.7 (3)	N3—C35—C34	122.6 (4)
C31—N3—C35	119.0 (3)	O31—C36—O32	123.2 (5)
C31—N3—Cu	120.2 (2)	O31—C36—C32	123.7 (5)
C35—N3—Cu	119.7 (3)	O32—C36—C32	113.1 (4)
N3—C31—C32	121.2 (4)	C36—O32—C37	116.1 (4)
C33—C32—C31	118.3 (4)	O32—C37—C38	107.7 (4)
C33—C32—C36	120.2 (4)		

Symmetry code: (i) $-x, 1-y, -z$.

Least-squares fit of (5) to the susceptibility data leads to the following parameters: $-J = 347 (1) \text{ cm}^{-1}$, $g = 2.037 (4)$ and $\rho = 0.65 (1)\%$; the value of the TIP is not meaningful because of uncertainties in the diamagnetic correction of such a large molecule. The corresponding $\chi(T)$ curve is compared in Fig. 3 with the experimental susceptibility.

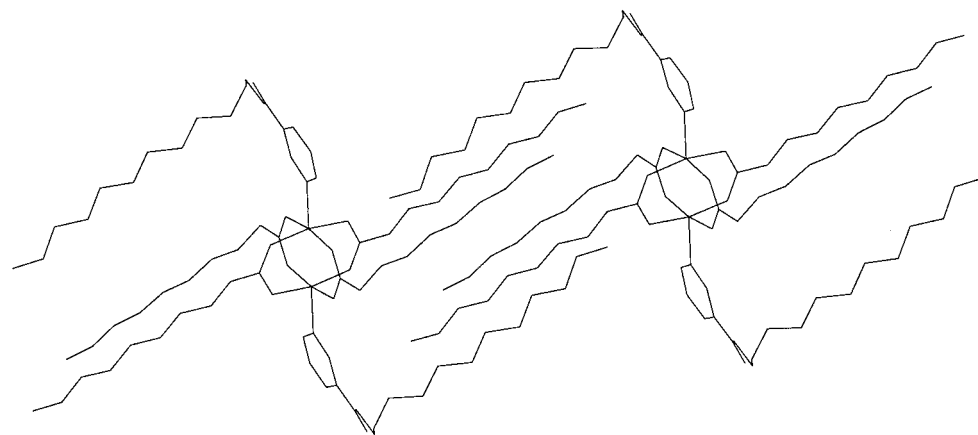


Figure 2
ORTEP (Johnson, 1965) drawing showing two dimeric units of $\text{Cu}_2(\text{O}_2\text{CC}_9\text{H}_{19})_4(\text{C}_5\text{H}_4\text{NCOOC}_{12}\text{H}_{25})_2$ symmetry related by a unit-cell translation along c .

Table 4
Transition temperature (K) and enthalpy (in kJ mol^{-1} within parentheses) for $\text{Cu}_2(\text{O}_2\text{CC}_n - 1\text{H}_{2n - 1})_4(\text{C}_5\text{H}_4\text{NCOOC}_{12}\text{H}_{25})_2$ derivatives.

n	Phase transition	Transition temperature (K)
$n = 10$	$C \rightarrow I$	315 (97)
$n = 12$	$C_1 \rightarrow C_2$	318 (<2)
	$C_2 \rightarrow I$	328 (90)
$n = 14$	$C_1 \rightarrow C_2$	316 (3)
	$C_2 \rightarrow I$	334 (162)
$n = 16$	$C_1 \rightarrow C_2$	310 (14)
	$C_2 \rightarrow C_3$	337 (153)
	$C_3 \rightarrow I$	364 (29)
$n = 18$	$C_1 \rightarrow C_2$	333 (73)
	$C_2 \rightarrow I$	370 (70)
$n = 20$	$C_1 \rightarrow C_2$	325 (13)
	$C_2 \rightarrow C_3$	339 (121)
	$C_3 \rightarrow I$	377 (32)

The compound shows a strong intramolecular anti-ferromagnetic interaction, which can be attributed to Cu—Cu superexchange through the carboxylate groups. In order to correlate these magnetic results with structural characteristics, we verified that the diffraction pattern of the bulk powdered sample used for magnetic experiments effectively agrees with that calculated from the crystallographic data. In general, the values of the magnetic parameters agree with those already published for $\text{Cu}_2(\text{O}_2\text{C}_n\text{H}_{2n-1})_4L_2$ compounds, where L coordinates *via* N atoms. The more relevant parameter, $|J|$, is near the upper limit of the thus far reported values ($308\text{--}353 \text{ cm}^{-1}$; Kato & Muto, 1988; Kawata *et al.*, 1992; Melník *et al.*, 1985), which seems in agreement with the short Cu—O bond distances. The correlation of $|J|$ with the basicity of the ligand has been the subject of some controversy in the literature. Indeed, some authors invoked an increase of the exchange parameter as the axial ligand becomes a stronger electron donor (Melník *et al.*, 1985; Melník, 1982); contrarily, others claimed (Kato & Muto, 1988) that the AF superexchange increases as the axial ligand basicity decreases, a fact that has been interpreted as a consequence of the strong Cu—

O bonds (shorter distances) present in order to maintain the local electroneutrality of copper(II) ions, which enhance the overlap between magnetic orbitals through the bridge's HOMO, resulting in a stronger AF interaction. Our results seem to support the conclusion drawn by Kato and Muto. However, the magnetic coupling described by J depends simultaneously on several factors and it cannot be assigned only to the ligand basicity. For example, it has been found that there is no correlation between J and the Cu—Cu distance (as expected for a bridged-mediated coupling

mechanism; Melník, 1982), but a marked one with the bend angle defined by the O—Cu—Cu—O and O—C—O planes (Kawata *et al.*, 1992): the smaller the bend angle, the stronger the superexchange coupling through the carboxylates, due to a good overlap between magnetic orbitals in the Cu—O—C—O—Cu fragment. In the present compound the very small bend angle (0.8°) agrees with the high $|J|$ value.

4. Conclusions

The use of long-chain nicotinate derivatives as axial ligands was successful in order to break out the polymeric structure of dicopper carboxylates.

Even if they did not present liquid crystal phases, there is evidence in the crystal structure of the decylcarboxylate adduct of a microsegregation of the alkyl chains from the remaining, and more polar, part of the molecule. In fact, the compounds with the longest alkyl chains exhibit a lamellar crystalline structure whose interlamellar distance was related to the largest dimension of one molecule. Those results indicate that this system could be slightly modified to induce the appearance of liquid crystal smectic phases.

The structural analysis of the decyl derivative allowed the interpretation of the magnetic behavior of this compound, which exhibits one of the strongest antiferromagnetic interactions among this type of dinuclear copper compounds. The short Cu—O distance, in addition to the coplanarity of the Cu—O—C—O—Cu group, warrant a large overlap between the $d_{x^2-y^2}$ magnetic orbital of the Cu^{2+} cation and the lone pair of the O atom that leads to a strong superexchange interaction between the metallic atoms through the carboxylate's bridge.

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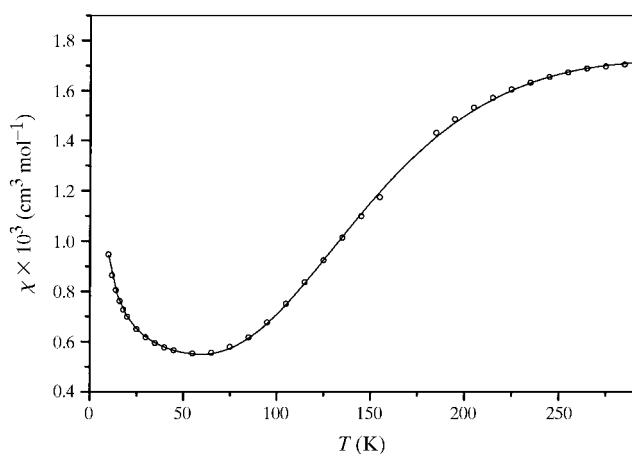


Figure 3 Temperature dependence of the measured molar magnetic susceptibility χ (○) for the complex $\text{Cu}_2(\text{O}_2\text{CC}_9\text{H}_{19})_4(\text{C}_5\text{H}_4\text{NCOOC}_{12}\text{H}_{25})_2$. The solid line corresponds to the value calculated with the model described in the text.

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