TRANSACTIONS OF THE AMERICAN MATHEMATICAL SOCIETY Volume 359, Number 5, May 2007, Pages 2137-2154 S 0002-9947(06)04100-6 Article electronically published on December 15, 2006

ON RIESZ TRANSFORMS AND MAXIMAL FUNCTIONS IN THE CONTEXT OF GAUSSIAN HARMONIC ANALYSIS

H. AIMAR, L. FORZANI, AND R. SCOTTO

ABSTRACT. The purpose of this paper is twofold. We introduce a general maximal function on the Gaussian setting which dominates the Ornstein-Uhlenbeck maximal operator and prove its weak type (1,1) by using a covering lemma which is halfway between Besicovitch and Wiener. On the other hand, by taking as a starting point the generalized Cauchy-Riemann equations, we introduce a new class of Gaussian Riesz Transforms. We prove, using the maximal function defined in the first part of the paper, that unlike the ones already studied, these new Riesz Transforms are weak type (1,1) independently of their orders.

1. INTRODUCTION AND MAIN RESULTS

Hermite polynomials play a central role in the context of Gaussian Harmonic Analysis. They are also the building blocks for the eigenfuctions of the harmonic oscillator in Quantum Mechanics. In this context (see [16]), let us denote by \mathcal{P} the one-dimensional momentum operator defined on a test function u as $\mathcal{P}u = -i\frac{\partial u}{\partial x}$ and by \mathcal{Q} the position operator defined by $\mathcal{Q}u = xu$. When solving the harmonic oscillator the underlying Hamiltonian is essentially given by

$$\frac{1}{2}(\mathcal{P}^2+\mathcal{Q}^2),$$

the quantum mechanical problem is then to find all the eigenvalues and eigenfunctions of the differential operator

$$\frac{1}{2}\left[-\frac{d^2}{dx^2} + x^2\right].$$

Its eigenfunctions are $e^{-x^2/2}H_k$, where H_k denotes the Hermite polynomial of degree k. They can be defined through the Rodrigues formula as

$$H_k(x) = (-1)^k e^{x^2} \frac{d^k}{dx^k} e^{-x^2}$$

for $x \in \mathbb{R}$ and k = 0, 1, ... They are also the eigenfunctions of the differential operator

$$\frac{1}{2}\frac{d^2}{dx^2} - x\frac{d}{dx} = \frac{1}{2}e^{x^2}\frac{d}{dx}(e^{-x^2}\frac{d}{dx}).$$

©2006 American Mathematical Society

Reverts to public domain 28 years from publication

Received by the editors May 5, 2004 and, in revised form, March 4, 2005.

²⁰⁰⁰ Mathematics Subject Classification. Primary 42B20, 42B25; Secondary 42C10, 47D06, 42A50, 60H07.

Key words and phrases. Gaussian measure, Fourier analysis, Fourier analysis in several variables, singular integrals, maximal functions, Hermite expansions.

The *n*-dimensional Hermite polynomial of order $\alpha = (\alpha_1, \ldots, \alpha_n) \in \mathbb{N}_0^n$ and degree $|\alpha| = \sum_{j=1}^n \alpha_j$, denoted by H_{α} , is defined as the tensor product of the one-dimensional ones,

$$H_{\alpha}(x) = \bigotimes_{j=1}^{n} H_{\alpha_j}(x_j)$$

with $x \in \mathbb{R}^n$. They are orthogonal with respect to the Gaussian measure

$$d\gamma(x) = e^{-|x|^2} dx$$

Let us consider the normalization h_{α} of H_{α} given by $h_{\alpha}(x) = \frac{H_{\alpha}(x)}{(\sqrt{\pi \alpha ! 2^{|\alpha|}})^{1/2}}$, then the set $\mathcal{F} = \{h_{\alpha}\}_{\alpha \in \mathbb{N}_{0}^{n}}$ turns out to be an orthonormal basis in $L^{2}(d\gamma)$.

Let $f \in L^2(d\gamma)$, then $f = \sum_{\alpha} a_{\alpha} h_{\alpha}$ with $a_{\alpha} = \int f h_{\alpha} d\gamma$. It can be proved that Abel's expansion $\sum_{\alpha} e^{-|\alpha|t} a_{\alpha} h_{\alpha}$ converges absolutely to the Ornstein-Uhlenbeck semigroup

$$T^{t}f(x) = \int_{\mathbb{R}^{n}} M(t, x, y) f(y) \, dy$$

for almost every $x \in \mathbb{R}^n$, where

$$M(t, x, y) = \sum_{\alpha} e^{-|\alpha|t} h_{\alpha}(x) h_{\alpha}(y)$$

= $\pi^{-n/2} (1 - e^{-2t})^{-n/2} e^{-\frac{|e^{-t}x-y|^2}{1 - e^{-2t}}}, \quad t > 0;$

M(t, x, y) is called the Mehler kernel (see [14]).

By writing $u(x,t) = T^t f(x)$, u turns out to be the solution of the parabolic partial differential equation

$$\frac{\partial u}{\partial t} = \frac{1}{2}\Delta u - x \cdot \nabla u$$

with initial data $f \in L^2(d\gamma)$.

The Ornstein-Uhlenbeck differential operator is defined by $L = \frac{1}{2}\Delta - x \cdot \nabla$, with Δ the Laplace operator and $\nabla = (\frac{\partial}{\partial x_1}, \dots, \frac{\partial}{\partial x_1})$ the gradient. Thus $T^t = e^{Lt}$. The *n*-dimensional Hermite polynomials are the eigenfunctions of L, i.e., $Lh_{\alpha} = -|\alpha|h_{\alpha}$. The Ornstein-Uhlenbeck semigroup T^t as well as L are selfadjoint operators with respect to the Gaussian measure.

If we reparametrize T^t with $r = e^{-t}$ and use the same notation for the reparametrized operator, then we have

$$T^{r}f(x) = \frac{e^{|x|^{2}}}{\pi^{1/2}(1-r^{2})^{n/2}} \int_{\mathbb{R}^{n}} e^{-\frac{|x-r|y|^{2}}{1-r^{2}}} f(y) \, d\gamma(y).$$

In 1969, C. Calderón [1] proved that the multiparametric maximal operator

$$\mathcal{T}^* f(x) = \sup_{\substack{0 < r_1 < 1 \\ 0 < r_n < 1}} \left| \frac{e^{|x|^2}}{\pi^{n/2}} \prod_{i=1}^n \frac{1}{(1 - r_i^2)^{1/2}} \int_{\mathbb{R}^n} e^{-\sum_{i=1}^n \frac{(x_i - r_i y_i)^2}{1 - r_i^2}} f(y) \, d\gamma(y) \right|$$

is bounded in $L^p(\mathbb{R}^n, d\gamma)$, p > 1. From this result, the $L^p(\mathbb{R}^n, d\gamma)$ strong type property p > 1 for the one-parameter maximal operator

$$T^*f(y) = \sup_{0 < r < 1} |T^r f(y)|$$

follows. It is worth mentioning that this result also follows from the general theory of symmetric diffusion semigroups and in this case the L^p constant obtained is independent of dimension. It is known that, for n > 1, T^* is not weak type (1,1) with respect to the Gaussian measure. The same question for T^* , with n >1, was an open problem until 1984, when P. Sjögren [22], proved that T^* is γ weak type (1,1). Sjögren's proof does not give pointwise estimates by means of average maximal operators on the global part; covering results, such as Besicovitch or Wiener Lemmas, are not used either. These are the basic classical tools used on the approximations of the identity with the Lebesgue measure and with doubling measures.

The ad hoc method developed by Sjögren is very useful and was used by other people in order to prove weak type inequalities of certain singular integral operators associated with this semigroup. But there are operators which cannot be handled likewise since their kernels exceed the bounds necessary to apply his "forbidden region" technique.

S. Pérez in [19], whose goal was to study the operators which could not be handled by Sjögren's technique, came back to the Ornstein-Uhlenbeck semigroup and gave an explicit formula for the maximal kernel of this semigroup and with that she associated the right geometry to get the weak type inequality.

Later on, P. Sjögren, with this explicit formula of the maximal kernel, gave a very simple and elegant proof of the weak-type (1,1) of T^* which can be found in [24].

In 1988 C. Gutiérrez and W. Urbina [13] came back to the problem of pointwise estimates for T^* and proved that

$$T^*f(y) \le M_{\gamma}f(x) + \max(2, |x|^n)e^{|x|^2} ||f||_{1,\gamma},$$

where

$$M_{\gamma}f(x) = \sup_{r>0} \frac{1}{\gamma(B(x,r))} \int_{B(x,r)} |f| \, d\gamma$$

is the centered Gaussian Hardy-Littlewood maximal function. By using Besicovitch's covering lemma, the γ -weak type (1,1) of M_{γ} follows. Nevertheless, this estimate does not give the weak type (1,1) inequality with respect to the Gaussian measure for T^* except for n = 1.

Since for r and x fixed, the maximum of the kernel of the operator T^r is attained at y = x/r, the centered maximal operator does not seem to be the best average maximal function to be used in order to get the γ -weak type (1,1) inequality.

We can estimate T^* by the non-centered Gaussian Hardy-Littlewood maximal function, but P. Sjögren proved in [23] that this maximal function is not weak type (1,1). The main difficulty with this maximal operator is that we cannot use Wiener's covering lemma since γ is far from being a doubling measure.

The first of the two basic goals of this article is to prove the weak type (1,1) inequality for T^* with respect to the Gaussian measure by using a covering lemma which is halfway between Besicovitch and Wiener and whose origin goes back to the Doctoral Dissertation of L. Forzani in [3]. We will prove something stronger than the weak type (1,1) inequality for T^* . First let us define

$$M_{\Phi}f(x) = \sup_{0 < r < 1} \frac{1}{\gamma((1+\delta)B(\frac{x}{r}, \frac{|x|}{r}(1-r)))} \int_{\mathbb{R}^n} \Phi\left(\frac{|x-ry|}{\sqrt{1-r^2}}\right) |f(y)| \, d\gamma(y)$$

where $\Phi : \mathbb{R}^+_o \to \mathbb{R}^+_o$ is a non-increasing function such that

$$S = \sum_{\nu \ge 1} \Phi(\frac{1}{2}(\nu - 1)) \nu^{2n} < \infty \quad \text{and} \quad \delta = \delta_{r,x} = \frac{r}{|x|(1 - r)} \min\{\frac{1}{|x|}, \sqrt{1 - r}\}.$$

Now by taking $\Phi(t) = \frac{1}{\pi^{n/2}} \exp(-t^2)$, it will be proved that

 $T^*f(x) < CM_{\Phi}f(x)$

(see the proof of Corollary 1.1 in $\S4$).

In $\S2$ we will prove the following theorem and its corollary in $\S4$.

Theorem 1.1. There exists a constant C depending only on S and n, such that for all $f \in L^1(d\gamma)$, $\lambda > 0$, we have

$$\gamma\{x \in \mathbb{R}^n : M_{\Phi}f(x) > \lambda\} \le \frac{C}{\lambda} \int_{\mathbb{R}^n} |f(y)| \, d\gamma(y)$$

i.e., $M_{\Phi}f$ is γ -weak type (1, 1).

Corollary 1.1. T^* is γ -weak type (1,1).

This result relies strongly on the following subtle covering lemma which will be proved in §4, where a polynomial growth for the overlapping of a special family of dilations for the covering balls is obtained.

Lemma 1.1. Let $A = \{x_{\alpha} : \alpha \in I\}$ be a subset of $\mathbb{R}^n \setminus \overline{B}(0, 2\zeta)$, with $\zeta > 2$ fixed and I a finite set of indices. For each $x \in A$ a number $r = r(x) \in (\frac{3}{4}, 1 - \frac{\zeta^2}{|x|^2})$ is given. Let B_j and B_j^{ν} be the balls $B(\frac{x_j}{r_j}, \frac{|x_j|}{r_j}(1-r_j))$ and $B(\frac{x_j}{r_j}, \nu\rho_j)$ respectively, with $\nu \ge 1$ and $\rho_j = \sqrt{1-r_j}$, and $\delta_j = \frac{r_j}{|x_j|(1-r_j)} \min\{\frac{1}{|x_j|}, \sqrt{1-r_j}\} = \frac{r_j}{|x_j|^2\rho_j^2}$. Then there exist a positive constant C, depending only on n, and a subset J of Isuch that

- i) $A \subset \bigcup_{j \in J} (1 + \delta_j) B_j;$ ii) $\sum_{j \in J} \chi_{B_j^{\nu}}(z) \leq C \nu^{2n}.$

On the other hand, the proof of Corollary 1.1 is based on the following lemma where we compute explicitly the Gaussian measure of a ball.

Lemma 1.2. There exists a constant C depending on n such that for all $x \in$ $\mathbb{R}^n \setminus \{0\}, r \in (1/2, 1) \text{ and } s \in (0, 1/2) \text{ the following inequality holds:}$

$$\gamma\left(B\left(\frac{x}{r}, \frac{|x|}{r}s\right)\right) \le C s^{\frac{n-1}{2}} \exp\left(-\frac{|x|^2}{r^2}(1-s)^2\right) \min\left\{\frac{1}{|x|^2}, s^{\frac{1}{2}}\right\}.$$

Let us now introduce the second problem of this paper.

At the end of the last century great efforts have been made in order to get a general singular integral theory in the context of the Ornstein-Uhlenbeck differential operator L. By analogy with the classical harmonic analysis the Gaussian Riesz Potentials were defined as $I_{\eta} = (-L)^{-\eta}$ with $\eta > 0$ over the orthogonal complement of the eigenspace associated with the eigenvalue 0. Formally,

$$I_{\eta} = \frac{1}{\Gamma(\eta)} \int_0^{\infty} t^{\eta - 1} T^t dt.$$

These operators turned out to be not weak type (1,1) (see [8]). They are indeed bounded on $L^p(d\gamma)$ for 1 but, unlike the classical Riesz Potentials, these

Licensed to Instituto de Matematica Aplicada Litoral IMAL. Prepared on Thu May 10 15:39:22 EDT 2012 for download from IP 200.9.237.254. License or copyright restrictions may apply to redistribution; see http://www.ams.org/journal-terms-of-use

do not improve integrability on the L^p scale. They do though on the $L^p LogL$ scale (see [6]).

By following [25] it is possible to define the higher order Riesz Transforms as

$$\mathcal{R}_{\alpha}f(x) = (-1)^{|\alpha|} d^{\alpha} I_{|\alpha|/2} f(x)$$

with $\alpha = (\alpha_1, \ldots, \alpha_n) \in \mathbb{N}_0^n$, $|\alpha| = \sum_{j=1}^n \alpha_j$, and $d^{\alpha} = \frac{\partial^{|\alpha|}}{\partial_{x_1}^{\alpha_1} \cdots \partial_{x_n}^{\alpha_n}}$. These operators were proved to be bounded on $L^p(d\gamma)$ by several people from different points of view; see [17], [10], [21], [26], [11], [12], [19], and [7]. But surprisingly the weak type (1, 1) case of these operators need not be true for all α . These operators are weak type (1, 1) if and only if $|\alpha| \leq 2$; see [18], [5], [2], [19], [9], and [4].

Let us go back and review the relationship between L, the infinitesimal generator of T^t , and the derivative operator d which defines the higher order Riesz Transforms \mathcal{R}_{α} through the potentials $I_{\underline{|\alpha|}}$.

B. Muckenhoupt in [18] defined in this context the Poisson integral u and its conjugate function v in L^2 through Hermite expansions, and as integral operators otherwise, so that they satisfy the following generalized Cauchy-Riemann equations

$$\begin{cases} \frac{\partial u}{\partial x} &= -\frac{\partial v}{\partial t} \\ \frac{\partial u}{\partial t} &= \frac{1}{2}e^{x^2}\frac{\partial}{\partial x}(e^{-x^2}v) \end{cases}$$

with u verifying the following second order elliptic differential equation

$$\frac{\partial^2 u}{\partial t^2} + Lu = 0.$$

In [15], K. Itô factors L out in terms of two derivative operators which are in duality with respect to the Gaussian measure, $L = \delta d$, which in the finite dimensional case d is just the usual gradient and $\delta = \frac{1}{2}e^{|x|^2}d e^{-|x|^2}$ is the *Gaussian* gradient. If we use δ instead of d in Muckenhoupt's approach, we get the new generalized Cauchy-Riemann system

$$\begin{cases} \delta \bar{u} = -\frac{\partial \bar{v}}{\partial t}, \\ \frac{\partial \bar{u}}{\partial t} = d\bar{v}, \end{cases}$$

with the function \bar{u} satisfying the second order partial differential equation

$$\frac{\partial^2 \bar{u}}{\partial t^2} + \bar{L}\bar{u} = 0,$$

where $\bar{L} = L - I = d \delta$.

From the Quantum Mechanics point of view, what we are doing is substituting the pair of operators $(\mathcal{P}, \mathcal{Q})$ by the real one $(i\mathcal{P} - 2\mathcal{Q}, i\mathcal{P}) = (\delta, d)$. Since the commutator $[\delta, d]$ is the identity operator, we again have that $\overline{L} = d\delta$.

If in the construction of the Gaussian Riesz Transforms we use δ instead of d and \overline{L} instead of L, we obtain an awesome result: the Riesz transforms associated with these new operators are all weak type (1, 1) independently of their orders.

For $1 , the <math>L^p(d\gamma)$ boundedness of these new operators follows from P. A. Meyer's Multiplier Theorem in [17], which cannot be applied to prove the weak type (1, 1) inequality.

If in \mathbb{R}^n we use the gradient

$$\delta^{\alpha} = \frac{1}{2^{|\alpha|}} e^{|x|^2} d^{\alpha} e^{-|x|^2}$$

and the Riesz potentials associated with \bar{L} , then these new singular integral operators are defined by

$$\bar{\mathcal{R}}_{\alpha}f(x) = (-1)^{|\alpha|}\delta^{\alpha}(-\bar{L})^{-|\alpha|/2}f(x).$$

The action of one of these operators over a Hermite polynomial is as follows:

(1.1)

$$\bar{\mathcal{R}}_{\alpha}H_{\beta} = \frac{(-1)^{|\alpha|}}{2^{|\alpha|}(|\beta|+1)^{|\alpha|/2}}e^{|x|^{2}}d^{\alpha}(e^{-|x|^{2}}H_{\beta}(x)) \\
= \frac{(-1)^{|\alpha+\beta|}}{2^{|\alpha|}(|\beta|+1)^{|\alpha|/2}}e^{|x|^{2}}d^{\alpha+\beta}(e^{-|x|^{2}}) \\
= \frac{1}{2^{|\alpha|}(|\beta|+1)^{|\alpha|/2}}H_{\alpha+\beta}(x).$$

On the first line we use that

$$\bar{L}H_{\beta} = -(|\beta|+1)H_{\beta}$$
 and $(-\bar{L})^{-|\alpha|/2}H_{\beta} = \frac{1}{(|\beta|+1)^{|\alpha|/2}}H_{\beta}$.

The singular integral operators $\overline{\mathcal{R}}_{\alpha}$ are weak type (1, 1) for all α . More precisely, in §3 we will prove the following theorem.

Theorem 1.2. There exists a constant C depending only on n and α such that for all $f \in L^1(d\gamma)$, $\lambda > 0$, we have

$$\gamma\{x \in \mathbb{R}^n : \bar{\mathcal{R}}_{\alpha}f(x) > \lambda\} \le \frac{C}{\lambda} \int_{\mathbb{R}^n} |f(y)| \, d\gamma(y)$$

i.e., $\overline{\mathcal{R}}_{\alpha}f$ is γ -weak type (1,1).

The main feature, in order to prove this theorem, will be to apply Theorem 1.1 with a special Φ .

2. A New maximal function M_{Φ}

In this section we will prove the γ -weak type of the operator

$$M_{\Phi}f(x) = \sup_{0 < r < 1} \frac{1}{\gamma((1+\delta)B(\frac{x}{r}, \frac{|x|}{r}(1-r)))} \int_{\mathbb{R}^n} \Phi\left(\frac{|x-ry|}{\sqrt{1-r^2}}\right) |f(y)| \, d\gamma(y)$$

where $\Phi : \mathbb{R}_o^+ \to \mathbb{R}_o^+$ is a non-increasing function such that

$$S = \sum_{\nu \ge 1} \Phi(\frac{1}{2}(\nu - 1)) \nu^{2n} < \infty \quad \text{and} \quad \delta = \delta_{r,x} = \frac{r}{|x|(1 - r)} \min\{\frac{1}{|x|}, \sqrt{1 - r}\}.$$

Proof of Theorem 1.1. We consider only $r > \frac{3}{4}$, since the maximal operator is trivially γ -weak type (1,1) for $0 < r \leq \frac{3}{4}$ (see [3]). Let us denote with the same letter M_{Φ} the maximal operator restricted to the interval $\frac{3}{4} < r < 1$, with M_{Φ}^{1} the maximal operator for $\frac{3}{4} < r < 1 - \frac{\zeta^{2}}{|x|^{2}}$ and M_{Φ}^{2} the corresponding one for $1 - \frac{\zeta^{2}}{|x|^{2}} < r < 1$. (ζ is the constant chosen in Lemma 1.1).

First we will prove that for $|x| \leq 2\zeta$, $M_{\Phi}f(x) \leq CM_{\gamma}f(x)$, where M_{γ} is the centered Gaussian Hardy-Littlewood maximal function and which is known to be

 γ -weak type (1, 1). Indeed, let us denote $R_{x,r} = \frac{|x|}{r}(1-r) + \min\{\frac{1}{|x|}, \sqrt{1-r}\}$. Then for $|x| \leq 2\zeta$,

$$\begin{split} M_{\Phi}f(x) &= \sup_{3/4 < r < 1} \frac{1}{\gamma(B(\frac{x}{r}, R_{x,r}))} \int_{\mathbb{R}^{n}} \Phi\left(\frac{|x - ry|}{\sqrt{1 - r^{2}}}\right) |f(y)| \, d\gamma(y) \\ &\leq C \sup_{3/4 < r < 1} \frac{e^{|x|^{2}}}{|B(x, R_{x,r})|} \sum_{\nu=0}^{\infty} \int_{\nu R_{x,r} \leq |y - \frac{x}{r}| \leq (\nu+1)R_{x,r}} \Phi\left(\frac{r|y - \frac{x}{r}|}{\sqrt{1 - r^{2}}}\right) \\ &\quad |f(y)| d\gamma(y) \\ &\leq C \sum_{\nu=0}^{\infty} \Phi(\nu/8\zeta)(\nu+2)^{n} \\ &\quad \sup_{3/4 < r < 1} \frac{1}{\gamma(B(x, (\nu+2)R_{x,r}))} \int_{B(x, (\nu+2)R_{x,r})} |f(y)| \, d\gamma(y) \\ &\leq C M_{\gamma} f(x). \end{split}$$

For $|x| \ge 2\zeta$, $M_{\Phi}f(x) \le M_{\Phi}^1f(x) + M_{\Phi}^2f(x)$ and the γ -weak type (1,1) of M_{Φ} will follow once we prove that both M_{Φ}^1 and M_{Φ}^2 are γ -weak type (1,1).

In order to prove the weak type (1,1) of M^1_{Φ} it is enough to prove that

$$\gamma(E_N^{1,\lambda}) \leq \frac{C}{\lambda} \int_{\mathbb{R}^n} |f(y)| \, d\gamma(y)$$

with the constant C independent of N and f, where

$$E_N^{1,\lambda} = \{x \in \mathbb{R}^n : |x| \ge 2\zeta \text{ and } M_{\Phi}^1 f(x) > \lambda\} \cap B(0,N).$$

For each $x \in E_N^{1,\lambda}$ there exists a $r = r(x) \in (\frac{3}{4}, 1 - \frac{\zeta^2}{|x|^2})$ such that

(2.1)
$$\frac{1}{\gamma((1+\delta)B(\frac{x}{r},\frac{|x|}{r}(1-r)))} \int_{\mathbb{R}^n} \Phi(\frac{|ry-x|}{\sqrt{1-r^2}}) |f(y)| d\gamma(y) \ge \lambda.$$

For every $x \in E_N^{1,\lambda}$, we have that $|x|\frac{1-r}{r}$ is bounded above and below by positive numbers, and the centers $\frac{x}{r}$ are a bounded subset of \mathbb{R}^n . Hence, there exists $\epsilon > 0$ such that for all $0 < \alpha < 1$,

$$\gamma\left(B(\frac{x}{r},(1+\alpha)|x|\frac{1-r}{r}+\epsilon)\right) \le 2\gamma\left(B(\frac{x}{r},(1+\alpha)|x|\frac{1-r}{r})\right)$$

for all $x \in E_N^{1,\lambda}$. Let A be a subset of $E_N^{1,\lambda}$ which is a maximal set with the property $|x - \bar{x}| > \frac{\epsilon}{2}$ for $x \neq \bar{x}, x \in A, \bar{x} \in A$. Since $E_N^{1,\lambda}$ is bounded, A is a finite set $A = \{y_1, \ldots, y_L\}$. If we apply Lemma 1.1 to the set A, we get a family of balls $\left\{B_j = B(\frac{x_j}{r_j}, |x_j| \frac{1-r_j}{r_j})\right\}_{j \in J \subset \{1, \ldots, L\}}$ such that $A \subset \bigcup_{j \in J} (1 + \delta_j) B_j$ and ii) of Lemma 1.1 also holds. Thus,

$$E_N^{1,\lambda} \subset \bigcup_{j \in J} B\left(\frac{x_j}{r_j}, (1+\delta_j)\frac{|x_j|}{r_j}(1-r_j)+\epsilon\right).$$

Then

$$\gamma(E_N^{1,\lambda}) \le 2\sum_{j \in J} \gamma\left(B(\frac{x_j}{r_j}, (1+\delta_j)|x_j|\frac{1-r_j}{r_j})\right).$$

From (2.1), since Φ is a non-increasing function such that

$$\sum_{v \ge 1} \Phi(\frac{1}{2}(v-1))v^{2n} < \infty,$$

we have, using ii) of Lemma 1.1, that

$$\begin{split} \gamma(E_N^{1,\lambda}) &\leq 2\sum_{j\geq 1} \gamma((1+\delta_j)B_j) \\ &\leq \frac{C}{\lambda}\sum_{j\geq 1} \int_{\mathbb{R}^n} \Phi\left(\frac{|r_jy-x_j|}{(1-r_j^2)^{1/2}}\right) |f(y)| \, d\gamma(y) \\ &\leq \frac{C}{\lambda}\sum_{j\geq 1}\sum_{\nu\geq 1} \int_{B(\frac{x_j}{r_j},\nu\rho_j)-B(\frac{x_j}{r_j},(\nu-1)\rho_j)} \Phi\left(\frac{|r_jy-x_j|}{(1-r_j^2)^{1/2}}\right) |f(y)| \, d\gamma(y) \\ &= \frac{C}{\lambda}\sum_{j\geq 1}\sum_{\nu\geq 1} \Phi(\frac{1}{2}(\nu-1)) \int_{B_j^{\nu}} |f(y)| \, d\gamma(y) \\ &\leq \frac{C}{\lambda} \int_{\mathbb{R}^n}\sum_{\nu\geq 1} \Phi(\frac{1}{2}(\nu-1)) \sum_j \chi_{B_j^{\nu}}(y) |f(y)| \, d\gamma(y) \\ &\leq \frac{C}{\lambda} \int_{\mathbb{R}^n}\sum_{\nu\geq 1} \Phi(\frac{1}{2}(\nu-1)) \nu^{2n} |f(y)| \, d\gamma(y) \\ &\leq \frac{C}{\lambda} \int_{\mathbb{R}^n} |f(y)| \, d\gamma(y). \end{split}$$

Now, we will prove that M_{Φ}^2 is weak type (1,1). First, let us observe that, if $r > 1 - \frac{\zeta^2}{|x|^2}$, then, for all $y \in (1+\delta)B(\frac{x}{r}, \frac{|x|}{r}(1-r)) = B(\frac{x}{r}, \frac{|x|}{r}(1-r) + \sqrt{1-r})$, the values of $e^{-|y|^2}$ are equivalent.

Now, let us define

$$E_N^{2,\lambda} = \{ x \in \mathbb{R}^n : |x| \ge 2\zeta \text{ and } M_{\Phi}^2 f(x) > \lambda \} \cap B(0,N).$$

The weak type (1,1) for M_{Φ}^2 follows once we prove the inequality

(2.2)
$$\gamma(E_N^{2,\lambda}) \le \frac{C}{\lambda} ||f||_{1,\gamma}$$

with C being a constant independent of N and f. For each $x \in E_N^{2,\lambda}$, we have $\gamma((1+\delta)B(\frac{x}{r},\frac{|x|}{r}(1-r))) \simeq e^{-|x|^2}(1-r)^{\frac{n}{2}}$. To prove (2.2), we will divide the integral in $M_{\Phi}^2 f$ into two parts: one given by $|y-x| < 2\frac{C}{|x|r}$ and the other one by $|y - x| > 2\frac{C}{|x|r}$.

For the first region we have

$$(2.3) \qquad \frac{e^{|x|^2}}{(1-r)^{\frac{n}{2}}} \int_{|y-x|<2\frac{C}{|x|r}} \Phi\left(\frac{|ry-x|}{\sqrt{1-r^2}}\right) |f(y)| \, d\gamma(y) \\ \leq C \frac{e^{|x|^2}}{(1-r)^{\frac{n}{2}}} \int_{|y-x|$$

with M_T the truncated non-centered Gaussian maximal function defined by

$$M_T f(x) = \sup_{\substack{x \in B(y,t) \\ 0 < t < \min\{1, \frac{1}{|x|}\}}} \frac{e^{|x|^2}}{\omega_n t^n} \int_{B(y,t)} |f(z)| \, d\gamma(z),$$

where ω_n is the volume of the unit ball in \mathbb{R}^n . The first inequality follows from the fact that Φ is bounded and $|ry - x| \ge r|y - x| - (1 - r)|x| \ge \frac{r}{2}|y - x|$. The second inequality follows from the fact that Φ is a Lebesgue integrable, non-increasing function and hence, it is a good approximation of the identity.

The truncated non-centered Gaussian maximal function is bounded by the centered Gaussian Hardy-Littlewood maximal function and therefore is γ -weak type (1,1).

For the second region, we have that |ry - x| > C|y - x|. Therefore,

$$\begin{aligned} \frac{1}{(1-r^2)^{\frac{n}{2}}} \Phi\left(\frac{|ry-x|}{\sqrt{1-r^2}}\right) &\leq & \frac{1}{|y-x|^n} \left(\frac{|y-x|}{\sqrt{1-r^2}}\right)^n \Phi\left(\frac{c|y-x|}{\sqrt{1-r^2}}\right) \\ &\leq & C \frac{(\sqrt{1-r^2})^n}{|y-x|^{2n}} \\ &\leq & \frac{C}{|x|^n |y-x|^{2n}}, \end{aligned}$$

since $\Phi(\frac{1}{2}(\nu - 1)) \le S$. Then,

$$\begin{split} & \frac{e^{|x|^2}}{(1-r)^{\frac{n}{2}}} \int_{|y-x|>2\frac{c}{|x|r}} \Phi\left(\frac{|ry-x|}{\sqrt{1-r^2}}\right) |f(y)| \, d\gamma(y) \\ & \leq \quad C \frac{e^{|x|^2}}{|x|^n} \int_{|y-x|>2\frac{c}{|x|}} \frac{|f(y)|}{|y-x|^{2n}} \, d\gamma(y), \end{split}$$

but

$$\frac{e^{|x|^2}}{|x|^n} \int_{|y-x|>2\frac{c}{|x|}} \frac{|f(y)|}{|y-x|^{2n}} \, d\gamma(y) \in L^1(d\gamma).$$

So, inequality (2.2) follows.

3. New higher order Gaussian Riesz Transforms

The new higher order Gaussian Riesz Transforms are defined as

$$\bar{\mathcal{R}}_{\alpha}f(x) = \text{p.v. } e^{|x|^2} \int_{\mathbb{R}^n} \bar{\mathcal{K}}_{\alpha}(x,y)f(y) \, d\gamma(y)$$

where

$$\bar{\mathcal{K}}_{\alpha}(x,y) = C_{\alpha} \int_{0}^{1} \left(\frac{-\log r}{1-r^{2}}\right)^{\frac{|\alpha|-2}{2}} H_{\alpha}\left(\frac{x-ry}{\sqrt{1-r^{2}}}\right) \frac{e^{-\frac{|x-ry|^{2}}{1-r^{2}}}}{(1-r^{2})^{\frac{n}{2}+1}} \, dr.$$

Licensed to Instituto de Matematica Aplicada Litoral IMAL. Prepared on Thu May 10 15:39:22 EDT 2012 for download from IP 200.9.237.254. License or copyright restrictions may apply to redistribution; see http://www.ams.org/journal-terms-of-use

Formally $\bar{\mathcal{K}}_{\alpha}$ is obtained by differentiating with the dual derivative the kernel corresponding to the Riesz potentials associated with \bar{L} ,

$$\begin{aligned} (-\bar{L})^{-|\alpha|/2} f(x) &= \frac{1}{\Gamma(|\alpha|/2)} \int_0^1 (-\log r)^{\frac{|\alpha|}{2} - 1} T^r f(x) dr \\ &= \tilde{C}_\alpha e^{|x|^2} \int_{\mathbb{R}^n} \left(\int_0^1 (-\log r)^{\frac{|\alpha| - 2}{2}} \frac{e^{-\frac{|x - ry|^2}{1 - r^2}}}{(1 - r^2)^{\frac{n}{2}}} dr \right) f(y) \, d\gamma(y). \end{aligned}$$

Proof of Theorem 1.2. For each $x \in \mathbb{R}^n$ we view this operator as the sum of two operators which are obtained, as is usual in this context, by splitting \mathbb{R}^n into a local part, B_x , the Euclidean ball centered at x and radius min $(1, \frac{1}{|x|})$, and its complement called the global part. Thus,

$$\bar{\mathcal{R}}_{\alpha}f(x) = \bar{\mathcal{R}}_{\alpha,l}f(x) + \bar{\mathcal{R}}_{\alpha,g}f(x)$$

where $\bar{\mathcal{R}}_{\alpha,l}f(x) = \bar{\mathcal{R}}_{\alpha}(f\mathcal{X}_{B_x})$ and

$$\bar{\mathcal{R}}_{\alpha,g}f(x) = \bar{\mathcal{R}}_{\alpha}(f(1 - \mathcal{X}_{B_x}))$$

We will prove that these two operators are γ -weak type (1,1) and so will be $\bar{\mathcal{R}}_{\alpha}$.

In order to prove that $\overline{\mathcal{R}}_{\alpha,l}$ is γ -weak type (1,1) we state the following theorem whose proof can be found in either [6] or [20].

Theorem 3.1. Let $\mathcal{K}(x, y)$ be a C^1 function off the diagonal of $\mathbb{R}^n \times \mathbb{R}^n$ which satisfies

$$|\mathcal{K}(x,y)| \le \frac{C}{|x-y|^n} \quad and \quad |D_y\mathcal{K}(x,y)| \le \frac{C}{|x-y|^{n+1}}$$

for $|x - y| \leq \min(1, \frac{1}{|x|})$, and the principal value of the integral operator T with kernel \mathcal{K} is bounded on $L^p(d\gamma)$ for some $1 . Then <math>T_l$, defined as $T_l(f)(x) = T(fB_x)(x)$, is γ -weak type (1, 1).

In our case

$$Tf(x) = p.v. \int \mathcal{K}(x, y)f(y)dy$$

with

$$\begin{aligned} \mathcal{K}(x,y) &= e^{|x|^2} \bar{\mathcal{K}}_{\alpha}(x,y) e^{-|y|^2} \\ &= C_{\alpha} \int_0^1 \left(\frac{-\log r}{1-r^2}\right)^{\frac{|\alpha|-2}{2}} H_{\alpha}\left(\frac{x-ry}{\sqrt{1-r^2}}\right) \frac{e^{-\frac{|rx-y|^2}{1-r^2}}}{(1-r^2)^{\frac{n}{2}+1}} dr \end{aligned}$$

and, therefore,

$$\frac{\partial \mathcal{K}}{\partial y_j}(x,y) = 2C_{\alpha} \int_0^1 \left(\frac{-\log r}{1-r^2}\right)^{\frac{|\alpha|-2}{2}} \left[\frac{-r\alpha_j}{\sqrt{1-r^2}} H_{\alpha-e_j}\left(\frac{x-ry}{\sqrt{1-r^2}}\right) \frac{(x-ry)}{\sqrt{1-r^2}} + H_{\alpha}\left(\frac{x-ry}{\sqrt{1-r^2}}\right) \frac{(rx_j-y_j)}{1-r^2} \left[\frac{e^{-\frac{|rx-y|^2}{1-r^2}}}{(1-r^2)^{\frac{n}{2}+1}} dr.$$

In the following two claims we will show that the hypotheses of Theorem 3.1 are fulfilled for this operator.

Claim 1. On
$$B_x$$
, $|\mathcal{K}(x,y)| \leq \frac{C}{|x-y|^n}$ and $|\frac{\partial \mathcal{K}}{\partial y_j}(x,y)| \leq \frac{C}{|x-y|^{n+1}}$

Proof. For every $y \in B_x$ there exists a constant C > 0 such that $C^{-1} \le e^{|y|^2 - |x|^2} \le C$, then

$$|\mathcal{K}(x,y)| \le C|e^{-|x|^2 + |y|^2} \mathcal{K}(x,y)| = C|\mathcal{K}_{\alpha}(x,y)|$$

and

$$\left|\frac{\partial \mathcal{K}}{\partial y_j}(x,y)\right| \le C \left|e^{-|x|^2 + |y|^2} \frac{\partial \mathcal{K}}{\partial y_j}(x,y)\right|.$$

On the other hand, on B_x ,

$$e^{-c\frac{|x-ry|^2}{1-r^2}} = e^{-c\frac{|x-y|^2}{1-r^2}} e^{-c\frac{1-r}{1+r}|y|^2} e^{-c\frac{(x-y)\cdot y}{1-r}} \le C e^{-c\frac{|x-y|^2}{1-r}},$$

and thus with this inequality and taking into account that $t^m e^{-ct^2} \leq C, \ \forall t \geq 0$, we get

$$\left| H_{\alpha}\left(\frac{x-ry}{\sqrt{1-r^2}}\right) \right| e^{-\frac{|x-ry|^2}{1-r^2}} \le C \sum_{m=0}^{|\alpha|} \left| \frac{x-ry}{\sqrt{1-r^2}} \right|^m e^{-\frac{|x-ry|^2}{2(1-r^2)}} e^{-\frac{|x-ry|^2}{2(1-r^2)}} \le C e^{-c\frac{|x-y|^2}{1-r}}$$

Therefore, by combining all the above remarks, on B_x we have

$$\begin{aligned} |\mathcal{K}(x,y)| &\leq C \int_0^1 \left(\frac{-\log r}{1-r^2}\right)^{\frac{|\alpha|-2}{2}} \frac{e^{-c\frac{|x-y|^2}{1-r}}}{(1-r)^{\frac{n}{2}+1}} \, dr \\ &\leq C \left[\int_0^{\frac{1}{2}} (-\log r)^{\frac{|\alpha|-2}{2}} \, dr + \int_{\frac{1}{2}}^1 \frac{e^{-c\frac{|x-y|^2}{1-r}}}{(1-r)^{\frac{n}{2}+1}} \, dr\right] \\ &\leq C \left(1 + \frac{1}{|x-y|^n}\right) \leq \frac{C}{|x-y|^n} \end{aligned}$$

and

$$\begin{aligned} \left| \frac{\partial \mathcal{K}}{\partial y_j}(x,y) \right| &\leq C \int_0^1 \left(\frac{-\log r}{1-r^2} \right)^{\frac{|\alpha|-2}{2}} \frac{e^{-c\frac{|x-y|^2}{1-r}}}{(1-r)^{\frac{n+3}{2}}} dr \\ &\leq C \left[\int_0^{\frac{1}{2}} (-\log r)^{\frac{|\alpha|-2}{2}} dr + \int_{\frac{1}{2}}^1 \frac{e^{-c\frac{|x-y|^2}{1-r}}}{(1-r)^{\frac{n+3}{2}}} dr \right] \\ &\leq C \left(1 + \frac{1}{|x-y|^{n+1}} \right) \leq \frac{C}{|x-y|^{n+1}}. \end{aligned}$$

Claim 2. The operator $T = \overline{\mathcal{R}}_{\alpha}$ is bounded on $L^2(d\gamma)$.

Proof. Let $f \in L^2(d\gamma)$, $f = \sum_{\beta} a_{\beta} h_{\beta}$ with $a_{\beta} = \int f h_{\beta} d\gamma$. From the action of $\bar{\mathcal{R}}_{\alpha}$ over Hermite polynomials (1.1),

$$\bar{\mathcal{R}}_{\alpha}h_{\beta}(x) = \left[\frac{\prod_{j=1}^{n}\prod_{k=0}^{\alpha_{j}-1}(\beta_{j}+\alpha_{j}-k)}{(2(|\beta|+1))^{|\alpha|}}\right]^{\frac{1}{2}}h_{\beta+\alpha}(x)$$

Licensed to Instituto de Matematica Aplicada Litoral IMAL. Prepared on Thu May 10 15:39:22 EDT 2012 for download from IP 200.9.237.254. License or copyright restrictions may apply to redistribution; see http://www.ams.org/journal-terms-of-use and, therefore,

$$\begin{aligned} \|\bar{\mathcal{R}}_{\alpha}f\|_{L^{2}(d\gamma)}^{2} &= \sum_{\beta} \frac{\prod_{j=1}^{n} \prod_{k=0}^{\alpha_{j}-1} (\beta_{j} + \alpha_{j} - k)}{(2(|\beta| + 1))^{|\alpha|}} |a_{\beta}|^{2} \\ &\leq \sum_{\beta} \prod_{j=1}^{n} (1 + \alpha_{j})^{\alpha_{j}} |a_{\beta}|^{2} \\ &\leq (1 + |\alpha|)^{|\alpha|} \sum_{\beta} |a_{\beta}|^{2} = C \|f\|_{L^{2}(d\gamma)}^{2}. \end{aligned}$$

Now if we apply these two claims to Theorem 3.1, the γ -weak type (1,1) of $\overline{\mathcal{R}}_{\alpha,l}$ follows.

In order to prove that $\overline{\mathcal{R}}_{\alpha,g}$ is also γ -weak type (1,1) we will prove

Claim 3. On $\mathbb{R}^n \setminus B_x$, $|\bar{\mathcal{R}}_{\alpha,g}f(x)| \leq CM_{\Phi}f(x)$ with $\Phi(t) = e^{-ct^2}$.

This together with Theorem 1.1 give the weak type (1,1) inequality for $\overline{\mathcal{R}}_{\alpha,q}$.

Proof of Claim 3.

$$\begin{split} |\bar{\mathcal{K}}_{\alpha}(x,y)| &= \left| \int_{0}^{1} \left(\frac{-\log r}{1-r^{2}} \right)^{\frac{|\alpha|-2}{2}} H_{\alpha} \left(\frac{x-ry}{\sqrt{1-r^{2}}} \right) \frac{e^{-\frac{|x-ry|^{2}}{1-r^{2}}}}{(1-r^{2})^{\frac{n}{2}+1}} \, dr \right| \\ &\leq C \int_{0}^{\frac{3}{4}} (-\log r)^{\frac{|\alpha|-2}{2}} \frac{e^{-\frac{|x-ry|^{2}}{2(1-r^{2})}}}{(1-r^{2})^{\frac{n}{2}}} \, dr \\ &+ C \int_{\frac{3}{4}}^{1-\zeta/|x|^{2}} \frac{e^{-\frac{|x-ry|^{2}}{2(1-r^{2})}}}{(1-r^{2})^{\frac{n-1}{2}}} \left(|x| \vee (1-r^{2})^{-\frac{1}{2}} \right) \frac{dr}{|x|(1-r^{2})^{3/2}} \\ &+ C \int_{1-\zeta/|x|^{2}}^{1} \frac{e^{-c\frac{|x-ry|^{2}}{1-r^{2}}}}{(1-r^{2})^{\frac{n-1}{2}}} \left(|x| \vee (1-r^{2})^{-\frac{1}{2}} \right) \frac{e^{-\overline{c}\frac{|x-y|^{2}}{1-r}}}{1-r} \, dr \\ &= C \left(\bar{\mathcal{K}}_{\alpha}^{1}(x,y) + \bar{\mathcal{K}}_{\alpha}^{2}(x,y) + \bar{\mathcal{K}}_{\alpha}^{3}(x,y) \right) \end{split}$$

where the inequality is obtained by annihilating the Hermite polynomial with part of the exponential, then splitting the unit interval of the integral into three subintervals [0, 3/4], $[3/4, 1 - \zeta/|x|^2]$, and $[1 - \zeta/|x|^2, 1]$ and taking into account that on the second one $|x| \vee (1-r^2)^{-1/2} \ge |x|$, on the third one $|x| \vee (1-r^2)^{-1/2} \ge (1-r^2)^{-1/2}$ and $|x - ry| \ge \bar{c}|x - y|$, and on the last two intervals $-\log r/(1-r^2)$ is bounded by a constant.

Thus, by using the definition of kernels $\bar{\mathcal{K}}^j_{\alpha}$ with j = 1, 2, 3, interchanging the order of integration on each operator $\bar{\mathcal{R}}^j_{\alpha,g}$ with j = 1, 2, 3, using Lemma 1.2 and

setting in this context $\Phi(t) = e^{-ct^2}$, we get

$$\begin{split} \bar{\mathcal{R}}^{1}_{\alpha,g}f(x) &= e^{|x|^{2}} \int_{\mathbb{R}^{n}} \bar{\mathcal{K}}^{1}_{\alpha}(x,y) |f(y)| d\gamma(y) \\ &= e^{|x|^{2}} \int_{\mathbb{R}^{n}} \int_{0}^{\frac{3}{4}} (-\log r)^{\frac{|\alpha|-2}{2}} \frac{e^{-\frac{|x-ry|^{2}}{2(1-r^{2})}}}{(1-r^{2})^{\frac{n}{2}}} dr |f(y)| d\gamma(y) \\ &= \int_{0}^{\frac{3}{4}} (-\log r)^{\frac{|\alpha|-2}{2}} e^{|x|^{2}} \int_{\mathbb{R}^{n}} \frac{e^{-\frac{|x-ry|^{2}}{2(1-r^{2})}}}{(1-r^{2})^{\frac{n}{2}}} |f(y)| d\gamma(y) dr \\ &\leq C \int_{0}^{\frac{3}{4}} (-\log r)^{\frac{|\alpha|-2}{2}} dr \quad M_{\Phi}f(x) \\ &\leq C M_{\Phi}f(x), \end{split}$$

$$\begin{split} \bar{\mathcal{R}}_{\alpha,g}^{2}f(x) &= e^{|x|^{2}} \int_{\mathbb{R}^{n}} \bar{\mathcal{K}}_{\alpha}^{2}(x,y) |f(y)| d\gamma(y) \\ &= e^{|x|^{2}} \int_{\mathbb{R}^{n}} \int_{\frac{3}{4}}^{1-\zeta/|x|^{2}} \frac{e^{-\frac{|x-ry|^{2}}{2(1-r^{2})}}}{(1-r^{2})^{\frac{n-1}{2}}} \left(|x| \vee (1-r^{2})^{-\frac{1}{2}}\right) \frac{dr}{|x|(1-r^{2})^{3/2}} \\ &= \int_{3/4}^{1-\zeta/|x|^{2}} e^{|x|^{2}} \int_{\mathbb{R}^{n}} \frac{e^{-c\frac{|x-ry|^{2}}{(1-r^{2})}}}{(1-r^{2})^{(n-1)/2}} (|x| \vee (1-r^{2})^{-1/2}) |f(y)| d\gamma(y) \\ &\qquad \frac{dr}{|x|(1-r^{2})^{3/2}} \\ &\leq C \frac{1}{|x|} \int_{3/4}^{1-\zeta/|x|^{2}} \frac{dr}{(1-r)^{3/2}} M_{\Phi}f(x) \\ &\leq C M_{\Phi}f(x), \end{split}$$

and, finally,

$$\begin{split} \bar{\mathcal{R}}^{3}_{\alpha,g}f(x) &= e^{|x|^{2}} \int_{\mathbb{R}^{n}} \bar{\mathcal{K}}^{3}_{\alpha}(x,y) \; |f(y)| \; d\gamma(y) \\ &= e^{|x|^{2}} \int_{\mathbb{R}^{n}} \int_{1-\zeta/|x|^{2}}^{1} \frac{e^{-c\frac{|x-ry|^{2}}{1-r^{2}}}}{(1-r^{2})^{\frac{n-1}{2}}} \; (|x| \lor (1-r^{2})^{-\frac{1}{2}}) \\ &\quad \frac{e^{-\bar{c}\frac{|x-y|^{2}}{1-r}}}{1-r} \; dr \; |f(y)| \; d\gamma(y) \\ &= \int_{1-\zeta/|x|^{2}}^{1} e^{|x|^{2}} \int_{\mathbb{R}^{n}} \frac{e^{-c\frac{|x-ry|^{2}}{(1-r^{2})}}}{(1-r^{2})^{(n-1)/2}} (|x| \lor (1-r^{2})^{-1/2}) \\ &\quad \frac{e^{-\bar{c}\frac{|x-y|^{2}}{1-r}}}{1-r} \; |f(y)| \; d\gamma(y) \; dr \\ &\leq \int_{1-\zeta/|x|^{2}}^{1} e^{|x|^{2}} \int_{\mathbb{R}^{n}} \frac{e^{-c\frac{|x-ry|^{2}}{(1-r^{2})}}}{(1-r^{2})^{(n-1)/2}} (|x| \lor (1-r^{2})^{-1/2}) \end{split}$$

$$\frac{1}{|x-y|^2} |f(y)| \, d\gamma(y) \, dr$$

$$\leq C|x|^2 \int_{1-\zeta/|x|^2}^1 dr \, M_{\Phi}f(x)$$

$$\leq C \, M_{\Phi}f(x).$$

So, since $|\bar{\mathcal{R}}_{\alpha,g}f(x)| \leq C_{\alpha} \sum_{j=1}^{3} \bar{\mathcal{R}}_{\alpha,g}^{j}f(x)$, Claim 3 holds.

4. Proof of Lemmas 1.1 and 1.2 and Corollary 1.1

Proof of Lemma 1.1. Let $I_1 = I$, $\alpha_1 \in I_1$ such that $|x_{\alpha_1}| = \min\{|x_{\alpha}| : \alpha \in I_1\}$. Let $x_1 = x_{\alpha_1}$ and $B_1 = B_{\alpha_1}$. Let $I_1, \ldots, I_{k-1}; x_1, \ldots, x_{k-1}; B_1, \ldots, B_{k-1}$ be chosen; we define $I_k = \{\alpha \in I : x_\alpha \notin \bigcup_{j=1}^{k-1} (1+\delta_j)B_j\}$, and we choose $\alpha_k \in I_k$ such that $|x_{\alpha_k}| = \min\{|x_{\alpha}| : \alpha \in I_k\}$. Let $x_k = x_{\alpha_k}$ and $B_k = B_{\alpha_k}$. Let $J = \{\alpha_1, \ldots, \alpha_N\}$ where N is the first integer for which $I_{N+1} = \emptyset$. Then i) is immediate. Before proving ii) let us make some remarks.

(1) x_j was chosen so that $x_j \notin (1 + \delta_s)B_s$ for all s < j. Hence

$$\frac{\left|\frac{x_s}{r_s} - x_j\right|^2 = \frac{|x_s|^2}{r_s^2} + |x_j|^2 - 2|x_j|\frac{|x_s|}{r_s}\cos\langle\frac{x_s}{r_s}, x_j\rangle \\ \ge R_s^2(1+\delta_s)^2;$$

- (2) $|x_j| \ge |x_s|$ for s < j; i.e., $|x_j|$ is increasing with j;
- (3) $|\frac{x_s}{r_s} \frac{x_j}{r_j}|^2 \ge \frac{1}{r_j} [\frac{|x_s|^2}{r_s^2 r_j} (r_j r_s)^2 + 2\frac{(1-r_s)}{r_s}] \ge \theta^2 \max^2(\rho_j, \rho_s)$ for s < j. In fact, using (1) and (2), and that $2R_s^2 \delta_s = 2\frac{(1-r_s)}{r_s}$,

$$\begin{split} \frac{x_s}{r_s} - \frac{x_j}{r_j} \Big|^2 &= \frac{|x_s|^2}{r_s^2} + \frac{|x_j|^2}{r_j^2} - 2\frac{|x_j|}{r_j} \frac{|x_s|}{r_s} \cos\langle \frac{x_s}{r_s}, \frac{x_j}{r_j} \rangle \\ &\geq \frac{|x_s|^2}{r_s^2} + \frac{|x_j|^2}{r_j^2} + \frac{1}{r_j} [R_s^2 (1 + \delta_s)^2 - |x_j|^2 - \frac{|x_s|^2}{r_s^2}] \\ &= \frac{1}{r_j} R_s^2 (1 + \delta_s)^2 - \frac{1}{r_j} [-|x_j|^2 (\frac{1 - r_j}{r_j}) + \frac{|x_s|^2}{r_s^2} (1 - r_j)] \\ &\geq \frac{1}{r_j} [R_s^2 (1 + \delta_s)^2 + |x_s|^2 (1 - r_j) [\frac{1}{r_j} - \frac{1}{r_s^2}]] \\ &\geq \frac{1}{r_j} [\frac{|x_s|^2}{r_s^2} (1 - r_s)^2 + 2\frac{(1 - r_s)}{r_s} + |x_s|^2 (1 - r_j) [\frac{1}{r_j} - \frac{1}{r_s^2}]] \\ &= \frac{1}{r_j} [\frac{|x_s|^2}{r_s^2 r_j} (r_j - r_s)^2 + 2\frac{(1 - r_s)}{r_s}] \\ &\geq \theta^2 \max^2(\rho_j, \rho_s). \end{split}$$

To obtain the last inequality we consider two cases:

i) $\rho_s^2 \ge \frac{1}{2}\rho_j^2$. Because of $\rho_s^2 = (1 - r_s)$ and the nonnegativity of the first term, the inequality follows.

ii) $\rho_i^2 \ge 2\rho_s^2$. We have that $(r_j - r_s)^2 = (\rho_j^2 - \rho_s^2)^2 \ge \frac{1}{4}\rho_j^4$. Using the fact that $|x_s|\rho_s \geq \zeta$ the inequality follows. (Recall that by hypothesis $r_s \leq 1 - \frac{\zeta^2}{|x_s|^2}$.) Now in order to prove ii) we define

$$I_1 = \{j : j \in J \text{ and } \nu \rho_j \ge \kappa\},$$

$$I_2 = \{j : j \in J \text{ and } \frac{R_j}{2} < \nu \rho_j < \kappa\}$$

$$I_3 = \{j : j \in J \text{ and } \nu \rho_j \le \frac{R_j}{2}\},$$

where $R_j = \frac{|x_j|}{r_j}(1 - r_j)$. We will prove

(4.1)
$$\sum_{j \in I_i} \chi_{B_j^{\nu}}(z) \le C\nu^{2n} \quad \text{for } i = 1, 2, 3,$$

from which ii) follows.

Now we prove (4.1). Let us consider $I_1(z) = \{j \in I_1 : z \in B_j^{\nu}\}$. In order to obtain the desired estimate all we need is to find a sequence of pairwise disjoint measurable sets $\{S_j\}_{j \in I_1(z)}$ such that

- (4) $S_j \subset B(z, C\nu);$ (5) $|S_j| \ge \frac{C}{\nu^n}$ for some constant C.
- The case i = 1 in (4.1) follows from (4) and (5).

We define $S_j = B(\frac{x_j}{r_i}; \frac{\theta}{2}\rho_j)$; (5) is immediate since $j \in I_1$ which implies $\rho_j \geq \frac{\kappa}{\nu}$. In order to get (4) let us take $h \in S_j$. Since $z \in B_j^{\nu}$, we have

$$|h-z| \leq |h - \frac{x_j}{r_j}| + |\frac{x_j}{r_j} - z|$$

$$\leq \frac{\theta}{2}\rho_j + \nu\rho_j$$

$$\leq C\nu.$$

That $\{S_j\}_{j \in I_1(z)}$ is a family of pairwise disjoint sets follows from (3).

Now, consider $I_2(z) = \{j \in I_2 : z \in B_j^{\nu}\}$. In order to obtain the desired estimate we just need to find a sequence of pairwise disjoint measurable sets $\{S_j\}_{j \in I_2(z)}$ such that

(6) $S_j \subset B(z, C\nu^2);$

(7) $|S_j| \simeq 1$ for some constant C.

The case i = 2 in (4.1) follows from (6) and (7).

We define $S_j = B(z + (\frac{x_j}{r_j} - z) \frac{|x_j|}{r_j}, C)$, therefore, (7) is immediate. Let us prove (6). Take $h \in S_j$, using the fact that $z \in B_j^{\nu}$ and $\frac{R_j}{2} \leq \nu \rho_j$ or equivalently $\rho_j \frac{|x_j|}{r_i} \leq 2\nu$, we get

$$\begin{split} |h - z| &\leq C + |(\frac{x_j}{r_j} - z)| \frac{|x_j|}{r_j} \\ &\leq C + \nu \rho_j \frac{|x_j|}{r_j} \\ &\leq C + 2\nu^2 \\ &\leq C\nu^2. \end{split}$$

To prove that the S_j are pairwise disjoint we will use consecutively the following facts:

i)
$$\left|\frac{x_s}{r_s} - \frac{x_j}{r_j}\right| \ge \theta \rho_s;$$

ii) $\left|x_s\right| \rho_s \ge \sqrt{\zeta};$
iii) $\left|\frac{x_j}{r_j} - z\right| \le \nu \rho_j \le \kappa \ (j \in I_2);$ and
iv) $\left|\frac{|x_s|}{r_s} - \frac{|x_j|}{r_j}\right| \le \left|\frac{x_s}{r_s} - \frac{x_j}{r_j} \le \left|\nu \rho_s + \nu \rho_j \le 2\kappa.$
So

$$\begin{vmatrix} \left(\frac{x_j}{r_j} - z\right) \frac{|x_j|}{r_j} - \left(\frac{x_s}{r_s} - z\right) \frac{|x_s|}{r_s} \end{vmatrix} & \geq \quad \frac{|x_s|}{r_s} \left| \frac{x_j}{r_j} - \frac{x_s}{r_s} \right| - \left| \frac{x_j}{r_j} - z \right| \left| \frac{|x_s|}{r_s} - \frac{|x_j|}{r_j} \right| \\ & \geq \quad \sqrt{\zeta} \quad \frac{\theta}{r_s} - 2 \kappa^2 \\ & \geq \quad C \end{aligned}$$

after choosing ζ and κ properly.

Finally, consider $I_3(z) = \{j \in I_3 : z \in B_j^{\nu}\}$. In order to obtain the desired estimate all we need is to find a sequence of pairwise disjoint measurable sets $\{S_j\}_{j \in I_3(z)}$ such that

 $\begin{array}{ll} (8) & S_j \subset B(z, C\nu\rho_{\tau}); \\ (9) & |S_j| \geq C\rho_{\tau}^n \text{ for some constant } C, \text{ where } \rho_{\tau} = \rho_{\min\{j:\, j \in I_3(z)\}}. \end{array}$

The case i = 3 in (4.1) follows from (8) and (9). We define $S_j = B(\frac{x_j}{r_j}, \frac{\theta}{2}\rho_j)$. We will prove that

(4.2)
$$\frac{1}{2}\rho_{\tau} \le \rho_j \le \frac{3}{2}\rho_{\tau} \quad \text{for all} \quad j \in I_3(z).$$

From (4.2) we have (8) and (9). That S_j are disjoint follows from (3).

Let us prove (4.2). From (3), $|\frac{x_{\tau}}{r_{\tau}} - \frac{x_j}{r_j}|^2 \ge \frac{1}{r_j} \left[\frac{|x_{\tau}|^2}{r_{\tau}^2 r_j} (r_j - r_{\tau})^2 + 2\frac{(1 - r_{\tau})}{r_{\tau}} \right]$, then $|\frac{x_j}{r_j} - \frac{x_{\tau}}{r_{\tau}}|^2 \ge \frac{|x_{\tau}|^2}{r_{\tau}^2} (\rho_j^2 - \rho_{\tau}^2)^2$ and since $\tau, j \in I_3(z)$, we have $\frac{|x_{\tau}|}{r_{\tau}} \ge 2\frac{\nu}{\rho_{\tau}}$. Therefore,

$$\begin{aligned} (\nu \rho_j + \nu \rho_\tau)^2 &\geq |\frac{x_j}{r_j} - \frac{x_\tau}{r_\tau}|^2 \\ &\geq \frac{|x_\tau|^2}{r_\tau^2} (\rho_j^2 - \rho_\tau^2)^2 \\ &\geq 4 \frac{\nu^2}{\rho_\tau^2} (\rho_j^2 - \rho_\tau^2)^2. \end{aligned}$$

Then $1 \ge 2 \frac{|\rho_j - \rho_\tau|}{\rho_\tau}$. Now, this inequality is equivalent to $|\rho_j - \rho_\tau| \le \frac{1}{2}\rho_\tau$ which in turn is equivalent to (4.2).

Proof of Lemma 1.2. We can write every $y \in \mathbb{R}^n$ as $y = (\xi + |x|\frac{1-s}{r})\frac{x}{|x|} + v$, with $\langle v, x \rangle = 0$. It is clear that $y \in B(\frac{x}{r}, \frac{|x|}{r}s)$ if and only if $\xi \in (0, 2s\frac{|x|}{r})$ and |v| < 0

License or copyright restrictions may apply to redistribution; see http://www.ams.org/journal-terms-of-use

$$\begin{split} \gamma \big(B\big(\frac{x}{r}, \frac{|x|}{r}s\big) \big) &= \int_{B(\frac{x}{r}, \frac{|x|}{r}s)} e^{-|y|^2} \, dz \\ &= e^{-\frac{|x|^2}{r^2}(1-s)^2} \int_0^{2s\frac{|x|}{r}} e^{-2\xi\frac{|x|}{r}(1-s)} e^{-|\xi|^2} \\ &\int_{\{v \in \mathbb{R}^{n-1}: \, |v| < \sqrt{2\frac{|x|}{r}s \, \xi - \xi^2}\}} e^{-|v|^2} \, dv \, d\xi \\ &\leq C_n e^{-\frac{|x|^2}{r^2}(1-s)^2} \int_0^{2s\frac{|x|}{r}} e^{-2\xi|x|\frac{1-s}{r}} \left(2\frac{|x|}{r}s \, \xi - \xi^2\right)^{\frac{n-1}{2}} \, d\xi \\ &\leq C_n e^{-\frac{|x|^2}{r^2}(1-s)^2} s^{\frac{n-1}{2}} \int_0^{2s\frac{|x|}{r}} e^{-2\xi|x|\frac{1-s}{r}} \left(2\frac{|x|}{r} \xi\right)^{\frac{n-1}{2}} \, d\xi \\ &\leq C_n e^{-\frac{|x|^2}{r^2}(1-s)^2} s^{\frac{n-1}{2}} \int_0^{2s\frac{|x|}{r}} e^{-2\xi|x|\frac{1-s}{r}} \left(2\frac{|x|}{r} \xi\right)^{\frac{n-1}{2}} \, d\xi \\ &\leq C_n e^{-\frac{|x|^2}{r^2}(1-s)^2} \frac{s^{\frac{n-1}{2}}}{|x|(1-s)^{\frac{n+1}{2}}} \int_0^{4s\frac{|x|^2}{r^2}(1-s)} e^{-t} t^{\frac{n-1}{2}} \, dt \\ &\leq C_n e^{-\frac{|x|^2}{r^2}(1-s)^2} \frac{s^{\frac{n-1}{2}}}{|x|} \min (1, s|x|^2) \\ &\leq C_n e^{-\frac{|x|^2}{r^2}(1-s)^2} s^{\frac{n-1}{2}} \min \left(\frac{1}{|x|}, s|x|\right) \\ &\leq C_n e^{-\frac{|x|^2}{r^2}(1-s)^2} s^{\frac{n-1}{2}} \min \left(\frac{1}{|x|}, s^{\frac{1}{2}}\right). \end{split}$$

Proof of Corollary 1.1. We choose $\Phi(t) = \frac{1}{\pi^{n/2}} e^{-t^2}$. From Lemma 1.2,

$$\gamma((1+\delta)B(\frac{x}{r},\frac{|x|}{r}(1-r)) \le C e^{-|x|^2}(1-r)^{\frac{n}{2}}$$

Then $T^*f \leq CM_{\Phi}f(x)$, and therefore the γ -weak type (1,1) inequality for T^* follows from Theorem 1.1.

References

- Calderón C. Some remarks on the multiple Weierstrass transform and Abel summability of multiple Fourier-Hermite series. Studia Mathematica, T. XXXII (1969), 119-148. MR0262770 (41:7375)
- Fabes E., Gutiérrez C., and Scotto R. Weak-type estimates for the Riesz transforms associated with the Gaussian measure. Rev. Mat. Iberoamericana 10 (1994), no. 2, 229–281. MR1286476 (95m:42023)
- [3] Forzani L. Lemas de cubrimiento de tipo Besicovitch y su aplicación al estudio del operador maximal de Ornstein-Uhlenbeck. (Besicovitch type covering lemmas and their applications to the study of the Ornstein-Uhlenbeck maximal operator). Ph.D. thesis disertation. Universidad Nacional de San Luis, Argentina.
- [4] Forzani L., Harboure E., and Scotto R. Weak type inequality for a family of singular integral operators related with the gaussian measure. Preprint.
- [5] Forzani L. and Scotto R. The higher order Riesz transform for Gaussian measure need not be of weak type (1,1). Studia Math. 131 (1998), no. 3, 205–214. MR1644460 (99h:42034)
- [6] Forzani L., Scotto R., and Urbina W. Riesz and Bessel potentials, the g^k functions and an area function for the Gaussian measure γ. Rev. Un. Mat. Argentina 42 (2000), no. 1, 17—37 (2001). MR1852728 (2002h:42022)

2153

- [7] Forzani L., Scotto R., and Urbina W. A simple proof of the L^p continuity of the higher order Riesz transforms with respect to the Gaussian measure γ^d. Sèminaire de Probabilitès, XXXV, 162–166, Lecture Notes in Math., 1755, Springer, Berlin, 2001. MR1837285 (2002c:42019)
- [8] García-Cuerva J., Mauceri G., Sjögren P., and Torrea J. Spectral multipliers for the Ornstein-Uhlenbeck semigroup. J. Anal. Math. 78 (1999), 281–305. MR1714425 (2000m:42015)
- [9] García-Cuerva J., Mauceri G., Sjögren P., and Torrea J. Higher-order Riesz operators for the Ornstein-Uhlenbeck semigroup. Potential Anal. 10 (1999), no. 4, 379–407. MR1698617 (2000j:42024)
- [10] Gundy R. Sur les transformations de Riesz pour le semi-groupe d'Ornstein-Uhlenbeck. C. R. Acad. Sci. Paris Sér. I Math. 303 (1986), no. 19, 967–970. MR0877182 (88c:60108)
- [11] Gutiérrez C. On the Riesz transforms for Gaussian measures. J. Funct. Anal. 120 (1994), no. 1, 107–134. MR1262249 (95c:35013)
- [12] Gutiérrez C., Segovia C., and Torrea J. On higher Riesz transforms for Gaussian measures. J. Fourier Anal. Appl. 2 (1996), no. 6, 583–596. MR1423529 (98m:42017)
- [13] Gutiérrez C. and Urbina W. Estimates for the maximal operator of the Ornstein-Uhlenbeck semigroup. Proc. Amer. Math. Soc. Proceedings of the American Mathematical Society 113 (1991), no. 1, 99-104. MR1068123 (92a:42023)
- [14] Hille E. A class of reciprocal functions. Ann. Math. 27 (1926), 426-464. MR1502747
- [15] Itô K. Malliavin's C[∞] Functionals of a Centered Gaussian System. IMA Preprint Series 327, (1987), University of Minnesota.
- [16] Messiah A. Quantum Mechanics. Vol. I. North-Holland, Publishing Co., Amsterdam. MR0129790 (23:B2826)
- [17] Meyer P.A. Transformations de Riesz pour les lois gaussiennes. Seminar on probability, XVIII, 179–193, Lecture Notes in Math., 1059, Springer, Berlin, 1984. MR0770960 (86i:60150)
- [18] Muckenhoupt B. Hermite conjugate expansions. Trans. Amer. Math. Soc. 139 (1969), 243–260. MR0249918 (40:3159)
- [19] Pérez S. Boundedness of Littlewood-Paley g-functions of higher order associated with the Ornstein-Uhlenbeck semigroup. Indiana Univ. Math. J. 50 (2001), no. 2, 1003–1014. MR1871397 (2002h:42026)
- [20] Pérez S. The local part and the strong type for operators related to the Gaussian measure.
 J. Geom. Anal. 11 (2001), no. 3, 491–507. MR1857854 (2002h:42027)
- [21] Pisier G. Riesz transforms: a simpler analytic proof of P.-A. Meyer's inequality. Seminaire de Probabilitès, XXII, 485–501, Lecture Notes in Math., 1321, Springer, Berlin, 1988. MR0960544 (89m:60178)
- [22] Sjögren P. On the maximal functions for the Mehler kernel. Springer-Lecture Notes in Math. 992 (1983), 73-82. MR0729346 (85j:35031)
- [23] Sjögren P. A remark on the maximal function for measures in Rⁿ. Amer. J. Math. 105 (1983), no. 5, 1231-1233. MR0714775 (86a:28003)
- [24] Sjögren P. Personal communication.
- [25] Stein E. Topics in harmonic analysis related to the Littlewood-Paley theory. Annals of Mathematics Studies, No. 63 Princeton University Press, Princeton, N.J., University of Tokyo Press, Tokyo 1970. MR0252961 (40:6176)
- [26] Urbina W. On singular integrals with respect to the Gaussian measure. Ann. Scuola Norm. Sup. Pisa Cl. Sci. (4) 17 (1990), no. 4, 531–567. MR1093708 (92d:42010)

Departamento de Matemáticas, Universidad Nacional del Litoral-CONICET, Santa Fe 3000, Argentina

E-mail address: haimar@ceride.gov.ar

Departamento de Matemáticas, Universidad Nacional del Litoral-CONICET, Santa Fe 3000, Argentina – and – School of Statistics, University of Minnesota, Minneapolis, Minnesota 55455

E-mail address: lforzani@math.unl.edu.ar

Departamento de Matemáticas, Universidad Nacional del Litoral, Santa Fe 3000, Argentina

E-mail address: scotto@math.unl.edu.ar