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Normal state electronic properties of $\text{LaO}_{1-x}\text{F}_x\text{BiS}_2$ superconductors

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ABSTRACT

A good description of the electronic structure of BiS_2 -based superconductors is essential to understand their phase diagram, normal state and superconducting properties. To describe the first reports of normal state electronic structure features from angle resolved photoemission spectroscopy (ARPES) in $\text{LaO}_{1-x}\text{F}_x\text{BiS}_2$, we used a minimal microscopic model to study their low energy properties. It includes the two effective tight-binding bands proposed by Usui et al., Phys. Rev. B, 86, 2012, 220501(R), and we added moderate intra- and inter-orbital electron correlations related to $\text{Bi}(p_y, p_x)$ and $\text{S}(p_y, p_x)$ orbitals. We calculated the electron Green's functions using their equations of motion, which we decoupled in second-order of perturbations on the correlations. We determined the normal state spectral density function and total density of states for $\text{LaO}_{1-x}\text{F}_x\text{BiS}_2$, focusing on the description of the k -dependence, effect of doping, and the prediction of the temperature dependence of spectral properties. Including moderate electron correlations, improves the description of the few experimental ARPES and soft X-ray photoemission data available for $\text{LaO}_{1-x}\text{F}_x\text{BiS}_2$. Our analytical approximation enabled us to calculate the spectral density around the conduction band minimum at $k_0 = (0.45\pi, 0.45\pi)$, and to predict the temperature dependence of the spectral properties at different BZ points, which might be verified by temperature-dependent ARPES.

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1. Introduction

In 2012, Mizuguchi et al. [1] discovered novel layered superconductors which have a crystal structure similar to that of the cuprate and iron-based superconductors, the so-called BiS_2 -based superconductors. Experimental and theoretical studies have been carried out in order to establish the basic properties of these new materials and identify the underlying mechanism for superconductivity [2,4,3]. Substantial enhancement of superconductivity under moderate pressures in BiS_2 -based compounds has been reported [5]. Up to date, eleven compounds have been discovered in the BiS_2 family [4,3], and the highest T_c is 10.6 K, reported in $\text{LaO}_{0.5}\text{F}_{0.5}\text{BiS}_2$ under 2 GPa of applied pressure [1]. The crystal lattice of BiS_2 -based superconductors consists of consecutive superconducting BiS_2 planes and blocking layers. The conduction planes can be viewed as a square array of Bi atoms, each of them with a basis of two S atoms attached [1,4,3]. The

same crystal structure was found in $\text{LaO}_{0.5}\text{F}_{0.5}\text{BiSe}_2$ [6], in this case alternating superconducting BiSe_2 and blocking LaO layers, with a smaller $T_c = 2.6$ K.

For the discovered superconducting materials within the BiS_2 family, superconductivity emerges from a metallic normal state [7,8] while the parent compounds are band insulators. In fact, $\text{LaO}_{1-x}\text{F}_x\text{BiS}_2$ changes from a band insulator at $x=0$ to a metal by means of chemical substitution: upon electron doping [1,7,8]. In contrast to iron-based superconductors [9], in the LnOBiS_2 ($\text{Ln} = \text{La, Ce, Pr, Nd}$) parent compounds neither magnetic nor structural transitions have yet been reported, indicating that magnetism would be of less relevance to superconductivity in the BiS_2 -based compounds as in iron-based superconductors [10]. Recently, the coexistence of superconductivity and ferromagnetism has been reported in $\text{Ce}_{1-x}\text{F}_x\text{BiS}_2$ compounds with $x > 0.4$, while at $x < 0.4$ a paramagnetic behavior was observed [11]. Very recently, an analysis comparing growth and characterization of LnOBiS_2 single crystals [12] highlights the fact that sizeable differences may exist between the nominal F-composition (x) and the analytical F-composition (y), which is obtained characterizing the samples by X-ray diffraction, scanning electron microscopy and electron

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probe microanalysis. For instance, for $\text{LaO}_{1-x}\text{F}_x\text{BiS}_2$ single crystals, to nominal: $x=0.3, 0.4, 0.5, 0.6$, respectively correspond analytical values: $y=0.23, 0.37, 0.43, 0.45$ [12].

A number of density functional theory (DFT) band-structure calculations have been reported [14,13,22,15–18,21,19,20], in particular for $\text{LaO}_{0.5}\text{F}_{0.5}\text{BiS}_2$ and its parent compound. The valence bands extend from -6 to 0 eV and consist of p-electron states from the O and S atoms, while the bands related to La-*f*, La-*d* lie far above the Fermi level (E_F), near 4 eV, and the conduction bands are dominated by Bi-*p* and S-*p* electron states [7]. Based on DFT calculations, a minimal two-orbital model for $\text{LaO}_{1-x}\text{F}_x\text{BiS}_2$ was proposed by Usui et al. [13] predicting a change of topology of the Fermi surface (FS) when electron doping x is increased. Recently, the FS topology proposed for higher doping [13] was confirmed by ARPES in $\text{LaO}_{0.54}\text{F}_{0.46}\text{BiS}_2$ [23]. Usui et al. [13] also suggested the presence of nesting of the FS with wave vector $\vec{k} = (\pi, \pi, 0)$ at $x \sim 0.5$, due to the quasi-one-dimensional nature of the conduction bands.

Core-level and valence-band soft X-ray photoemission spectroscopy (SXPS) [24] were used to investigate the electronic structure of $\text{LaO}_{1-x}\text{F}_x\text{BiS}_2$ ($x = 0, 0.3, 0.5$), which was found to be mainly consistent with predictions of DFT calculations, including the doping dependence [25,14,7]. Nevertheless, a deviation was found between the experimental intensity and spectral shape of the states near E_F in comparison with those calculated [25,14,7]. As x increases from 0 to 0.5 , a rigid-band shift of the whole band-structure is predicted, based on the valence band shift observed: towards higher binding energy by about ~ 0.3 eV, a value which is much smaller than expected from DFT calculations [25]. Recently, the electronic structure of nearly optimally doped superconductor $\text{LaO}_{0.54}\text{F}_{0.46}\text{BiS}_2$ was investigated using ARPES [23], finding relatively weak electron correlation effects and a marked influence of spin-orbit coupling on the BiS_2 planes. A square-like ARPES intensity distribution centered at the BZ center was found, consistent with the prediction of the effective two-band model by Usui et al. [13], who indicated that the Fermi surface was close to a topological transition.

Other theoretical studies have focused on the study of electron-electron correlations and the possibility of unconventional superconductivity in BiS_2 -based compounds [22,26,27,21,28]. In this context, superconductivity was investigated by Zhou et al. [22] using a microscopic model for LaOBiS_2 including the effective two-bands proposed by Usui et al. [13] and an on-site intra-orbital Coulomb repulsion U . Concretely, the spin excitations in the superconducting state were studied calculating the spin susceptibility in random phase approximation (RPA), and three potential pairing symmetries (*d*-, *s*- and *p*-wave) were analyzed, in a weak correlation scenario: $U \sim 0.8$ – 2.5 eV. With the further addition of an on-site interorbital Coulomb repulsion as well as Hund coupling to the two-orbital model, Martins et al. [26] investigated the superconductive pairing properties of BiS_2 -based superconductors. Employing a weak-coupling multiorbital RPA analysis, a clear relationship between quasi nesting in the Fermi surface, spin fluctuations, and superconductivity was found [26], predicting that pairing symmetry measurements may contain a mixture of both A_{1g} and B_{2g} symmetries, suggesting $U \sim 1.08$ – 1.8 eV. In Ref. [29] the pairing symmetry of BiS_2 compounds was investigated using the two-band model in Ref. [13] by assuming that electron-electron correlation are relevant. It was also assumed [29] that short range superconductive pairing stems from short antiferromagnetic exchange couplings, finding that extended *s*-wave pairing symmetry is always more favourable than *d*-wave pairing. The normal state spectral properties using the two-band minimal model of Ref. [13], were not investigated in RPA yet.

Many aspects of the normal state of $\text{LaO}_{1-x}\text{F}_x\text{BiS}_2$, such as the

effect of electron correlations on the electronic structure, the doping and temperature dependence of the spectral properties, remain unstudied, which has prompted our investigation. Here, we present our study of the paramagnetic normal state of $\text{LaO}_{1-x}\text{F}_x\text{BiS}_2$ compounds. In particular, we analyze the effect of electron correlations, believed to be moderate in these compounds, as well as of doping and of temperature on the spectral properties. We do this employing a minimal microscopic model, which includes the two effective uncorrelated orbitals proposed by H. Usui et al. [13] for $\text{LaO}_{1-x}\text{F}_x\text{BiS}_2$ as well as intra- and inter-orbital Coulomb interactions, and using the method of equations of motion for the Green's functions which we solved in second-order of perturbations on the correlations, to determine the temperature, doping and *k*-dependent spectral density function and density of states for $\text{LaO}_{1-x}\text{F}_x\text{BiS}_2$. In a previous work [30], we used this analytical approach to study the normal state spectral properties of ferropnictides [30], using for the non-interacting bandstructure the two effective tight-binding orbitals proposed by Raghu et al. [31]. Our paper is organized as follows: in Section 2 we describe the microscopic correlated two-orbital model adopted for our study of the normal state spectral properties of $\text{LaO}_{1-x}\text{F}_x\text{BiS}_2$ compounds, and the analytical method employed (which is complemented by Appendix A, where we state the detailed expressions obtained for the relevant Green's functions). In Section 3 we present and discuss our results at different temperature and doping values, for relevant BZ points, and compare them with the relatively few available theoretical and experimental results for the normal state of $\text{LaO}_{1-x}\text{F}_x\text{BiS}_2$. In Section 4 we summarize our main conclusions, prompting for experiments to test our predictions, and mention possible future applications of our work.

2. Microscopic model and analytical approach

2.1. Correlated effective model

As mentioned in Introduction, an effective non-interacting two-orbital tight-binding model to describe the low energy physics of $\text{LaO}_{1-x}\text{F}_x\text{BiS}_2$, was proposed by Usui et al. in Ref. [13]. Starting from a first-principles band calculation for LaOBiS_2 , where they found that the conduction bands near the Fermi level consist mainly of a mixture of in-plane Bi-6*p* and S-3*p* orbitals, they constructed maximally localized Wannier orbitals to obtain effective tight-binding models, describing the kinetic energy part of the effective Hamiltonian. Extracting the portion of the bands which is relevant to the BiS_2 layers, assumed to be related to superconductivity, and neglecting interlayer hoppings, when focusing only on the bands that intersect the Fermi level Usui et al. [13] obtained a reduced Hamiltonian consisting of a two-dimensional effective two-orbital tight-binding model. The model [13] yields two bands (per BiS_2 layer) which have a quasi-one-dimensional character, providing good nesting of the Fermi surface, with a dispersion characterized by two conduction band minima [13] located at $X = (\pi, 0)$ and $k_0 = (0.45, 0.45)\pi$. DFT calculations [13] indicated that the FS topology should undergo considerable changes with doping as discussed in previous section, which the two-orbital model could describe. Usui et al [13] also pointed out that the total electron filling in the two-orbital model, n , should be taken as: $n=x$ to describe $\text{LaO}_{1-x}\text{F}_x\text{BiS}_2$.

Concretely, the effective two-orbital tight-binding model of Usui et al. [13] to describe the kinetic energy part of the effective Hamiltonian for $\text{LaO}_{1-x}\text{F}_x\text{BiS}_2$, is given by:

$$\mathcal{H}_0 = \sum_{k,\sigma} \left[E_c(\vec{k}) c_{k\sigma}^\dagger c_{k\sigma} + E_d(\vec{k}) d_{k\sigma}^\dagger d_{k\sigma} \right] \quad (1)$$

where the operators $c_{\vec{k},\sigma}^\dagger$ and $d_{\vec{k},\sigma}^\dagger$ create an electron in the respective orbitals c and d , with spin $\sigma = \uparrow, \downarrow$, and crystal momentum \vec{k} , while the respective energies are:

$$E_c(\vec{k}) = \frac{1}{2} \left[\epsilon_X + \epsilon_Y \pm \sqrt{(\epsilon_X - \epsilon_Y)^2 + \epsilon_{XY}^2} \right] - \mu \quad (2)$$

where μ is the chemical potential, and:

$$\begin{aligned} \epsilon_X(\mathbf{k}) = & t_0 + 2t_1(\cos k_x + \cos k_y) + 2t_3 \cos(k_x - k_y) + 2t_4 \cos(k_x + k_y) \\ & + 2t_6[\cos(2k_x + k_y) + \cos(k_x + 2k_y)] \\ & + 2t_8[\cos(2k_x - k_y) + \cos(k_x - 2k_y)] \end{aligned}$$

$$\begin{aligned} \epsilon_Y(\mathbf{k}) = & t_0 + 2t_1(\cos k_x + \cos k_y) + 2t_3 \cos(k_x + k_y) + 2t_4 \cos(k_x - k_y) \\ & + 2t_6[\cos(2k_x - k_y) + \cos(k_x - 2k_y)] \\ & + 2t_8[\cos(2k_x + k_y) + \cos(k_x + 2k_y)] \end{aligned}$$

$$\begin{aligned} \epsilon_{XY}(\mathbf{k}) = & 2t_2(\cos k_x - \cos k_y) + 2t_5(\cos 2k_x - \cos 2k_y) \\ & + 4t_7[\cos(2k_x + k_y) - \cos(k_x + 2k_y)] \end{aligned} \quad (3)$$

The hopping parameters up to fourth neighbors in Eqs. (3), between sites on the square lattice formed by the Bi atoms, are taken from Refs. [13,26]: $t_0 = 2.811$, $t_1 = -0.167$, $t_2 = 0.107$, $t_3 = 0.880$, $t_4 = 0.094$, $t_5 = -0.028$, $t_6 = 0.014$, $t_7 = 0.020$, $t_8 = 0.069$, where all energy parameters are given in eV (as in the rest of our paper).

In our present work, to describe analytically the normal state properties of $\text{LaO}_{1-x}\text{F}_x\text{BiS}_2$ compounds, we will consider a minimal model preserving the essential physics of the problem given by:

$$\mathcal{H} = \mathcal{H}_0 + V_{int} \quad (4)$$

where the kinetic energy is described by \mathcal{H}_0 of Eq. (1) including the two effective orbitals proposed by Usui et al. [13] detailed above, while we add electron correlations in V_{int} . In fact, assuming that short-range Hubbard-like electron–electron interactions are present, at each site we will consider an intra-orbital Coulomb repulsion U , and an inter-orbital repulsion V , so that V_{int} reads:

$$V_{int} = \sum_i [U(n_{i\uparrow}n_{i\downarrow} + N_{i\uparrow}N_{i\downarrow}) + V(n_{i\uparrow} + n_{i\downarrow})(N_{i\uparrow} + N_{i\downarrow})] \quad (5)$$

where $n_{i\sigma} = c_{i\sigma}^\dagger c_{i\sigma}$ and $N_{i\sigma} = d_{i\sigma}^\dagger d_{i\sigma}$, for spins $\sigma = \uparrow, \downarrow$.

Thus, we will be modelling $\text{LaO}_{1-x}\text{F}_x\text{BiS}_2$ by an effective extended Hubbard model consisting of two correlated electron orbitals per site, to study the electronic properties of the paramagnetic normal state of these compounds: in particular, the k -dependence, effect of doping and temperature dependence of the spectral properties.

2.2. Analytical calculation of the spectral properties

To determine the spectral density of c and d electrons in our effective extended Hubbard model, we calculated the corresponding retarded Green's functions at finite temperature introduced by Zubarev [32] defined by:

$$G_\sigma^{ret}(\vec{k}, \omega) = G(k, \omega + i\delta) = \langle\langle c_{k\sigma}; c_{k\sigma}^\dagger \rangle\rangle(\omega + i\delta) \quad (6)$$

$$F_\sigma^{ret}(\vec{k}, \omega) = F(k, \omega + i\delta) = \langle\langle d_{k\sigma}; d_{k\sigma}^\dagger \rangle\rangle(\omega + i\delta) \quad (7)$$

where δ is an infinitesimal positive number, and $\langle\langle \dots \rangle\rangle$ represents the usual notation for Zubarev's Green's functions [32] which for fermionic operators \hat{A}, \hat{B} are defined as the time Fourier-

transform of the retarded Green's function:

$$\langle\langle \hat{A}(t); \hat{B}(t') \rangle\rangle = -i\theta(t-t')\langle \hat{A}(t)\hat{B}(t') - \hat{B}(t')\hat{A}(t) \rangle \quad (8)$$

where the time-dependent operators appear in Heisenberg representation, and $\theta(t)$ is the Heaviside step function. The expectation values $\langle \dots \rangle$ of quantum observables at finite temperature are calculated as the trace of the product of the density operator and the observable, with the density operator written in terms of the appropriate statistical ensemble at temperature T [32]. In our case, we study the paramagnetic normal state of the system at temperature T with fixed number of electrons n , thus evaluating the required expectation values using the grand canonical ensemble. (The technique can also be extended to study zero-temperature properties, in which case the Green's function definition involves calculating the expectation values in the ground state of the system.)

We calculated $G_\sigma(k, \omega)$ and $F_\sigma(k, \omega)$ using Zubarev's equations of motion (EOM) formalism, [32,33] in particular from the following exact coupled set of equations:

$$\begin{aligned} [\omega - E_c(k)]G_\sigma(k, \omega) = & \frac{1}{2\pi} + \sum_{k_1, k_2} \left[\frac{U}{N} \Gamma_1(k_1, k_2, k, \omega) \right. \\ & \left. + \frac{V}{N} \Gamma_2(k_1, k_2, k, \omega) + \frac{V}{N} \Gamma_3(k_1, k_2, k, \omega) \right] \end{aligned} \quad (9)$$

$$\begin{aligned} [\omega - E_d(k)]F_\sigma(k, \omega) = & \frac{1}{2\pi} + \sum_{k_1, k_2} \left[\frac{U}{N} \Gamma_4(k_1, k_2, k, \omega) \right. \\ & \left. + \frac{V}{N} \Gamma_5(k_1, k_2, k, \omega) + \frac{V}{N} \Gamma_6(k_1, k_2, k, \omega) \right] \end{aligned} \quad (10)$$

where in Eq. (9) we denoted the higher-order Green's functions which encode the charge, spin and orbital fluctuations:

$$\begin{aligned} \Gamma_1(k_1, k_2, k, \omega) & \equiv \Gamma_{\sigma, \bar{\sigma}, \bar{\sigma}}^{ccc}(k_1, k_2, k, \omega) \\ \Gamma_2(k_1, k_2, k, \omega) & \equiv \Gamma_{\sigma, \sigma, \sigma}^{cdd}(k_1, k_2, k, \omega) \\ \Gamma_3(k_1, k_2, k, \omega) & \equiv \Gamma_{\sigma, \bar{\sigma}, \bar{\sigma}}^{cdd}(k_1, k_2, k, \omega) \end{aligned} \quad (11)$$

and in Eq. (10):

$$\begin{aligned} \Gamma_4(k_1, k_2, k, \omega) & \equiv \Gamma_{\sigma, \bar{\sigma}, \bar{\sigma}}^{ddd}(k_1, k_2, k, \omega) \\ \Gamma_5(k_1, k_2, k, \omega) & \equiv \Gamma_{\sigma, \sigma, \sigma}^{dcc}(k_1, k_2, k, \omega) \\ \Gamma_6(k_1, k_2, k, \omega) & \equiv \Gamma_{\sigma, \bar{\sigma}, \bar{\sigma}}^{dcc}(k_1, k_2, k, \omega) \end{aligned} \quad (12)$$

while the following notation was introduced above:

$$\Gamma_{\sigma\alpha\sigma\beta\sigma\gamma}^{\alpha,\beta,\gamma}(k_1, k_2, k, \omega) \equiv \langle\langle c_{\alpha k_2, \sigma\alpha} c_{\beta k_1, \sigma\beta}^\dagger c_{\gamma k_1-k_2+k, \sigma\gamma}; c_{k\sigma}^\dagger \rangle\rangle(\omega) \quad (13)$$

The Hartree–Fock solution to this problem is obtained by closing the set of equations of motion in first-order of perturbations on the correlations, performing a mean-field approximation of the Γ_i ($i = 1, 6$) Green's functions in Eqs. (9) and (10), by which each is expressed only in terms of $G_\sigma(k, \omega)$ and $F_\sigma(k, \omega)$. In Appendix A.1 we state the Hartree–Fock solutions for the c and d electron Green's functions.

In the present work, instead, we solved the problem at a higher level of approximation: we went on to calculate the three exact equations of motion for the Γ_i ($i = 1, 3$) Green's functions, coupled to $G(\vec{k}, \omega)$ in first order through Eq. (9); and proceeded likewise for Γ_i ($i = 4, 6$) coupled to $F(\vec{k}, \omega)$ through Eq. (10). Each of the six second-order equations of motion for Γ_i ($i = 1, 6$) introduces three new coupled higher-order Green functions in the problem.

To close the coupled set of equations of motion in second-order

of perturbations on the correlations, we used the approximation previously employed in our investigation of ferropnictides in Ref. [30] (for full details of our calculations, see Appendix A of Ref. [30]). Briefly, our approach consists in approximating all the higher-order Green's functions introduced in each subset of equations of motion for Γ_i ($i = 1, 3$), expressing them in mean-field only in terms of Γ_i ($i = 1, 3$) and $G(\vec{k}, \omega)$. We proceeded likewise for the subset of equations for Γ_i ($i = 4, 6$) related to $F(\vec{k}, \omega)$. This second-order approximation yields a closed system of eqs. of motion for $G, \Gamma_1, \Gamma_2, \Gamma_3$; and an analogous one for $F, \Gamma_4, \Gamma_5, \Gamma_6$, which we solved to determine $G(\vec{k}, \omega)$ and $F(\vec{k}, \omega)$ obtaining the expressions presented in Appendix A.2. Notice that the evaluation of c and d Green's functions obtained in our second-order perturbative approach requires performing double and triple summations over the first Brillouin zone (BZ) of the crystal lattice. For simplicity, we have assumed a square lattice, as done previously, [13] and for the BZ summations used the Chadi-Cohen BZ sampling method [35]. As well known for other Hubbard-like models (see e.g. Refs. [34,33]), the decoupling scheme adopted for the higher order Green's functions coupled in the second-order EOM, accounts for non-trivial k -dependent self-energy effects of the single-particle Green function, as our results in next section will show (and we previously demonstrated for ferropnictides: in particular, in Appendix B of Ref. [30]).

With the c and d electron Green's functions determined for the paramagnetic normal state of our effective extended Hubbard model, we calculate the total spectral density function, to be compared with ARPES experiments, $A(\vec{k}, \omega)$, as follows:

$$A(\vec{k}, \omega) = \sum_{\sigma} A_{c\sigma}(\vec{k}, \omega) + A_{d\sigma}(\vec{k}, \omega) \quad (14)$$

where

$$A_{c\sigma}(\vec{k}, \omega) = -\frac{1}{\pi} \text{Im} C_{\sigma}^{\text{ret}}(k, \omega); \quad A_{d\sigma}(\vec{k}, \omega) = -\frac{1}{\pi} \text{Im} F_{\sigma}^{\text{ret}}(k, \omega) \quad (15)$$

The total density of states (DOS), denoted as $A(\omega)$, is obtained by integrating over the whole BZ the spectral density function:

$$A(\omega) = \sum_{k \in \text{BZ}} A(k, \omega) \quad (16)$$

Finally, the self-consistent set of coupled equations to be solved in our second-order perturbative approach is completed by the equation which yields the total electron band filling n (or chemical potential μ) at temperature T :

$$\begin{aligned} n(\mu) &= x = n_c + n_d \\ &= \int \frac{d\vec{k}}{(2\pi)^2} \frac{d\omega}{2\pi} \left[A_c(\vec{k}, \omega) \frac{1}{e^{\beta E_c(\vec{k})} + 1} + A_d(\vec{k}, \omega) \frac{1}{e^{\beta E_d(\vec{k})} + 1} \right] \end{aligned} \quad (17)$$

where $\beta = 1/k_B T$ and k_B is Boltzmann's constant.

3. Results and discussion

In this section we present the electronic structure results obtained in our approach for the paramagnetic normal state of $\text{LaO}_{1-x}\text{F}_x\text{BiS}_2$ superconductors, at different temperature and doping values, using the model presented in Section 2.1. We compare our results with the relatively few experimental ARPES and soft X-ray photoemission data and previous theoretical results which are available for these compounds. We solved the self-consistent equations for the Green's functions detailed in Section 2.2 numerically, using the tight-binding parameters of Ref. [13] for the non-interacting effective bands proposed to describe $\text{LaO}_{1-x}\text{F}_x\text{BiS}_2$.

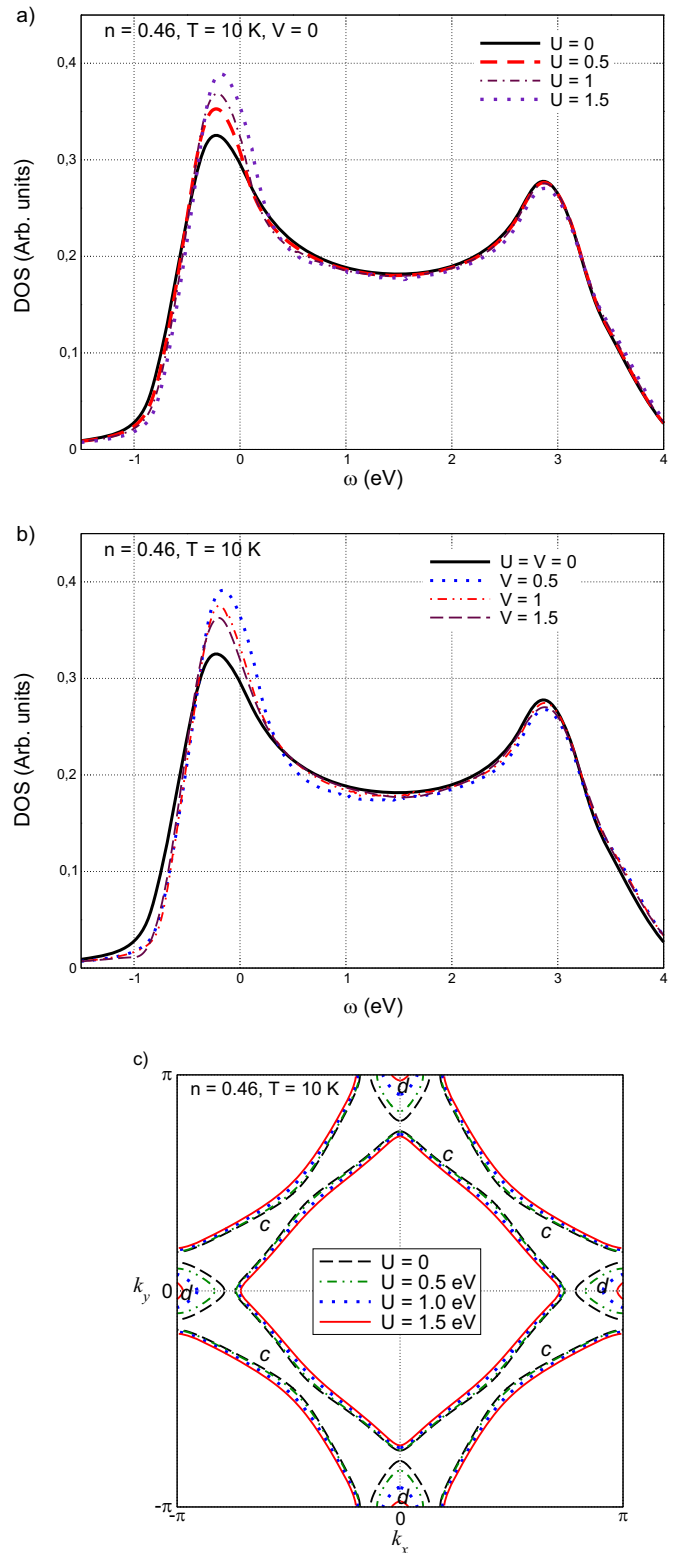


Fig. 1. (a) Total density of states as a function of energy, $A(\omega)$ (ω measured w.r.t the Fermi level, E_F), for different values of U (as indicated in the plot), and at fixed $V=0$; (b) $A(\omega)$ for different values of V (as indicated in the plot), at fixed $U=1.5$ eV (except for the solid line: depicting the uncorrelated case); (c) Effect of correlations on the FS: U as indicated in plot and: $V=U/2$, c and d labels indicate the orbital character of each set of Fermi surface sheets. In this figure: electron band filling $n = x = 0.46$ and temperature: $T = 10$ K. Non-interacting tight-binding parameters from Ref. [13]: $t_0 = 2.811$, $t_1 = -0.167$, $t_2 = 0.107$, $t_3 = 0.880$, $t_4 = 0.094$, $t_5 = -0.028$, $t_6 = 0.014$, $t_7 = 0.020$, $t_8 = 0.069$, and in the units of eV. $k_0 = (0.45, 0.45)\pi$ (in eV). Chadi-Cohen BZ summations order: $\nu=8$.

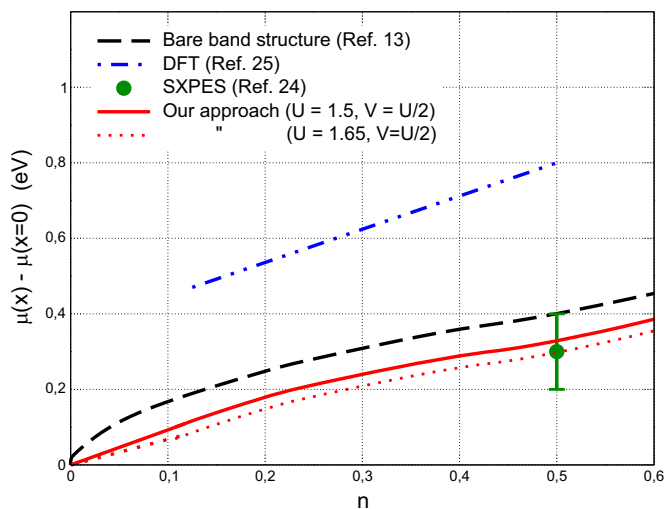


Fig. 2. Dependence of the chemical potential shift on total electron band filling ($n=x$): comparison between soft X-ray photoemission data at $T=300$ K in $\text{LaO}_{1-x}\text{F}_x\text{BiS}_2$ [24], DFT predictions, [25] the non-interacting two-orbital model [13] and the results including correlations with our second-order perturbative approach at $T=300$ K.

Regarding the precision of the Chadi–Cohen sampling method [35] used for the evaluation of Brillouin zone summations in the Green's functions, it is specified by the order ν : we found good convergence in most results using the seventh order ($\nu=7$) for the square-lattice BZ Chadi–Cohen sampling, [35] but in this paper all spectral density function and total density of states results presented were obtained using ninth order for improved accuracy, except for Figs. 1 and 2 where $\nu=8$ was used (as indicated in the respective figure captions). Notice that $\nu=8$ order for the square-lattice BZ Chadi–Cohen sampling implies using 8256 special BZ points, while $\nu=9$ implies using 32,896 ones [35]. For the retarded Green's functions $\delta=10^{-6}$ was used in the present work, except for the DOS results plotted: with $\delta=10^{-4}$ (these are typical values, for instance in Ref. [26] $\delta=10^{-5}$ was used).

We present our results divided in subsections, each of them focusing on different aspects of the problem. First, we discuss the separate effect of the intra- and inter-orbital electron correlations in our approach for $\text{LaO}_{1-x}\text{F}_x\text{BiS}_2$, in particular on the density of states. We also analyze the renormalization by correlations of the Fermi surface and effective bands. Comparison with chemical potential shift data for specific electron doping, allow us to identify the range of moderate correlation values to be used in the rest of our paper for a proper description of the compounds. Next, in Section 3.2 we focus on the momentum dependence of the spectral density function along symmetry paths of the square lattice Brillouin zone, comparing our results with the available reported ARPES data at X, and predicting the spectral density to be expected near $k_0=(0.45, 0.45)\pi$ (the second relevant conduction band minimum for $\text{LaO}_{1-x}\text{F}_x\text{BiS}_2$). In Section 3.3 we discuss the doping dependence of the total density of states. In Section 3.4 we predict the temperature dependence of the total density of states and of the spectral density function at different k-points for $\text{LaO}_{1-x}\text{F}_x\text{BiS}_2$, and compare with available data on $\text{NdO}_{1-x}\text{F}_x\text{BiS}_2$.

3.1. Effect of intra- and inter-orbital Coulomb interactions on the electronic structure

Though several scenarios have been proposed to describe the superconductivity in BiS_2 -based compounds, as discussed in Section 1, there is still no agreement on the relevant mechanism originating it. Similarities with the cuprate and pnictide

superconductors indicate the possibility of an unconventional pairing mechanism involving electron correlations [26], possibly not strong [13].

In this subsection we analyze the effect on the electronic structure of the intra- and inter orbital correlations included in the model for $\text{LaO}_{1-x}\text{F}_x\text{BiS}_2$ presented in Section 2.1. We compare the results obtained with our second order perturbative approach for the electron Green's functions, detailed in Section 2.2 and Appendix A, with other theoretical and the relatively few experimental data yet available, determining typical values for U and V to be used in the rest of the paper.

We first show the separate effects of the intra- and inter-orbital correlation parameters included in the model. The total density of states as a function of energy is shown in Fig. 1(a), for different values of intra-orbital correlation U between 0 and 1.5 eV, at fixed inter-orbital: $V=0$, and at the same temperature $T=10$ K and electron filling $n=0.46$ as in the only normal state ARPES experiment yet available for the $\text{LaO}_{1-x}\text{F}_x\text{BiS}_2$ family [23]. At those temperature and filling values, for $U=1.5$ eV $=2V$ the chemical potential is: $\mu=1.11$ eV. Being the total bandwidth $W\sim 4.5$ eV for the non-interacting (bare) effective band structure of Ref. [13], $U=1.5$ eV would correspond approximately to $0.33W$. Increasing U , we predict that the main difference in the DOS curves (see Fig. 1 (a)) would be a monotonous increase of the spectral weight of the peak located slightly below E_F w.r. to the non-interacting case ($U=0=V$).

In Fig. 1(b), we focalize on the effect of the inter-orbital Coulomb interaction V on the density of states, when a fixed $U=1.5$ eV is considered (the $U=0=V$ case is included for comparison). We find that the effect of increasing V is opposite to that of increasing U , effectively tending to screen it: the main effect of increasing V is a monotonous decrease of the spectral weight of the DOS peak located slightly below the Fermi level, evolving the resulting DOS towards the non-interacting one ($U=0=V$). Therefore, due to the differences in the relevant non-interacting effective bandstructure of $\text{LaO}_{1-x}\text{F}_x\text{BiS}_2$ [13] and ferropnictides [31], and the corresponding differences in location of the Fermi levels, we predict a striking difference in the effect of the inter-orbital electron correlation V which is relevant in BiS_2 -based superconductors, while largely irrelevant for ferropnictides ([30,37,36]). This is a fact already pointed out in connection with other correlated systems, [38,39], where the presence of two or more bands close in energy to E_F , with relatively low density of electrons (as would be the case for $\text{LaO}_{1-x}\text{F}_x\text{BiS}_2$ [13]), would amplify the effect of inter-orbital electron correlations.

The non-interacting Fermi surface of $\text{LaO}_{1-x}\text{F}_x\text{BiS}_2$ has been discussed in Ref. [13,26], it includes a topological FS change between one topology characteristic of low doping, experimentally confirmed by ARPES in related compounds: $\text{NdO}_{1-x}\text{F}_x\text{BiS}_2$ [41,40], and a FS topology characteristic of high doping, experimentally confirmed for $\text{LaO}_{0.54}\text{F}_{0.46}\text{BiS}_2$ by ARPES [23]. Using the non-interacting effective two-orbital model by Usui et al. [13], the FS change of topology would take place at the critical doping: $n_{crit}=0.44$, while we checked that this value is reduced when correlations are taken into account: for $U=1.5$ eV $=2V$ with our second-order perturbative approach we find $n_{crit}=0.35$, while in Hartree–Fock approximation $n_{crit}=0.22$. In Fig. 1(c) we exhibit the effect of the renormalization by correlations on the FS obtained in our approach for $\text{LaO}_{0.54}\text{F}_{0.46}\text{BiS}_2$ and compare it with the non-interacting result [13], at temperature $T=10$ K for the nominal composition $x=n=0.46$, exhibiting a FS with the characteristic high doping topology. We predict that the main effect of increasing electron correlations would be to enhance the nesting properties of large portions of the FS (in particular, of the large central sheet around Γ , related to c orbitals). Interestingly, a similar effect on nesting was reported by Martins et al. [26] but as a result from

an increase of doping, in the same paper where they established a relationship between the quasi-nesting properties of the FS, spin fluctuations and superconductivity in BiS_2 superconductors. Fig. 1 (c) confirms that, at fixed electron filling, the size of the FS is conserved upon the increase of electron correlations, as expected: we find that the decrease of the area of the electron pockets related to the d -band, is compensated by the increase of the FS area related to the c electrons.

In order to determine typical values for U and V for an adequate description of $\text{LaO}_{1-x}\text{F}_x\text{BiS}_2$, in Fig. 2 we compare different results for the chemical potential shift as a function of doping. We include the experimental SXPES estimation of Ref. [24] who, assuming that the Fermi level of the $x=0$ sample is located at the bottom of the conduction band, deduced a chemical potential shift of ~ 0.3 eV between the $x=0$ and 0.5 samples. This value is much smaller than the expected one from band calculations in Ref. [25] (0.8 eV), also plotted. Notice in Fig. 2, also, that the non-interacting effective two-orbital model [13] yields a chemical potential shift of ~ 0.4 eV between $x=0$ and 0.5, and that adding correlations with our perturbative approach, the description of the experimental results can be improved. Therefore, in the following we will present our results for the description of $\text{LaO}_{1-x}\text{F}_x\text{BiS}_2$ fixing $U=1.5$ eV and $V=0.75$ eV, similar values to those used in Ref. [26] to study other properties.

Finally, in Fig. 3 we compare the renormalized band structure for $\text{LaO}_{1-x}\text{F}_x\text{BiS}_2$ obtained including correlations in our second-order perturbative approach (detailed in Section 2.2 and Appendix A) with the first-order Hartree-Fock results, and the effective bands of the non-interacting two-orbital model proposed in Ref. [13]. Notice that, by including correlations, at Γ and M the degeneracy of the non-interacting band structure is broken, as a result of the difference in the c - and d -band self-energies obtained in both the first order (Hartree-Fock) and the second order perturbative approaches whose results are detailed in Appendix A, i.e. in both approximations an orbital-dependent renormalization by correlations appears, which gives rise to the lifting of the degeneracy at Γ and M . While in Hartree-Fock approximation, the rigid-band (k -independent) renormalization yields a larger total bandwidth, the renormalization obtained with our second-order approximation is k -dependent and yields a narrowing of the total band structure w.r. to the non-interacting two-orbital model [13].

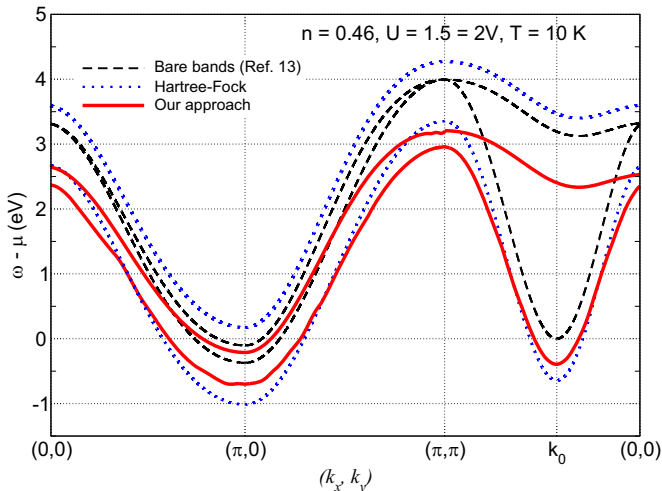


Fig. 3. Band structure renormalization: comparison of the renormalization by correlations ($U = 1.5$ eV = $2V$). Solid lines: our second-order perturbative approach; dotted lines: Hartree-Fock with the non-interacting (dashed lines [13]) band structure of the two-orbital model for $\text{LaO}_{1-x}\text{F}_x\text{BiS}_2$. In each case, the upper band corresponds to the d -orbital, and the lower band to the c -orbital. $k_0 = (0.45, 0.45)\pi$. Other parameters as in Fig. 1, except for: $\nu=9$.

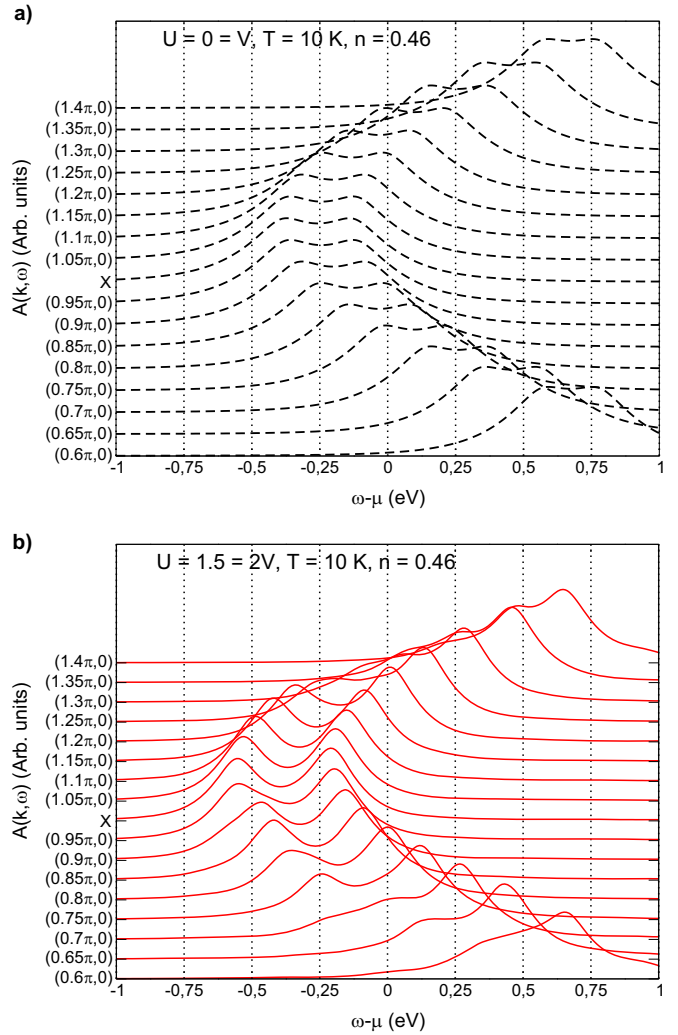


Fig. 4. Spectral density $A(\vec{k}, \omega)$ as a function of energy ω , for different BZ points (shown vertically displaced along the symmetry path $X - \Gamma$): (a) for the non-interacting two-orbital model [13], (b) present work: correlated two-orbital model, with our second-order perturbative approach. Parameters: $U = 1.5$ eV = $2V$, Other parameters as in Fig. 1, except for: $\nu=9$.

Notice in Fig. 3, also, that while the non-interacting model and our second-order perturbative approach to include correlations agree near $\vec{k} = (\pi, 0)$, in placing electrons in the c and d bands which cross the Fermi level, in Hartree-Fock only c -electrons appear; while at $\vec{k}_0 = (0.45, 0.45)\pi$: all three approaches agree in that only c electrons appear, though the non-interacting band barely touches E_F while more c electrons are predicted to be present when correlations are included.

3.2. Momentum dependence of the spectral density $A(\vec{k}, \omega)$

In this section, we exhibit the spectral density function results obtained with our approach along two BZ paths, around the two relevant minima of the $\text{LaO}_{1-x}\text{F}_x\text{BiS}_2$ bandstructure, shown in Fig. 3.

First, in Fig. 4 we exhibit $A(\vec{k}, \omega)$ as a function of crystal momentum \vec{k} around X , at low temperature: $T=10$ K, and doping $n = x = 0.46$, parameters chosen as to compare our results directly with the only available ARPES data for $\text{La}_{1-x}\text{F}_x\text{BiS}_2$ compounds [23]. Fig. 4(a) shows the spectral density function calculated with the non-interacting two-orbital model of Usui et al. [13]: notice that along the BZ path considered around X it exhibits a two-peak

structure. The two well-defined and separated peaks, possessing equal spectral weight, represent the non-interacting(bare) c and d bands crossing E_F near X shown in Fig. 3. ARPES experiments for $\text{LaO}_{0.54}\text{F}_{0.46}\text{BiS}_2$ [23] indicate that the bottom of the lower conduction band E_{\min} is located near -0.75 eV, while in Fig. 4(a) and Fig. 3, in the absence of correlations it is placed near -0.37 eV. Interestingly, including renormalization effects by correlations with our second-order perturbative approach, we find that the agreement with ARPES improves: concretely, in Fig. 4(b) we exhibit our calculated renormalized spectral density function around X , yielding $E_{\min} \sim -0.7$ eV for $U = 1.5$ eV = $2V$ (as also noticeable in Fig. 3), closer to the ARPES result. Meanwhile, notice that the Hartree-Fock approximation overestimates the correction, yielding $E_{\min} \sim -1$ eV, and equal spectral weight for the two peaks. As seen in Fig. 4(b), our approach including correlations also captures the shift in spectral weight distribution between the two peaks along the BZ path, which is present in ARPES results [23].

As a prediction, in Figs. 5(a) and (b) we exhibit the non-interacting spectral density and our calculated renormalized spectral density function along the symmetry path $M - \Gamma$, around $\vec{k}_0 = (0.45\pi, 0.45\pi)$, respectively. This BZ region has not yet been probed by ARPES for $\text{LaO}_{1-x}\text{F}_x\text{BiS}_2$ compounds. Comparing the results in Fig. 4(b) with those in Fig. 5(b), one sees that the

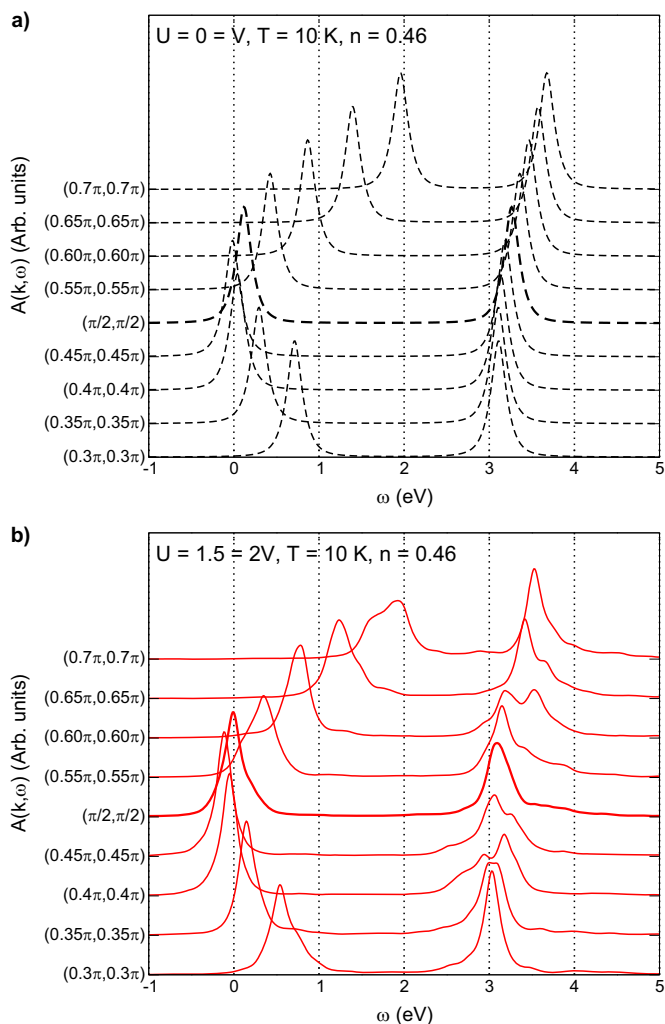


Fig. 5. Spectral density $A(\vec{k}, \omega)$ as a function of energy ω , for different BZ points (shown vertically displaced along the symmetry path $M - \Gamma$, around \vec{k}_0): (a) for the non-interacting two-orbital model [13], (b) present work: correlated two-orbital model, with our second-order perturbative approach. Parameters: $U = 1.5$ eV = $2V$, other parameters as in Fig. 1, except for: $\nu=9$.

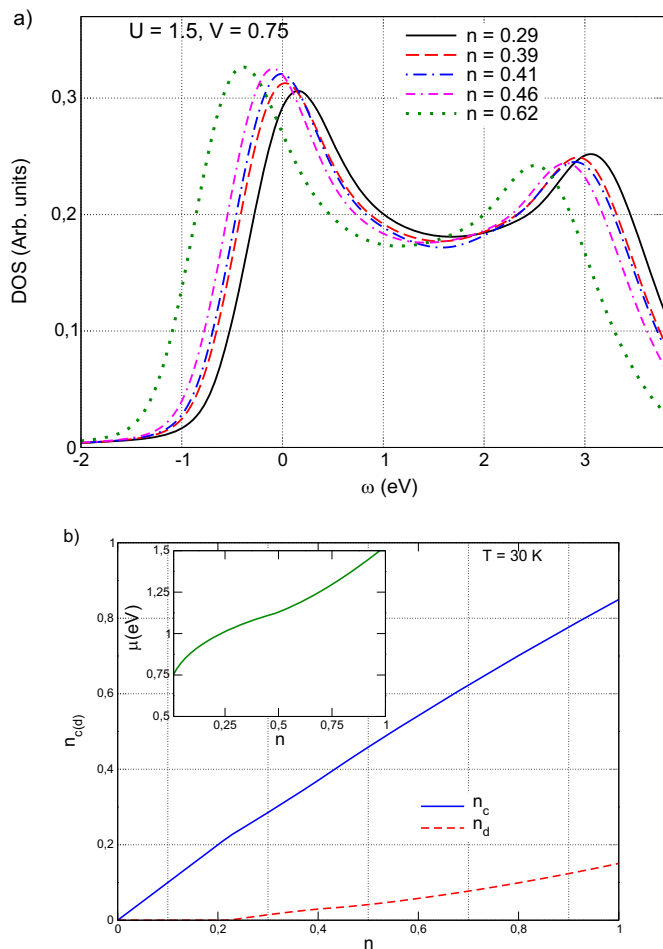


Fig. 6. Doping effect. (a) DOS $A(\omega)$, for different values of electron doping (indicated in the figure) at $T=10$ K. (b) Filling of the c and d bands, as a function of doping, at $T=30$ K. Inset: dependence of the chemical potential on total band filling at $T=30$ K. $U = 1.5$ eV = $2V$, and other parameters as in Fig. 1, except for: $\nu=9$.

renormalization effects by correlations around E_F are more relevant near X than near \vec{k}_0 .

3.3. Effect of electron doping

In this subsection, we focus on effects of electron doping on the spectral properties. In Fig. 6(a), we show our calculated total density of states for doping values n between 0.29 and 0.62. Increasing electron-doping, n , we find an almost rigid-band-like shift of the peak positions to lower energies, in agreement with results of SXPES [24] and DFT [25] near the Fermi level in $\text{LaO}_{1-x}\text{F}_x\text{BiS}_2$ with x between 0 and 0.5.

Our DOS results also describe an increase upon doping of the spectral weight of the peak located at lower energy (closest to E_F), as reported in the SXPES experiment [24]. This fact is reminiscent of the mentioned effect of the increase of U on the DOS and on the FS topology discussed in Section 3.1, also in connection with Ref. [26], and would justify the reduction upon increase of U of n_{crit} for the FS topology change. Nevertheless, in Fig. 6(a) we find a difference: because the increase of doping not only increases the spectral weight of the peak near E_F , as increasing U does, but it also shifts the peaks to lower energies as found in SXPES [24].

As mentioned before, the effective two-orbital model [13], which was determined taking into account mainly the Bi-6p and S-3p orbitals near E_F , is suitable to describe the low-energy properties of $\text{LaO}_{1-x}\text{F}_x\text{BiS}_2$. For a description of the DOS between

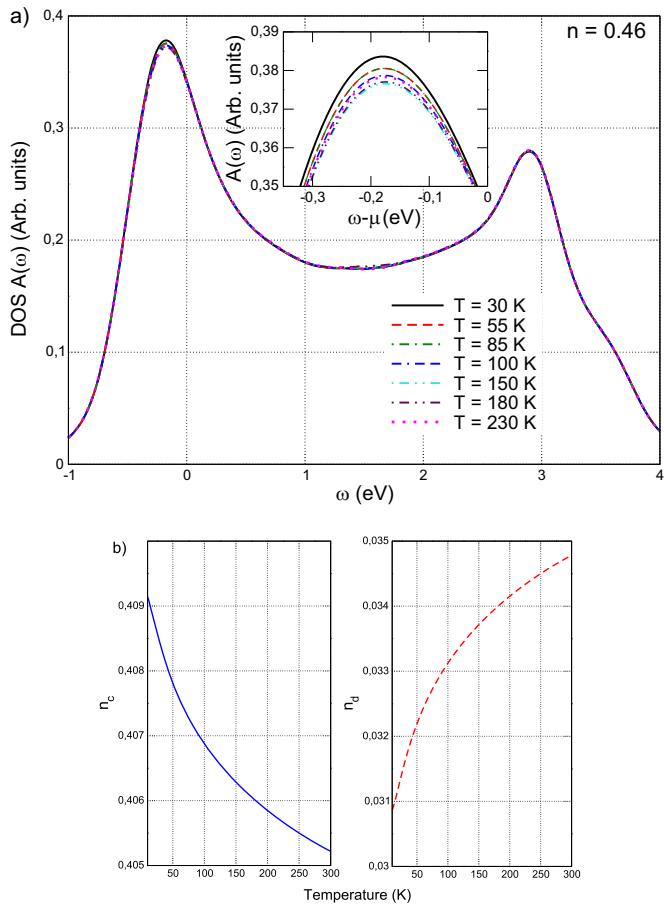


Fig. 7. Temperature dependence at $n=0.46$. (a) DOS at different temperatures (as indicated in the figure). Inset: amplification, showing the T-dependence of the peak near E_F . (b) c and d band fillings, as a function of temperature. Parameters: $U = 1.5$ eV = $2V$. Other parameters as in Fig. 1, except for: $\nu=9$.

1 eV and 4 eV, the inclusion of La and O orbitals would be important [25,3].

In Fig. 6(b), we exhibit the effect of doping on the filling of each effective band: though both show a monotonous increase of filling with n , as mentioned before most electrons occupy the lower c band. We find the d band almost empty below $n=0.2$. The inset of Fig. 6(b) shows the chemical potential as a function of n at low temperature $T=30$ K, exhibiting an inflection point at $n=0.44$.

3.4. Effect of temperature on the DOS and the spectral density

First, we will present our prediction for the temperature dependence of the total density of states of $\text{LaO}_{1-x}\text{F}_x\text{BiS}_2$ compounds. In Fig. 7(a), we show our calculated DOS for several temperatures at $n=0.46$ (doping value of sample studied by ARPES at $T=10$ K [23]).

Notice that the total DOS in Fig. 7(a), which includes the sum over contributions from the whole BZ, is nearly independent of temperature. The peak near E_F is shown amplified in the figure inset: increasing temperature a very slight reduction of spectral weight of the peak is observable. In Fig. 7(b) we further plot the corresponding band fillings, n_c and n_d , as function of temperature. Since for the doping value and temperatures considered, the Fermi level falls inside c and d bands, whose fillings change smoothly, the negligible temperature dependence of the total DOS is not surprising.

Next, we will focus on the temperature dependence of the spectral density function, predicting it for $\text{LaO}_{1-x}\text{F}_x\text{BiS}_2$ compounds

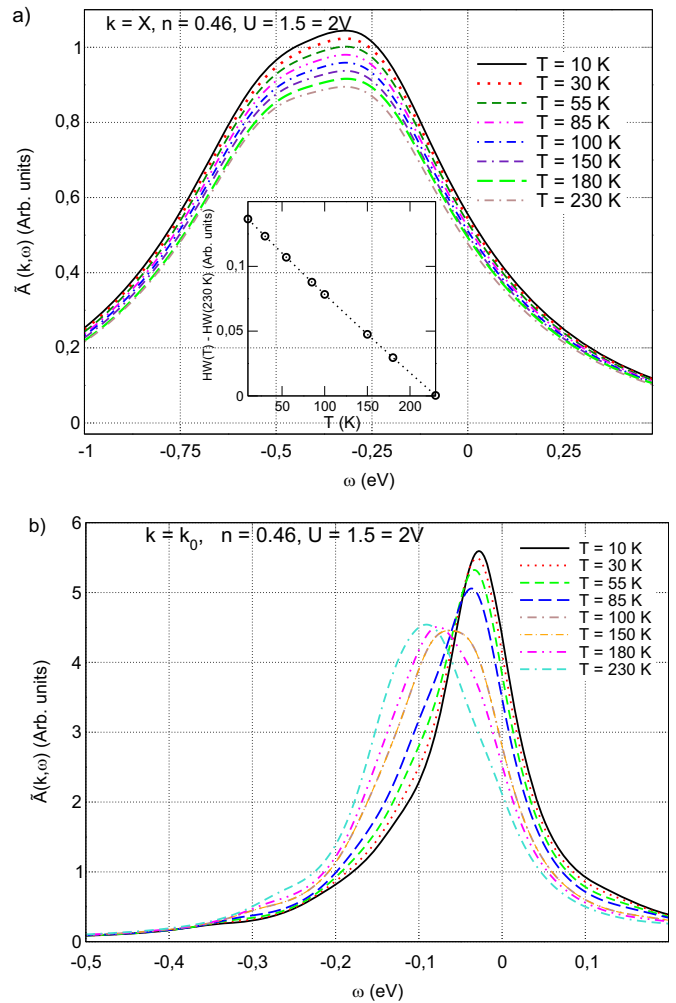


Fig. 8. (a) Temperature dependence of $\tilde{A}(\vec{k}, \omega)$, at $\vec{k} = X$ (for the temperatures indicated in the figure) and (b) Temperature dependence of $\tilde{A}(\vec{k}, \omega)$, at $\vec{k}_0 = (0.45\pi, 0.45\pi)$. $U=1.5$ eV, $V=0.75$, $n=0.46$. Other parameters as in Fig. 1, except for: $\nu=9$.

where no T-dependent ARPES results are available. We do make a comparison with the available temperature-dependent ARPES data for $\text{NdO}_{0.7}\text{F}_{0.3}\text{BiS}_2$ [41]. To allow for direct comparison with ARPES data [41], the spectral density results presented in Fig. 8, which are denoted by $\tilde{A}(\vec{k}, \omega)$, were obtained by the convolution of $A(\vec{k}, \omega)$ with the experimental energy resolution $\delta E = 25$ meV [41].

First, we predict the evolution with temperature of the spectral density function for $\text{LaO}_{1-x}\text{F}_x\text{BiS}_2$ compounds, fixing \vec{k} at the relevant high-symmetry points of the BZ, namely the conduction band minima X and \vec{k}_0 shown in Fig. 3. In Fig. 8(a), we show our calculated temperature dependence for $\tilde{A}(\vec{k} = X, \omega)$ at $n=0.46$, doping which corresponds to $\text{LaO}_{0.54}\text{F}_{0.46}\text{BiS}_2$ compound studied by ARPES only at the temperature $T=10$ K in Ref. [23]. In Fig. 8(a), the spectral density at X exhibits a broad slightly asymmetric hump, extending from about -1 eV to slightly above E_F , in agreement with Ref. [23] at $T=10$ K (and our results of Fig. 4b convoluted with $\delta E = 25$ meV). Increasing temperature, we find that the spectral weight is monotonically reduced, and no new peaks appear with temperature. The inset in Fig. 8(a) shows the integrated spectral weight of the hump as a function of temperature, with the subtraction of the integrated spectral weight of the hump at $T=230$ (denoted as $HW(T) - HW(230 \text{ K})$) as in the data analysis reported for the T-dependent ARPES experiment in

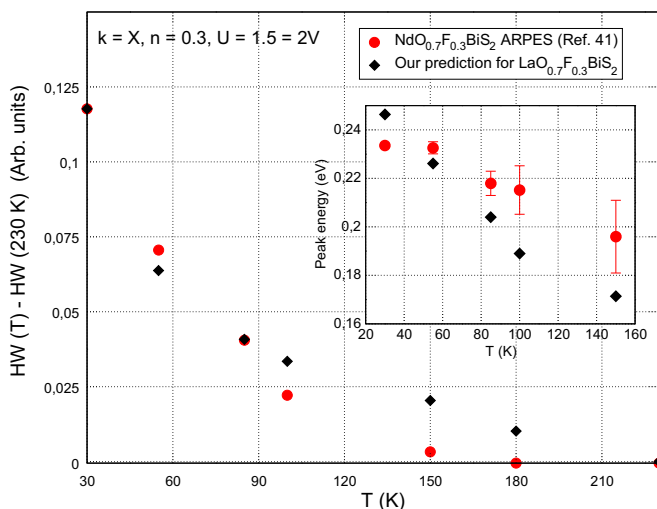


Fig. 9. Integrated spectral weight of $\vec{A}(\vec{k}, \omega)$, at $\vec{k} = X$ (temperatures are indicated in the figure). Inset: temperature dependence of the energy of the maximum of the hump. Experimental energy resolution of 25 meV, $U = 1.5$ eV, $V = 0.75$ eV, $n = 0.30$. Other parameters as in Fig. 1, except for: $\nu = 9$.

$\text{NdO}_{1-x}\text{F}_x\text{BiS}_2$ [41]. In Fig. 8(b), we predict the evolution of $\vec{A}(\vec{k}, \omega)$ with temperature at \vec{k}_0 , not yet probed by ARPES. The spectral function exhibits a single peak near E_F , which is shifted to lower energies as temperature is increased. Also, the standard thermal broadening effect with redistribution of the peak's spectral weight is visible. We checked that the total integrated spectral weight as a function of temperature remains constant. Between 100 and 150 K, we find negligible temperature dependence of $\vec{A}(\vec{k}, \omega)$ at \vec{k}_0 .

The temperature dependence of $\vec{A}(\vec{k}, \omega)$ we are predicting for $\text{LaO}_{0.54}\text{F}_{0.46}\text{BiS}_2$ in Fig. 8 could be probed by future ARPES experiments at X and at \vec{k}_0 .

Finally, in Fig. 9, we compare the temperature dependent spectral density results we predict for $\text{LaO}_{0.7}\text{F}_{0.3}\text{BiS}_2$, with the only available temperature-dependent ARPES experiment we found for a BiS_2 -based compound: $\text{NdO}_{0.7}\text{F}_{0.3}\text{BiS}_2$ [41]. This doping level corresponds to the optimal electron doping for superconductivity ($T_C = 4$ K) in the $\text{NdO}_{1-x}\text{F}_x\text{BiS}_2$ family. (To allow a direct comparison of the results included in Fig. 9, we have used the same data treatment adopted in Ref. [41]: to remove thermal broadening effects, the spectral data $\vec{A}(\vec{k}, \omega)$ were divided by resolution-convoluted Fermi functions at each temperature, and finally all spectra were multiplied by the Fermi function at $T = 10$ K.)

Concretely, in Fig. 9, we compare the temperature dependence of the integrated hump weight $HW(T) - HW(230 \text{ K})$ for both compounds, and in the inset compare the temperature dependence of the hump positions. Though a monotonous decrease of both magnitudes with temperature is observable, there are differences between our predictions for $\text{LaO}_{0.7}\text{F}_{0.3}\text{BiS}_2$, and the ARPES results for $\text{NdO}_{0.7}\text{F}_{0.3}\text{BiS}_2$ [41]. But these are to be expected, since our approach adds correlations to the effective two-orbital model with tight-binding parameters specifically determined for $\text{LaO}_{1-x}\text{F}_x\text{BiS}_2$ compounds [13]. No similar tight-binding parameters were determined specifically for $\text{NdO}_{1-x}\text{F}_x\text{BiS}_2$ compounds yet, and differences between the band structures are expected, according to DFT calculations [42]: while for $\text{LaO}_{1-x}\text{F}_x\text{BiS}_2$ the two band minima closest to E_F (the d -band minimum at X and the c -band minimum at \vec{k}_0) have very similar energies, this is not the case for $\text{NdO}_{1-x}\text{F}_x\text{BiS}_2$ where for $x = 0$ an energy difference of 0.8 eV was determined between the two relevant band minima by DFT [42].

4. Conclusions

We have studied the normal state spectral properties of $\text{LaO}_{1-x}\text{F}_x\text{BiS}_2$ superconductors, using an extended Hubbard model based on two correlated effective orbitals and an analytical approximation to decouple and solve the equations of motion for the electron Green's functions. Our results demonstrate that the inclusion of moderate electron correlations in $\text{LaO}_{1-x}\text{F}_x\text{BiS}_2$ improves the description of the available experimental results for these compounds. Concretely, including electron correlations with our second-order perturbative approach we could (i) improve the description of the chemical potential shift of approximately 0.3 eV between the $x = 0$ and $x = 0.5$ compounds reported in soft X-ray photoemission experiments at $T = 300$ K, by including moderate intra- and inter-orbital correlations: $U = 1.5$ eV, and $V = 0.75$ eV; (ii) describe the momentum dependence measured in ARPES around X ; (iii) and study the Fermi-surface topological change, which is independent of temperature: it occurs at $x = 0.44$ in the absence of correlations, while moderate correlations shift it to lower doping values.

Furthermore, our analytical approximation has the advantage of enabling us to describe temperature and k -dependent results, which allowed us to explore the spectral density at other BZ regions and in particular predict the momentum dependence around the second relevant band structure minimum located at $\vec{k}_0 = (0.45\pi, 0.45\pi)$. We also found that electron correlations enhance the nesting properties of the Fermi surface. Regarding the temperature dependence of the spectral function, our work predicts results which depend on the BZ point: at X we found that the position of the main hump near the Fermi level is unchanged, while an important hump spectral weight reduction takes place when temperature is increased. Instead at k_0 we found that temperature induces a shift of the position of the peak close to the Fermi level, while a spectral weight redistribution which leaves the total hump weight almost constant takes place. It would be interesting to have our predictions tested, by temperature-dependent ARPES experiments in $\text{LaO}_{1-x}\text{F}_x\text{BiS}_2$.

Our calculated Green's functions could be used to evaluate other low energy normal state physical properties of $\text{LaO}_{1-x}\text{F}_x\text{BiS}_2$: being our analytical approach especially useful for analyzing k -dependent results, in contrast to other theoretical techniques where self-energies are local.

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Appendix A. Analytical results for the c and d electron Green's functions

A.1. First order solution of EOM: Hartree-Fock

Below, we state the Hartree-Fock solutions obtained for the c and d electron Green's functions, as mentioned in Section 2.2:

$$G_{\sigma}^{H.F.}(k, \omega) \cong \frac{1}{2\pi \left[\omega - E_c(k) - (U\langle n_{c,\sigma} \rangle - V\langle n_d \rangle) \right]} \quad (\text{A.1})$$

and,

$$F_{\sigma}^{H.F.}(k, \omega) \cong \frac{1}{2\pi[\omega - E_d(k) - (U\langle n_{d,\sigma} \rangle - V\langle n_c \rangle)]} \quad (A.2)$$

A.2. Second-order solution of EOM

Below, we present the solution in second-order of perturbations on the electron correlations U and V for $G(\vec{k}, \omega)$ and $F(\vec{k}, \omega)$ obtained through the analytical approximation [30] described in Section 2.2:

$$G_{\sigma}(k, \omega) \cong \frac{\frac{1}{2\pi} + \sum_{k_2} \left\{ \frac{U}{N} \left(\frac{Y_1}{X_1} \right) + \frac{V}{N} \left[\frac{Y_2}{X_2} + [A_3]_{k_1} + [C_3]_{k_1} \left(\frac{Y_1}{X_1} + \frac{Y_2}{X_2} \right) \right] \right\}}{\omega - E_c(k) - \sum_{k_2} \left\{ \frac{U}{N} \left(\frac{Z_1}{X_1} \right) + \frac{V}{N} \left[\frac{Z_2}{X_2} + [B_3]_{k_1} + [C_3]_{k_1} \left(\frac{Z_1}{X_1} + \frac{Z_2}{X_2} \right) \right] \right\}} \quad (A.3)$$

To abbreviate, we use the following notation for summations: $[\alpha]_{k_1} \equiv \sum_{k_1} \alpha$, for different α coefficients. Meanwhile:

$$X_1(k_2, k) = \left(1 - \frac{U}{N} [C_2]_{k_1} [C_3]_{k_1} \right) (1 - [G_1]_{k_1} [C_3]_{k_1}) - [G_1]_{k_1} \left(\frac{U}{N} [C_3]_{k_1} + 1 \right) \left(\frac{U}{N} [C_3]_{k_1} + \frac{V}{N} \right) [C_2]_{k_1}$$

$$Y_1(k_2, k) = \left(1 - \frac{U}{N} [C_2]_{k_1} [C_3]_{k_1} \right) ([A_1]_{k_1} + [A_3]_{k_1} [G_1]_{k_1}) + \left([A_2]_{k_1} + \frac{U}{N} [C_2]_{k_1} [A_3]_{k_1} \right) [G_1]_{k_1} \left(\frac{U}{N} [C_3]_{k_1} + 1 \right)$$

$$Z_1(k_2, k) = \left(1 - \frac{U}{N} [C_2]_{k_1} [C_3]_{k_1} \right) ([B_1]_{k_1} + [B_3]_{k_1} [G_1]_{k_1}) + [B_2]_{k_1} + \frac{U}{N} [C_2]_{k_1} [B_3]_{k_1}$$

$$X_2(k_2, k) = (1 - [G_1]_{k_1} [C_3]_{k_1}) \left(1 - \frac{U}{N} [C_2]_{k_1} [C_3]_{k_1} \right) - \left(\frac{U}{N} [C_3]_{k_1} + V \right) [C_2]_{k_1} [G_1]_{k_1} ([C_3]_{k_1} + 1)$$

$$Y_2(k_2, k) = (1 - [G_1]_{k_1} [C_3]_{k_1}) \left([A_2]_{k_1} + \frac{U}{N} [C_2]_{k_1} [A_3]_{k_1} \right) + \left(\frac{U}{N} [C_3]_{k_1} + V \right) [C_2]_{k_1} ([A_1]_{k_1} + [A_3]_{k_1} [G_1]_{k_1})$$

$$Z_2(k_2, k) = (1 - [G_1]_{k_1} [C_3]_{k_1}) \left([B_2]_{k_1} + \frac{U}{N} [C_2]_{k_1} [B_3]_{k_1} \right) + \left(\frac{U}{N} [C_3]_{k_1} + V \right) [C_2]_{k_1} ([B_1]_{k_1} + [B_3]_{k_1} [G_1]_{k_1})$$

$$A_1 = \frac{\langle n_{k_1,\sigma} \rangle}{2\pi[\omega - E_c(k_2) - U(n_c + 1)]}$$

$$A_2 = \frac{\langle N_{k_1,\sigma} \rangle}{2\pi[\omega - E_c(k_2) - Un_c - 2n_d(U + V) + V]}$$

$$A_3 = \frac{\langle N_{k_1,\sigma} \rangle}{2\pi[\omega - E_c(k_2) - Un_c - 2Vn_d - V]}$$

$$B_1 = \frac{\{2V\langle n_{k_1,\sigma} \rangle n_d + (2 + \langle n_{k_1,\sigma} \rangle)Un_c + U(\langle n_{k_1,\sigma} \rangle - \langle n_{k_1-k_2+k,\sigma} \rangle)\}}{\omega - \omega_1}$$

$$B_2 = \frac{\{(U - V)\langle N_{k_1,\sigma} \rangle + 2V\langle N_{k_1,\sigma} \rangle n_d - V(\langle N_{k_1,\sigma} \rangle - \langle N_{k_1-k_2+k,\sigma} \rangle) + V\}}{\omega - \omega_2}$$

$$B_3 = \frac{\{U\langle N_{k_1,\sigma} \rangle n_c + V(1 - \langle n_{k_1,\sigma} \rangle - \langle n_{k_1-k_2+k,\sigma} \rangle)n_d + V\langle n_{k_1,\sigma} \rangle\}}{\omega - \omega_3}$$

$$G_1 = \frac{\frac{V}{N} [\langle n_{k_1,\sigma} \rangle - \langle n_{k_1-k_2+k,\sigma} \rangle]}{\omega - \omega_1}$$

$$G_2 = \frac{[\langle N_{k_1,\sigma} \rangle - \langle N_{k_1-k_2+k,\sigma} \rangle]}{\omega - \omega_2}$$

$$G_3 = \frac{[\langle N_{k_1,\sigma} \rangle - \langle N_{k_1-k_2+k,\sigma} \rangle]}{\omega - \omega_3}$$

$$\omega_1 = [E_c(k_1 - k_2 + k) - E_c(k_1) + E_c(k_2) + U(n_c + 1)]$$

$$\omega_2 = [E_d(k_1 - k_2 + k) - E_d(k_1) + E_c(k_2) + Un_c + 2n_d(U + V) - V]$$

$$\omega_3 = [E_d(k_1 - k_2 + k) - E_d(k_1) + E_c(k_2) + Un_c + 2Vn_d + V]$$

Analogously, for $F_{\sigma}(k, \omega)$ corresponding to the d-electrons we obtained:

$$F_{\sigma}(k, \omega) \cong \frac{\frac{1}{2\pi} + \sum_{k_2} \left\{ \frac{U}{N} \left(\frac{Y_1^*}{X_1^*} \right) + \frac{V}{N} \left[\frac{Y_2^*}{X_2^*} + [A_3^*]_{k_1} + [C_3^*]_{k_1} \left(\frac{Y_1^*}{X_1^*} + \frac{Y_2^*}{X_2^*} \right) \right] \right\}}{\omega - E_d(k) - \sum_{k_2} \left\{ \frac{U}{N} \left(\frac{Z_1^*}{X_1^*} \right) + \frac{V}{N} \left[\frac{Z_2^*}{X_2^*} + [B_3^*]_{k_1} + [C_3^*]_{k_1} \left(\frac{Z_1^*}{X_1^*} + \frac{Z_2^*}{X_2^*} \right) \right] \right\}} \quad (A.4)$$

where

$$X_1^*(k_2, k) = \left(1 - \frac{U}{N} [C_2^*]_{k_1} [C_3^*]_{k_1} \right) (1 - [G_1^*]_{k_1} [C_3^*]_{k_1}) - [G_1^*]_{k_1} \left(\frac{U}{N} [C_3^*]_{k_1} + 1 \right) \left(\frac{U}{N} [C_3^*]_{k_1} + \frac{V}{N} \right) [C_2^*]_{k_1}$$

$$Y_1^*(k_2, k) = \left(1 - \frac{U}{N} [C_2^*]_{k_1} [C_3^*]_{k_1} \right) ([A_1^*]_{k_1} + [A_3^*]_{k_1} [G_1^*]_{k_1}) + \left([A_2^*]_{k_1} + \frac{U}{N} [C_2^*]_{k_1} [A_3^*]_{k_1} \right) [G_1^*]_{k_1} \left(\frac{U}{N} [C_3^*]_{k_1} + 1 \right)$$

$$Z_1^*(k_2, k) = \left(1 - \frac{U}{N} [C_2^*]_{k_1} [C_3^*]_{k_1} \right) ([B_1^*]_{k_1} + [B_3^*]_{k_1} [G_1^*]_{k_1}) + [B_2^*]_{k_1} + \frac{U}{N} [C_2^*]_{k_1} [B_3^*]_{k_1}$$

$$A_1^* = \frac{\langle N_{k_1,\sigma} \rangle}{2\pi[\omega - E_d(k_2) - U(n_d + 1)]}$$

$$A_2^* = \frac{\langle n_{k_1,\sigma} \rangle}{2\pi[\omega - E_d(k_2) - Un_d - 2n_c(U + V) + V]}$$

$$A_3^* = \frac{\langle n_{k_1,\sigma} \rangle}{2\pi[\omega - E_d(k_2) - Un_d - 2Vn_c - V]}$$

$$B_1^* = \frac{\{2V\langle N_{k_1,\bar{\sigma}} \rangle n_c + (2 + \langle N_{k_1,\bar{\sigma}} \rangle)Un_d + U(\langle N_{k_1,\bar{\sigma}} \rangle - \langle N_{k_1-k_2+k,\bar{\sigma}} \rangle)\}}{\omega - \omega_1^*}$$

$$B_2^* = \frac{\{(U - V)\langle N_{k_1,\sigma} \rangle + 2V\langle n_{k_1,\sigma} \rangle n_c - V(\langle n_{k_1,\sigma} \rangle - \langle n_{k_1-k_2+k,\sigma} \rangle) + V\}}{\omega - \omega_2^*}$$

$$B_3^* = \frac{\{U\langle n_{k_1,\bar{\sigma}} \rangle n_d + V(1 - \langle N_{k_1,\bar{\sigma}} \rangle - \langle N_{k_1-k_2+k,\bar{\sigma}} \rangle)n_c + V(\langle N_{k_1,\bar{\sigma}} \rangle)\}}{\omega - \omega_3^*}$$

$$X_2^*(k_2, k) = \left(1 - [C_1^*]_{k_1} [C_3^*]_{k_1}\right) \left(1 - \frac{U}{N} [C_2^*]_{k_1} [C_3^*]_{k_1}\right) - \left(\frac{U}{N} [C_3^*]_{k_1} + V\right) [C_2^*]_{k_1} [C_1^*]_{k_1} ([C_3^*]_{k_1} + 1)$$

$$Y_2^*(k_2, k) = \left(1 - [C_1^*]_{k_1} [C_3^*]_{k_1}\right) \left([A_2^*]_{k_1} + \frac{U}{N} [C_2^*]_{k_1} [A_3^*]_{k_1}\right) + \left(\frac{U}{N} [C_3^*]_{k_1} + V\right) [C_2^*]_{k_1} ([A_1^*]_{k_1} + [A_3^*]_{k_1} [C_1^*]_{k_1})$$

$$Z_2^*(k_2, k) = \left(1 - [C_1^*]_{k_1} [C_3^*]_{k_1}\right) \left([B_2^*]_{k_1} + \frac{U}{N} [C_2^*]_{k_1} [B_3^*]_{k_1}\right) + \left(\frac{U}{N} [C_3^*]_{k_1} + V\right) [C_2^*]_{k_1} ([B_1^*]_{k_1} + [B_3^*]_{k_1} [C_1^*]_{k_1})$$

$$C_1^* = \frac{V [\langle N_{k_1,\bar{\sigma}} \rangle - \langle N_{k_1-k_2+k,\bar{\sigma}} \rangle]}{N (\omega - \omega_1^*)}$$

$$C_2^* = \frac{[\langle n_{k_1,\sigma} \rangle - \langle n_{k_1-k_2+k,\sigma} \rangle]}{\omega - \omega_2^*}$$

$$C_3^* = \frac{[\langle n_{k_1,\bar{\sigma}} \rangle - \langle n_{k_1-k_2+k,\bar{\sigma}} \rangle]}{\omega - \omega_3^*}$$

$$\omega_1^* = [E_d(k_1 - k_2 + k) - E_d(k_1) + E_d(k_2) + U(n_d + 1)]$$

$$\omega_2^* = [E_c(k_1 - k_2 + k) - E_c(k_1) + E_d(k_2) + Un_d + 2n_c(U + V) - V]$$

$$\omega_3^* = [E_c(k_1 - k_2 + k) - E_c(k_1) + E_d(k_2) + Un_d + 2Vn_c + V]$$

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