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Interaction of Glycine, Lysine, Proline and Histidine with Dipalmitoylphosphatidylcholine Lipid Bilayers: a Theoretical and Experimental Study[†]

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The interaction of unblocked glycine, lysine, proline, and histidine (in their three forms, namely two tautomers and the protonated form) with a dipalmitoylphosphatidylcholine (DPPC) bilayer was assessed using extensive atomistic Molecular Dynamics simulations. Free energy profiles for the insertion of each amino acid into the lipid bilayer were computed along an appropriated reaction coordinate. The simulation results for glycine in the presence of DPPC were compared with experimental data obtained by Fourier Transform Infrared Spectroscopy. Experimental results predict, in good agreement with simulations, the existence of intermolecular interactions between the DPPC head groups and glycine. Atomistic simulations were further extended to investigate the free energy profiles for lysine, proline and histidine, leading to the following conclusions: (*i*) lysine free energy profiles computed using a united atom force-field and an analog molecule, where the side-chain is truncated at the β -carbon atom, differ significantly from each other; (*ii*) the free energy profiles for the three forms of histidine are all very similar, although the charged form interacts mostly with the carbonyl groups of DPPC, while the tautomers interact with the phosphate groups; and (*iii*) proline does not show a minimum in the free energy profile, pointing to the absence of binding to the membrane lipids. Overall, this work contributes to our general understanding of the various factors affecting the interactions between amino acids and a model cell membrane, and may spur progress in the effort to develop new molecular models to study larger biological systems.

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1 Introduction

² Most of the experimental evidence regarding the interaction of $_{11}^{11}$ natural amino acids with lipid membranes is commonly inter-¹² preted in terms of the chemical nature of the side-chains. ^{1–8} ¹³ Hence, a usual approach in computer simulations of these sys-¹⁴ tems is to represent the amino acid as an "analog molecule", ¹⁵ ¹⁵ consisting of just the side-chain truncated at the β -carbon ¹⁶ ¹⁶ atom. ^{9–20} However, the use of the analog molecule approach ¹⁷

opens the question of to what extent the amino acid backbone influences its partitioning into a lipid bilayer. This problem is particularly important for residues that are not part of regular secondary structure elements in proteins, such as statistical coil fragments or loop regions, which constitute 50% of all residues in proteins.²¹ The fact that certain amino acids cannot be studied within the analog molecule approximation, e.g., glycine (Gly) and proline (Pro), only exacerbates the problem. Consequently, the aim of this work is threefold. First, to use Gly, which bears no side chain, as a reference compound for testing the additivity of backbone and side chain transfer free energies in all 20 naturally occurring amino acids (except for Pro). Therefore, atomistic Molecular Dynamics (MD) simulations and Fourier Transform Infrared spectroscopy (FTIR) are employed to investigate the nature of Gly-DPPC interactions at a molecular level. The experimental observations will enable us to assess the capabilities and potential limitations of the force-fields used in this work.

Second, to carry out MD simulations for the insertion of unblocked charged lysine into a DPPC bilayer, by using both a united-atom representation of the whole amino acid (Lys^+) and the analog molecule approach $(Lys^+-analog)$. The results of these simulations, together with those for Gly, will

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[†] Electronic Supplementary Information (ESI) available: Convergence criteria of the free energy profiles and experimental frequency values of the symmetric and antisymmetric stretching and bending modes, are provided. See DOI: 10.1039/b000000x/ ²¹

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be used to discuss the non-additivity of backbone and side 81 32 chain transfer free energies and, hence, the accuracy and limi- 82 33 tations of the analog molecule model. Although it is clear that 83 34 aqueous-organic transfer free energies cannot in general be de- 84 35 composed into molecular fragments' contributions, it is im- 85 36 portant to quantify non-additive effects for amino acid trans- 86 37 fers into lipid bilayers, given that analog molecule models are 87 38 widely used in biophysical simulations. As stated above, the 88 39 limitations of these models may be severe when transferring 89 40 residues which are not part of rigid structural motifs in pro-41 teins. 42

Third, to study the transfer free energy profile for both Pro⁹⁰ 43 and histidine (His). The reasons for choosing these two amino ₉₁ 44 acids are the following. Proline, which is an imino rather than " 45 an amino acid, does not admit an analog molecule represen-46 tation. Furthermore, His is special among all the ionizable 94 47 amino acids because it possesses a pK $^{\circ}$ = 6.6 and, hence, may $_{95}$ 48 be charged or neutral around pH 7.0 where most of the bio-96 49 logical processes occur. Moreover, for the neutral form of His $_{\rm q7}$ 50 two tautomers exist, namely N^{δ 1}-H and N^{ϵ 2}-H, which need to ₉₈ 51 be discussed separately. 52

The rest of the paper is organized as follows: simulation₁₀₀ and experimental methods are detailed in Section 2, results₁₀₁ are presented and discussed in Section 3, and conclusions are₁₀₂ summarized in Section 4.

57 2 Methodology

58 2.1 Experimental

108 Synthetic 1, $2-\frac{100}{109}$ 2.1.1 Lipid Sample Preparation. 59 dipalmitoyl-sn-glycero-3-phosphocholine and unblocked Gly 60 with > 99% and > 98% purity, respectively, were purchased $\frac{1}{111}$ 61 from Sigma-Aldrich and used without further purification. 62 The lipids dissolved in chloroform were dried to form a film $\frac{1}{113}$ 63 under a nitrogen stream to study the interaction of Gly with the $\frac{1}{114}$ 64 phospholipids. The lipid film was left 24 hours under vacuum 65 to ensure the proper removal of solvents. Lipids were rehy-66 drated in de-ionized triple-distilled water, and in solutions of 67 different concentrations (25, 50, 100, 150 and 200 mM) pre-68 pared in H_2O or D_2O , above the gel/liquid-crystal phase tran-69 sition temperature (323.2 K), gently shaking for 15 minutes to 70 produce multilamellar vesicles (MLV's). The final concentra-71 tion of MLV's was 0.05 mg/ μ l or 50 mg/ml.²² 72 122

2.1.2 Measurements. FTIR spectra were recorded in 123 73 transmission mode in a system continuously purged with dry124 74 air, on a Perkin Elmer 1600 spectrophotometer provided with 125 75 a DTGS detector. The equipment was coupled to a SPV1.0126 76 system that transfers energy by means of a semiconductor cell127 77 working with the Peltier effect. The infrared spectra of lipo-128 78 somes were obtained co-adding 64 scans with 1 cm^{-1} resolu-129 79 tion using ZnSe windows. The working temperature range130 80

was $298.2 - 323.2 \pm 0.5$ K, and the spectra were analyzed using the GRAMS/32 mathematical software (Hertfordshire, UK). The contours of the C=O stretching bands (vC=O) were obtained by Fourier self deconvolution using band width parameters between 18 and 20 cm⁻¹ and a band narrowing factor of 2, as defined by the mathematical software GRAMS/32 Spectral Notebase. Deconvolution was used to determine the position of the bands corresponding to the two populations of carbonyl groups in the gel state.^{23–25}

2.2 Molecular Dynamics Simulations

MD simulations were used to investigate the insertion of Gly, Pro, the three forms of His, and charged Lys (both the united atom model and the analog side-chain molecule), into a DPPC lipid bilayer. The unblocked amino acids were modeled with the GROMOS 53a6 force field.²⁶ For the Lys⁺ analog, the side chain was truncated at the β -carbon atom. In the united atoms representation of GROMOS 53a6, this was achieved by replacing the original methylene group associated to the β carbon with a united atom methyl group. The charged forms of the amino acids were neutralized by including a counterion (Cl⁻) in the simulation cell. DPPC was modeled using the force-field proposed by Berger et al.²⁷ combined with the Single Point Charge (SPC) water model.²⁸ Each simulation box contained 64 DPPC molecules (32 lipids per leaflet), approximately 3815 water molecules (full hydration), and one amino acid initially located at the center of the water slab (z =3.5 nm from the membrane center). The bilayer normal was perpendicular to the x - y plane of the coordinates system.

Since the timescale for the spontaneous penetration of the amino acid into the bilayer is large compared to the simulation time, an external force was applied to the amino acid in order to generate initial configurations for the subsequent free energy calculations. A harmonic potential with a force constant of 3000 kJ.mol⁻¹.nm⁻² was applied to the reaction coordinate, defined as the z-component of the distance vector between the center of mass of the amino acid and the center of mass of the lipid bilayer^{18,29,30}. The amino acid was thrust into the lipid bilayer at a rate of approximately 7 nm/ns, and was allowed to move freely on the x - y plane. The Potential of Mean Force (PMF) for the penetration of the amino acid was computed by Umbrella Sampling³¹ using a set of 36 windows spanning the reaction coordinate interval 0.0-3.5 nm. Each window was let to relax for 10 ns, and then simulated for over 100 ns. Free energy profiles were recovered with the Weighted Histogram Analysis Method (WHAM), ^{32,33}. Convergence was assessed by applying WHAM on consecutive trajectory blocks of 20 ns (see Figs. 1-6 in the Supporting Information).

All simulations were performed with the GROMACS-4.5.5 package, ^{34,35} using a time step of 2 fs. Lennard-Jones inter-

actions were cutoff at 1 nm, and dispersion corrections were 180 131 applied to energy and pressure in order to account for the181 132 pair-potential truncation. Long range electrostatic interactions182 133 were evaluated using the particle mesh Ewald method, ^{36,37}₁₈₃ 134 with real space interactions cutoff at 1 nm, and reciprocal₁₈₄ 135 space interactions computed on a 0.16 nm grid with a fourth-136 order spline interpolation. At the beginning of each simu-137 lation, a steepest descent minimization process was applied 138 to the whole system in order to remove any excess of strain 139 and potential overlaps between neighboring atoms. Produc-140 tion runs were performed in the NPT thermodynamics ensem-141 ble, using as a thermostat the velocity rescaling algorithm of 142 Bussi et al.,³⁸ and a weak pressure coupling algorithm for the 143 barostat.³⁹ The pressure was always set to 1 atm and the tem-144 perature to 323 K (above the phase transition temperature, 314 145 K, of DPPC).⁴⁰ The coupling constants for the thermostat and 146 the barostat were 0.1 ps and 1 ps, respectively. 147

3 Results and Discussion

149 3.1 FTIR Experiments

Gly is one of most abundant amino acids in nature, and is also 150 involved in several biological processes. More importantly, 151 Gly has no lateral chain and makes an appropriate model for 152 investigating the role of the backbone on the interaction of 153 amino acids with lipid bilayers. This is particularly relevant 154 from the point of view of molecular simulations, as the ab-155 sence of a side-chain allows us to assess the adequacy of the 156 analog molecule approach. Moreover, the comparison be-157 tween simulations and experiments delimits the scope of the 158 force-field and the computational strategy employed in this 159 work. 160

3.1.1 **Hydrophobic region.** The symmetric stretching of 161 the fatty acid methylene groups (v_s CH₂) was studied in or-162 der to determine the effect of Gly on the hydrophobic region 163 of the lipid bilayer. This vibrational mode is reported to oc-164 cur at 2850 cm⁻¹, and is of great importance due to its sensi-185 165 tivity to mobility changes and to the conformational disorder186 166 of the hydrocarbon chains. The maximum absorption of this187 167 band shifts to higher frequencies when the membrane becomes 188 168 fluid (for example, when the hydrocarbon chains gauche ro-189 169 tamer population increases with respect to the *trans* rotamer₁₉₀ 170 population). This frequency shift occurs at the phospholipid₁₉₁ 171 transition temperature $(Tm = 314.65 \text{ K})^{41,42}$. Fig. 1 shows¹⁹² 172 that the Tm of pure DPPC agrees with the value reported in₁₉₃ 173 the literature⁴². No substantial changes were observed for li-194 174 posomes prepared in H₂O or D₂O with different Gly:DPPC₁₉₅ 175 molar ratios. This indicates that the gel phase of DPPC is not₁₉₆ 176 altered by the presence of Gly (see Fig. 1 and Table 1). In197 177 addition, no significant shifts, within the experimental error \pm_{198} 178 1 cm⁻¹, were observed in the symmetric, antisymmetric and 199 179

bending modes of the methyl and methylene groups of the inner lipid bilayer in the gel phase (measured at 298.2 K), nor in the crystalline liquid phase (measured at 323.2 K). Frequencies and frequency changes at both temperatures are reported in Tables 1 and 2 of the Supporting Information.



Fig. 1 Changes in vibrational frequency of the CH₂ symmetric stretching mode in Gly:DPPC (at different molar ratios), as a function of temperature. Gly:DPPC molar ratios: (\Box) 0.0:1, (\bigcirc) 0.4:1, (\triangle) 0.9:1, (\bigtriangledown) 2.0:1 and (\diamondsuit) 3.0:1

 $\label{eq:Table 1} \begin{array}{l} \mbox{Table 1 Phase transition temperature (Tm) in Gly:DPPC liposomes} \\ \mbox{(at different molar ratios), both in H_2O and D_2O \\ \end{array}$

Molar ratio	H ₂ O	D ₂ O
Gly/DPPC	(K)	(K)
0.0:1	314.7	315.1
0.4:1	314.7	314.7
0.9:1	314.7	314.7
2.0:1	314.1	314.7
3.0:1	313.7	313.7
4.0:1	315.6	315.6

3.1.2 Hydrophilic or interphasial region. It has been reported that the carbonyl ester linking the glycerol backbone with the fatty acid chains and the phosphate groups are the main hydration sites of phosphatidylcholines^{24,43}. Gly-DPPC spectra were registered in D₂O to assign the C=O stretching mode frequency (vC=O), and in H₂O, to assign the PO₂⁻ vibrational mode frequencies.

It is well known that the main vC=O peak of diacyl lipids can be decomposed into at least two components. One of them corresponds to the H-bonded and the other to the nonbonded (free) conformers of the C=O group⁴⁴. The higher frequency band component (1740 - 1742 cm⁻¹) has been assigned to free vC=O groups ($vC=O_f$), whereas the lower frequency component (~1728 cm⁻¹) has been attributed to the vC=O vibration of H-bonded conformers ($vC=O_b$)⁴⁵. Deconvolution and curve fitting were performed to determine the position and
relative contribution of the two carbonyl populations. A large
set of spectra and fitting curves are shown in Fig. 1 of the
Supporting Information (SI).

Figs. 2, 3 and 4 depict the frequency shifts of bonded 204 and free C=O groups ($vC=O_b$ and $vC=O_f$, respectively) for 205 Gly:DPPC ratios between 0.0:1 and 4.0:1. Three different 206 temperatures are considered; 298.2K, corresponding to the gel 207 phase (Fig. 2); the transition temperature 314.2 K (Fig. 3); and 208 323.2K corresponding to the liquid crystalline phase (Fig. 4). 209 Numerical values are provided in Table 1 of the SI. Below a 210 Gly:DPPC molar ratio of 3.0:1, neither $vC=O_b$ nor $vC=O_f$ 211 depict a noticeable shift with respect to the pure lipid at both 212 298.2 and 323.2K. However, at the 4.0:1 molar ratio, a smooth 213 shift to lower frequencies is observed for both carbonyl popu-214 lations (Figs. 2 and 4 and Table 1 of the SI). 215



Fig. 2 Frequency shifts of: (\Box) $vC=O_f$, (\bigcirc) $vC=O_b$, (\blacktriangledown) $v_{as}PO_2^-$ 229 and (\blacktriangle) $v_sPO_2^-$, stretching vibrational mode as a function of the Gly:DPPC molar ratio at 298.2 K (gel state). 231



Fig. 3 Frequency shifts of: (\Box) $vC=O_f$, (\bigcirc) $vC=O_b$, (\blacktriangledown) $v_{as}PO_2^-$ ²⁴⁶ and (\blacktriangle) $v_sPO_2^-$, stretching vibrational mode as a function of the Gly:DPPC molar ratio at 314.2 K (transition state). ²⁴⁸



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Fig. 4 Frequency shifts of: (\Box) $vC=O_f$, (\bigcirc) $vC=O_b$, (\blacktriangledown) $v_{as}PO_2^$ and (\blacktriangle) $v_sPO_2^-$, stretching vibrational mode as a function of the Gly:DPPC molar ratio at 323.2 K (liquid crystalline state).

 $vC=O_f$ to the carbonyl stretching mode, taken from Fig. 1 of the SI, as a function of the Gly:DPPC molar ratio. In the gel state (298.2K) the contribution of $vC=O_b$ is greater than that of $vC=O_f$ for Gly:DPPC molar ratios between 0.0:1 and 3.0:1. However, the trend inverts at the molar ratio 4.0:1 indicating a saturation of the interphase with Gly molecules (see also Figs. 1a I to VI of SI) . At the transition temperature (314.2 K) and in the liquid crystalline phase (323.2K), the contributions of $vC=O_b$ and $vC=O_f$ in the pure lipid (Gly:DPPC: 0.0:1) are almost the same, but $vC=O_b$ becomes clearly dominant when increasing the Gly:DPPC ratio up to 0.9:1 (314.2K) and 2.0:1 (323.2K). It must be pointed out that each plot in Fig. 5 (a, b or c) shows the results of experiments carried out at the same temperature (298.2, 314.2 and 323.2K). In other words, for each set of experiments only the Gly concentration increased and the contribution of water did not change with respect to the pure lipid (Gly:DPPC 0.0:1). Therefore, one could infer that the increase in $vC=O_b$ contribution indicates the formation of hydrogen bonds between Gly and DPPC.

The observations reported in the previous paragraph can be summarized stating that the presence of Gly leads to noticeable changes in $vC=O_f$ and $vC=O_b$ (see Fig. 5 and Fig. 1 of the SI). The evolution of both carbonyl populations was more evident in the fluid state. Assuming that the relative area of a band component is proportional to the respective conformer population, it can be concluded that the populations of C=O_{bond} and C=O_{free} conformers change upon addition of Gly.

The asymmetric stretching mode of the phosphate group (PO_2^-) shifts to lower frequencies in hydrated lipids ^{23–25,46,47}. This shift has been ascribed to direct H-bonding of water molecules to the charged phosphate groups. Therefore, PO_2^- has been suggested to act as a sensor of the hydration level of

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Fig. 5 Contribution of $vC=O_f$ and $vC=O_b$ to the carbonyl population at different Gly:DPPC molar ratios and at (a) 298, (b) 314.2 and (c) 323.2 K. Symbols indicate: (\blacksquare) $vC=O_b$, and (\bigcirc) $vC=O_f$.

278 the interphase $^{23-25,46}$. Gly had a different quantitative effect₂₇₉ 251 on the PO₂⁻ stretching bands, depending on whether the bi-280 252 layer was in the gel (298.2 K) or in the liquid crystalline state₂₈₁ 253 (323.2 K) (see Figs. 2, 4 and 6). Independently of the phase₂₈₂ 254 state of the membrane, the presence of the amino acid did $\operatorname{not}_{_{283}}$ 255 show a substantial effect on the PO_2^- symmetric stretching₂₈₄ 256 band. However, the antisymmetric stretching mode showed 285 257 an important shift (Δv) from -7 to -9 cm⁻¹ in the gel phase, 258 and around -2 to -4 cm^{-1} in the liquid crystalline state (see₂₈₇ 259 Fig. 2 and 4). These observations suggest that in addition to_{288} 260 the replacement of water molecules, there is participation of_{289} 261 the PO_2^- groups in the interaction with Gly through H-bonds.₂₉₀ 262 As can be inferred from the data reported in Fig. 3, at the₂₉₁ 263 transition temperature (314.2 K) and at all Gly molar ratios2992 264 assayed, there are no significant changes on the antisymmetric 2983 265 and symmetric stretching mode frequencies of the PO₂⁻ group.₂₉₄ 266 Overall, the results reported in this section reveal that: $(i)_{295}$ 267 the addition of Gly does not alter the fluidity (order of the hy-296 268 drocarbon chains) of the lipid membrane, and (ii) the presence₂₉₇ 269 of specific interactions between the head groups of DPPC and₂₉₈ 270 Gly. These observations strongly suggest that the phosphate299 271



Fig. 6 IR spectra for various Gly:DPPC molar ratios in the $1300 - 1000 \text{ cm}^{-1}$ region.

groups of the lipid membrane form H-bonds with Gly, in replacement of the water molecules, both in the gel and in the liquid crystalline states.

3.2 Molecular Dynamics Simulations of Gly and Lys

In order to facilitate the description of the free energy profiles, partial local mass density profiles were computed for the simulated bilayer system, and are shown in Fig. 7. Based on these profiles the bilayer is divided into four regions, according to the model used by MacCallum et. al¹⁸. Region 1 corresponds to bulk water with a small population of DPPC head groups; region 2 contains most of the charged phosphate and choline atoms; region 3 is a mix between the final portion of the polar head groups and the beginning of the lipids tails; and region 4 includes only the hydrophobic tails.

The free energy profile for inserting Gly into the DPPC bilayer is plotted in Fig. 8. The curve displays a minimum at approximately 1.7 nm from the center of the bilayer and near the core of region 2 (head groups). The free energy gain to bring the amino acid from bulk water to the lipid surface is $\sim -40 \text{ kJ.mol}^{-1}$. After the minimum, the free energy rises up to $\sim 50 \text{ kJ.mol}^{-1}$, as the amino acid approaches the center of the bilayer. The general features of the free energy profile of Fig. 8 indicate that Gly adsorption on DPPC occurs spontaneously, while its partitioning to the center of the membrane is highly unfavorable. The strong surface binding of the amino acid to the bilayer agrees with the trends discussed in section 3.1, which suggested that the strongest Gly-DPPC interactions occur at the level of the polar head groups rather than in the



Fig. 7 Partial mass density profile of the simulated system. Whole system (- - -), DPPC (---), water $(\cdots \cdot \cdot)$, lipids' carbonyl groups $(- \cdot - \cdot)$, head groups (\diamondsuit) , Phosphorous (\bigcirc) and Nitrogen (\Box) . The vertical lines and numbers divide the system into four regions (see text for details).

300	membrane core. The adsorption of Gly on DPPC, and its spe-
301	cific interaction with the phosphate groups, is also supported
302	by the simulations and experiments reported in reference 48 .



 Fig. 8 Transfer free energy profile for Gly. Vertical lines divide the 331

 system into 4 regions (see Fig. 7). Error bars are standard errors calculated by splitting a 100 ns MD-US trajectory into 5

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 independent blocks.

In order to characterize changes in the bonding pattern as336 303 the amino acid penetrates into the membrane, the number H-337 304 bonds between Gly and water, Gly and DPPC and between338 305 DPPC and water was computed as a function of the reaction₃₃₉ 306 coordinate z. Fig. 9 demonstrates that the number of Gly-340 307 water H-bonds decreases as the amino acid moves into the bi-341 308 layer, *i.e.*, as it gets into the hydrophobic region of the mem-342 309 brane. On the other hand, Fig. 9 shows that the number Gly-343 310 DPPC H-bonds (including bonds to the phosphate and to the₃₄₄ 311

carbonyl groups), reaches a maximum in the region of the hydrophilic heads, and decreases slightly as the molecule moves towards the hydrophobic core. It is then clear from Fig. 9 that after traversing region 2 (see Fig. 7), Gly remains partially hydrated and coordinated to a single DPPC head group. This was also confirmed by the inspection of simulation snapshots (see panel A of Fig. 10). For completeness, Fig. 9 shows that the average number of DPPC-water hydrogen bonds changes very little during the insertion of the amino-acid. This can be attributed to the fact that the local perturbation induced by Gly on the DPPC-water interface, is small compared to the total extend of the interface.



Fig. 9 Calculated number of hydrogen bonds formed between: (\bigcirc) Gly-water (\diamondsuit) DPPC- water (see text for further explanation) and (\Box) Gly-DPPC. Error bars calculated from 5 independent simulations.

As stated in Section 1, one of the motivations of the present work is to assess the impact of the analog molecule approach, where amino acids are represented by their side chains, on the thermodynamic work required to bring the molecule from the bulk of the solvent to the surface, or to the center, of the membrane. With this purpose, lysine was chosen as a case study because it possesses a large and flexible side chain, and it also plays important roles in membrane protein activity²⁰.

The free energy profiles for the Lys⁺-analog and the corresponding whole molecule model, are shown in Fig. 11. As observed, the two free energy curves depict the same general trends (a deep minimum near the membrane surface, and a local maximum at the membrane center), but also significant quantitative differences. In particular, the Lys⁺-analog (solid curve in Fig. 11) shows almost no free energy change when transferring the molecule from water to the center of the lipid bilayer. Also the free energy minimum (of ~ -57 kJ mol⁻¹) is located within region 3 of the bilayer.

It must be noticed that there is currently a degree of dispersion in the binding energy of amino acid analogs to PC bilayers as predicted by different force-fields. For example, Mac-

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Fig. 10 Configurational snapshots extracted from Umbrella ³⁶⁹ Sampling windows located at the center (z = 0) of the DPPC bilayer.³⁷⁰ (A) Gly, (B) Lys⁺-analog and (C) whole molecule model of Lys⁺. ³⁷¹ The orange beads represent the phosphorous atoms of DPPC. Water ³⁷² molecules are represented by wires. ³⁷³

Callum et al. have reported a binding energy of around 20 kJ₃₇₆ 345 mol⁻¹ for the Lys⁺-analog on DOPC, based on a combination³⁷⁷ 346 of Berger's and the OPLS force-field¹⁸. On the other hand, Li³⁷⁸ 347 et al. have found no adsorption (no minimum in the PMF) of 379 348 the Lys⁺-analog on DPPC on the basis of CHARMM. Simi-380 349 lar trends have been reported for charged arginine analogs on³⁸¹ 350 PC lipids^{15,16,18,20}. Considering the current discussions in the³⁸² 351 literature, and the state of the art in computer simulations of³⁸³ 352 protein-lipids systems, we believe that the binding of charged³⁸⁴ 353 amino acids to zwitterionic membranes is still a topic under 354 scrutiny and debate. The present results provide further ele-386 355 ments for judgment. 356 387

For the whole-molecule model of Lys⁺, the maximum in₃₈₈ the free energy curve (see dashed line on Fig. 11) occurs at₃₈₉



Fig. 11 Free energy profiles for Lys^+ (- - -), the Lys^+ analog (_____), and the difference between the PMFs of Lys^+ and Gly (- · - ·). Vertical lines divide the system into 4 regions (see text for further explanation). Error bars are standard errors computed by splitting a 100 ns MD-US trajectory into 5 independent blocks.

the center of the bilayer (hydrophobic region) and, taking as reference the analog molecule model, the minimum is shifted towards the water-lipid interface. This points to the existence of strong interactions between Lys⁺ and the polar head groups of DPPC. Fig. 11 also shows a difference of $\sim 20 \text{ kJ}.\text{mol}^{-1}$ in the equilibrium adsorption energy, in favor of the whole Lys⁺ molecule. Such a difference could be traced to a combination factors. The presence of the zwitterionic backbone in the whole-molecule model introduces additional lipid amino acid interactions, and also leads to a larger and tighter hydration shell. Furthermore, the backbone may decrease the conformational freedom of the side-chain and, hence, reduce the entropic contribution to the free energy change. Also, a careful analysis of simulation snapshots reveals that when the molecule is inside the bilayer, the average orientation of the side chain predicted by the two models is different. The isolated chain orients its axis perpendicular to the membrane (x - y) plane, whereas in the whole molecule model the side chain lies on the bilayer plane. Also, panels B and C of Fig. 10 show that the presence of the zwitterionic backbone induces the formation of a water channel across the membrane, which is absent in the isolated chain simulations. Such a transmembrane defect not only modulates the free energy cost of transferring the molecule to the center of the bilayer, but also conditions the orientation of the amino acid inside the membrane.

Going back to Fig. 11, the dot-dashed line represents the difference between the PMF of the whole Lys⁺ model and that of Gly. Could the transfer free energy of the amino acid be unambiguously decomposed into a backbone and a side chain contribution, the difference in PMF of Fig. 11 should resem-

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ble the profile of the analog molecule model. Clearly, that is₄₂₀ 390 not the case, highlighting the non-additivity of backbone and₄₂₁ 391 side chain contributions to the transfer free energy. Although₄₂₂ 392 analog molecules can be a good approximation to study the₄₂₃ 393 membrane insertion of rigid portions of a macromolecule (*i.e.*:424 394 α -helix in proteins), they may be inaccurate to represent the₄₂₅ 395 energetic of transferring the most flexible parts (*i.e.*: loops,426 396 turns), which include $\sim 50\%$ of all amino acidic residues in 397 proteins. 398 427

399 3.3 Molecular Dynamics Simulations of Pro and His

The transfer free energy profile for Pro is plotted in Fig. 12.430 400 The maximum of the free energy curve occurs at the center of_{431} 401 the bilayer and is $\sim 60 \text{ kJ.mol}^{-1}$; while the minimum of $\sim -9_{432}$ 402 kJ.mol⁻¹, which is quite weak for a polyatomic molecule, is₄₃₃ 403 located at the boundary between regions 2 and 3, *i.e.* close to_{434} 404 the carbonyl groups, and at 1.5 nm from the center of the bi-435 405 layer. Clearly, these results point to an unfavorable interaction₄₃₆ 406 between Pro and DPPC. 407 437



Fig. 12 Free energy profiles for Pro $(- \cdot \cdot -)$ and the three forms of His, namely, two tautomers $[N^{\delta 1}$ -H (----) and $N^{\epsilon 2}$ -H (- - -)], and the protonated form $[His^+ (- \cdot - \cdot)]$. Error bars are standard errors ⁴⁴⁸ calculated by splitting a 100 ns MD-US trajectory into 5 ⁴⁴⁹ independent blocks. ⁴⁵⁰

The free energy profiles for the three forms of His are also⁴⁵² 408 shown in Fig. 12. Overall the two tautomers ($N^{\delta 1}$ -H and $N^{\epsilon 2}_{-453}$ 409 H) and the ionized form (His⁺) show very similar trends, with⁴⁵⁴ 410 a minimum near the DPPC head groups and a global max-455 411 imum at the center of the bilayer. Such level of similarity₄₅₆ 412 will be explained and discussed in detail in Section 3.4. In457 413 the mean time, a few minor difference between the curves of 458 414 Fig. 12 are worth mentioning. The free energy maxima for459 415 the three forms of His have values of \sim 50 kJ.mol⁻¹, both for₄₆₀ 416 $N^{\delta 1}H$ and $N^{\epsilon 2}H$, and $\sim 40 \text{ kJ.mol}^{-1}$ for His⁺. In the neu-461 417 tral tautomers, the depth of the minima differ in only $\sim 5_{462}$ 418 kJ.mol⁻¹, and are located close to 1.8 nm (near to the phos-463 419

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phate groups). However, for the ionized form (His⁺), the free energy minimum is ~ 5 kJ.mol⁻¹ and ~ 11 kJ.mol⁻¹ deeper than for N^{δ 1}-H and N^{ϵ 2}-H, respectively. Also this minimum is located at the boundary between regions 2 and 3, suggesting specific interactions with the carbonyl groups of DPPC. Overall, our results imply that the three forms of His adsorb spontaneously on the surface of the DPPC bilayer.

3.4 Analysis of Hydration/Dehydration

Previous simulations of amino acid insertion into lipid bilayers have shown the existence of water molecules trapped into the membrane^{15,18,49,50}, when the molecule reaches the center of the lipid bilayer. As this effect is also observed in simulations of whole-molecule models, it is worth investigating the impact of the amino acid representation on the amount of hydration water as a function of the reaction coordinate. For example, panels B and C of Fig. 10 already suggest some significant differences in solvation between a whole Lys⁺ molecule and its analog.

Once an atom "X" of a given amino acid is chosen as a reference point (*e.g.*, the carbonyl oxygen of the backbone), the number of hydrating water molecules can be calculated from the radial distribution function,

$$g_{XO}(r) = \frac{N_O(r)}{4\pi r^2 \rho \delta r} \tag{1}$$

integrated up to the first minimum^{51,52}. In this equation, *X* represents the reference atom type in the amino acid, *O* is the oxygen atom of water, $N_O(r)$ is the number of *O* atoms located in a spherical shell of thickness δr and radius *r* measured from *X*, while ρ is the *O* number density. We define the dehydration number (Δ_{HN}) of the amino acid as:

$$\Delta_{HN}(z) = N_O^i - N_O^{bulk} \tag{2}$$

where N_O^i is the number of water molecules coordinating atom X of the amino acid in the i - th Umbrella Sampling window, and N_O^{bulk} is the corresponding coordination number when the amino acid is at the center of the water slab. Coordination numbers were computed by numerical integration of $g_{XO}(r)$ up to the first minimum (r_{fm}) , according to $N_O^i = 4\pi\rho \int_0^{r_{fm}} g_{XO}(r)r^2 dr$.

The dehydration profile of Gly, Lys, Pro and His was determined using eqn 2 for each of the 36 Umbrella Sampling windows, taking as reference (X) the carbonyl oxygen atom of the backbone. The results shown in Fig. 13 indicate that, except for charged His, all the amino acids exhibit a similar dehydration pattern as they are inserted into the bilayer, *i.e.*, each amino acid looses a total of 10-12 water molecules after insertion. In contrast, charged His looses much less hydration water (only 3 water molecules on the average) than the corresponding neutral tautomers. The His⁺ ion seems to be₄₉₂
effectively shielded by a tightly bound layer of hydration wa-493
ter, which could explain the similarity between the free energy₄₉₄
profiles of His⁺ and His reported in Fig. 12.



Fig. 13 Dehydration (Δ_{HN}) of the amino acids as they are inserted ⁵¹¹ into the lipid bilayer (see text for further explanation). Vertical lines ⁵¹² divide the system into 4 regions as in Fig. 7. Upper panel: Δ_{HN} for ⁵¹³ Gly (\bigcirc), Lys⁺ (\lhd), Pro (\bigtriangledown), His⁺ (\triangle), N^{δ 1}-H (\square), and N^{ϵ 2}-H (\diamondsuit), ⁵¹⁴ measured from the carbonyl oxygen atom of the backbone. Bottom panel: Δ_{HN} for the analog (\blacklozenge) and whole molecule model (\bigcirc) of Lys⁺ measured from the nitrogen atom of the lateral chain.

In the case of Gly, Figs. 9 and 13 provide a complementary 468 view of the change in the bonding pattern as the amino acid 469 penetrates into the membrane. Naturally the overall decrease 470 in the number of hydrogen bonds observed in Fig. 9 (a), is 471 concomitant with the decrease in the number of hydrating wa-472 ter molecules shown in the upper panel Fig. 13. In particular, 473 when Gly reaches the center of the bilayer (see Fig. 7, of the 474 Supporting Information) it losses 11, but retain (on average) 475 \sim 2-3, water molecules; among the retained water molecules, 476 only one form hydrogen bond with the Gly (see Fig. 9 a). 477 Simultaneously, Gly forms one hydrogen-bond with the phos-478 phate group of a lipid molecule (see Figs. 9 b and 10 A). 479

Finally, the bottom panel of Fig. 13 shows $\Delta_{HN}(z)$ for both the whole and the analog molecule model of Lys⁺, computed from the nitrogen atom of the lateral chain ($X = N_{chain}$ in eqn⁵¹⁶ 1). Clearly, the two models lead to significant quantitative differences in the number of solvating water molecules as the⁵¹⁷ amino acid moves towards the center of the bilayer. ⁵¹⁸

486 3.5 Unbiased Simulations of Amino Acids Exclusion 520 from the Bilayer Centre 521 522

In order to test whether the amino-acids could remain trapped⁵²³
in a metastable state when reaching the centre of the bilayer⁵²⁴
(potential local minima not captured in the free-energy pro-⁵²⁵
files), unbiased MD trajectories were initiated from the top of⁵²⁶

the free-energy barrier. In all cases simulations were started from configurations that had evolved under Umbrella Sampling for 100ns. Fig. 14 shows the time evolution of the distance, along the bilayer normal, between the center of mass (COM) of the membrane and the COM of the amino acid, after removing the harmonic restraint. The results are presented in the following order: (a) Gly, (b) N^{ε^2} -H, (c) N^{δ^1} -H, (d) His⁺, (e) Lys⁺, (f) Lys⁺-analog and (g) Pro. In all the seven cases the amino acid spontaneously leaves the membrane core, and migrate towards the bilayer surface. This occurs within a time scale of a few tens of nanoseconds. Gly, His (in the three forms), Lys⁺ and the Lys⁺-analog end up exploring the minimum of the free energy profiles reported in sections 3.2 and 3.3. In the special case of Pro, where a weak interaction with the membrane was found, it can be appreciated that the amino acid leaves the bilayer core, moves freely into the solvent, and finally gets in contact with the bilayer surface. The horizontal dashed lines in Fig. 14 represent the average distance to the bilayer centre, once the time series has stabilised. These values are collected in Table 2 (column d_{τ}^{a}) and compared with the position of the minima of the corresponding free-energy profiles (column d_z^b). Clearly, within a timescale of at most 50ns all amino acids reach the thermodynamic equilibrium position.

Table 2 Distance between the center of mass of DPPC and the corresponding amino acids. d_z^a corresponds to the average value obtained from the unrestrained molecular dynamics simulations (dashed lines in Fig. 14). d_z^b corresponds to the minimum of the free energy profiles reported in Sections 3.2 and 3.3.

Amino Acid	d_z^a (nm)	d_z^b (nm)
Gly	1.7 ± 0.2	~ 1.7
N^{ϵ_2} -H	2.1 ± 0.3	~ 1.9
$N^{\delta 1}$ -H	1.7 ± 0.3	~ 1.8
His ⁺	1.5 ± 0.2	~ 1.5
Lys ⁺	1.5 ± 0.2	~ 1.6
Lys ⁺ -analog	1.1 ± 0.2	~ 1.1
Pro	1.6 ± 0.2	~ 1.5

4 Conclusions

Molecular Dynamics simulations and FTIR experiments were used to investigate the interaction of a selected set of amino acids (Gly, Lys, Pro and three forms of His) with a dipalmitoylphosphatidylcholine (DPPC) bilayer. All the amino acids were considered to be in the zwitterionic form. Free energy profiles for the insertion of the amino acids into the membrane were computed by Umbrella Sampling, using as reaction coordinate the z-distance between the center of the bilayer and the amino acid. Given that Gly bears no side chain, it was taken as a reference system to investigate the backbone and side chain

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Fig. 14 Distance, between the center of mass of DPPC and the corresponding amino acids, along the normal to the bilayer plane. (a) Gy, (b) N^{ε_2} -H, (c) N^{δ_1} -H, (d) His⁺, (e) Lys⁺, (f) Lys⁺-analog and (g) Pro. The horizontal dashed lines, represents the average value of such distance after discarding the transient.

contributions to the free energy cost for transferring amino
acids from the aqueous phase to the surface and bulk of the
lipid membrane. Both simulations and experiments showed
that Gly adsorbs spontaneously on the surface of DPPC, forming distinguishable hydrogen-bonds with the lipids' phosphate
groups.

The analysis of free energy profiles for the insertion of Lys⁺, computed with a whole molecule model of the amino acid and the commonly used analog molecule approach, showed that Lys⁺ adsorbs strongly on DPPC and its insertion into the bilayer incurs a high energy penalty. More importantly, the comparison between the two PMFs for Lys⁺ and the PMF for Gly (backbone analog) demonstrated that the waterlipid transfer free energy of Lys⁺ can not be decomposed into additive side chain and backbone contributions. This puts a note of caution on the use of analog molecules when computing the transfer energy of peptides and flexible portions of proteins, such as statistical coil fragments or loop regions, which involve ~50% of all residues in proteins²¹.

Finally, Pro and His exhibit their own peculiarities and were discussed separately. Due to its chemical structure, Pro does not admit an analog molecule representation. Our calculations showed that this imino acid only exhibits unfavorable interactions with DPPC. At the same time the free energy profiles for the three forms of histidine resulted to be all very similar, although the charged form interacts mostly with the carbonyl groups of DPPC, while the tautomers do with the phosphate groups. Also, when entering the bilayer, the charged form of His preserves a significantly larger amount of hydration water than the two neutral tautomers.

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