NEW VERSION

How coagulation zone size is underestimated in computer modeling of RF ablation by ignoring the cooling phase just after RF power is switched off

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SUMMARY

All the numerical models developed for radiofrequency (RF) ablation so far have ignored the possible effect of the cooling phase (just after RF power is switched off) on the dimensions of the coagulation zone. Our objective was thus to quantify the differences in the minor radius of the coagulation zone computed by including and ignoring the cooling phase. We built models of RF tumor ablation with two needle-like electrodes: a dry electrode (5 mm long and 17G in diameter) with a constant temperature protocol (70°C) and a cooled electrode (30 mm long and 17G in diameter) with a protocol of impedance control. We observed that the computed coagulation zone dimensions were always underestimated when the cooling phase was ignored. The mean values of the differences computed along the electrode axis were always lower than 0.15 mm for the dry electrode and 1.5 mm for the cooled electrode, which implied a value lower than 5% of the minor radius of the coagulation zone (which was 3 mm for the dry electrode, and 30 mm for the cooled electrode). The underestimation was found to be dependent on the tissue characteristics: being more marked for higher values of specific heat and blood perfusion and less marked for higher values of thermal conductivity.

KEY WORDS: ablation, radiofrequency heating, tumor ablation

1. INTRODUCTION

Radiofrequency (RF) ablation is a high-temperature focal treatment aimed at destroying irreversibly localized zones of biological tissue. Electrical current is delivered to the tissue through metal electrodes to produce a local temperature increase and the consequent creation of a thermal coagulation zone. To date, numerous theoretical models have been developed to study different issues involved in the RF ablation procedure and most of these have been based on numerical methods [1]. One of the most important functions of these models is to predict the size of the coagulation zones created. To do this, computer simulations are conducted to obtain the evolution of tissue temperature and then to estimate the thermal damage contour by using functions which take the temperature-time relationship into account [2].

To the best of our knowledge, all the computer models so far developed to predict the size of the coagulation zone created by radiofrequency energy only take the RF application phase into account when estimating the accumulated thermal damage. However, it is known that tissue temperature does not drastically decrease to a baseline, but that there is a thermal latency, which implies that temperatures at deep locations increase slightly after RF power has been switched off [3]. Furthermore, after switching off the RF power the temperature decrease is especially slow in deep tissue, where powerful heat sink mechanisms (such as the cooling inside an RF applicator during RF tumor ablation), producing a "tail" in the temperature evolution, which could contribute to the accumulated thermal damage and hence to increasing the size of the coagulation zone. Our goal was to quantify the underestimation of the coagulation zone dimensions in computer modeling studies by ignoring the cooling phase after RF power has been switched off. Although the models

were not based on a specific tissue, the tissue properties varied within a broad range taken from literature data. We built RF ablation models for both dry and cooled needle-like electrodes and quantified the changes in the predicted coagulation zone dimensions by computer models when the cooling phase is considered in the thermal damage calculations and also identified the factors that contribute to these changes.

2. METHODS

2.1. Physical situation modeled

The physical situation modeled was a monopolar RF ablation conducted with needle-type electrodes, which are typically used in tumor ablation, assuming that the electrode was completely inserted in hepatic tissue (Figure 1A). The problem presents axial symmetry and hence a two-dimensional analysis was possible. Two types of needle-type electrodes were evaluated: dry (Fig. 1B) and internally cooled (Fig. 1C). In the former, the applied voltage was automatically modulated in order to keep the electrode temperature more or less constant. This was done with a proportional-integral (PI) control (see Section 2.4). These electrode and application modes are typically employed in bone tumor ablation [4]. In the second, an impedance control based on voltage pulses was implemented (see Section 2.5). These electrode type and application modes are typically employed in hepatic tumor ablation [5].

The dispersive electrode was modeled as an electrical condition on boundaries at a distance from the active electrode. A sensitivity analysis was used to check that tissue dimensions were large enough to keep tissue surfaces at boundary tissue temperature and voltage. The value of the maximum temperature (T_{max}) reached in the tissue at 60 s of RF

ablation was a control parameter in these analyses. The model dimensions were considered appropriate when there was a difference of less than 0.5% in T_{max} between the simulations.

2.2. Governing equations and model characteristics

All the models were based on a coupled electric-thermal problem, which was solved numerically on COMSOL Multiphysics software (COMSOL, Burlington MA, USA). The governing equation for the thermal problem was the bioheat equation modified by the enthalpy method to take the vaporization phenomenon into account [6]:

$$\frac{\partial(\rho h)}{\partial t} = \nabla \cdot (k\nabla T) + q + Q_p + Q_{met}$$
(1)

where ρ is tissue density, *h* enthalpy, *k* thermal conductive, *T* temperature and *t* time. *q* refers to the heat source from RF power, Q_p refers to the blood perfusion heat and Q_{met} refers to the metabolic heat which is negligible in RF ablation. For biological tissues enthalpy is related to tissue temperature by the following expression [6]:

$$\frac{\partial(\rho h)}{\partial t} = \frac{\partial T}{\partial t} \cdot \begin{cases} \rho_l c_l & 0 < T \le 99^{\circ} C \\ h_{fg} C & 99 < T \le 100^{\circ} C \\ \rho_g c_g & T > 100^{\circ} C \end{cases}$$
(2)

where ρ_i and c_i are density and specific heat of tissue respectively at temperatures below 100°C (*i*=*l*) and at temperatures above 100°C (*i*=*g*), h_{fg} is the product of water latent heat of vaporization and water density at 100°C, and *C* is water content inside the tissue (68%) [7].

We simulated an in vivo situation in which Q_p was computed from:

$$Q_{p} = \beta \rho_{b} c_{b} \omega_{b} (T_{b} - T)$$
(3)

where ρ_b is density of blood, c_b specific heat of blood, T_b blood temperature (37°C), ω_b blood perfusion coefficient ($\omega_b = 0.0155 \text{ s}^{-1}$ [8]) and β is a coefficient which took the values of 0 and 1, according to the value of the thermal damage: $\beta=0$ for $\Omega \ge 4.6$, and $\beta=1$ for Ω < 4.6. The parameter Ω is assessed by the Arrhenius damage model [9], which associates temperature with exposure time using a first-order kinetics relationship:

$$\Omega(t) = \int_{0}^{t} A e^{\frac{\Delta E}{RT(\tau)}} d\tau$$
(4)

where *R* is the universal gas constant, *A* $(7.39 \times 10^{39} \text{ s}^{-1})$ is a frequency factor and ΔE $(2.577 \times 10^5 \text{ J/mol})$ is the activation energy for the irreversible damage reaction [9,10]. We assessed the thermal damage by the D99 isoline [11], which corresponds to $\Omega = 4.6$ (99% probability of cell death). The location of the D99 isoline defined the size of the coagulation zone. Although some previous studies have used an isotherm to compute lesion size, in the context of our work it was impossible to use the isotherm to compute the damage (as in [12]) because it does not reproduce the accumulated damage when the RF power is turned off.

The heat source from RF power q (Joule losses) was given by $q=\sigma \cdot |\mathbf{E}|^2$, where \mathbf{E} is the electric field which was obtained from the electrical problem and σ was the electrical conductivity. The equation $\nabla \cdot \sigma \nabla V=0$ was the governing equation for the electrical problem, V being the voltage. The electric field was calculated by means of $\mathbf{E}=-\nabla V$. A quasi-static approach was used since the resistive current is much lower than the displacement current at RF \approx 500 kHz) [13]. Electrical conductivity σ is a temperature dependent piecewise function:

$$\sigma(T) = \begin{cases} \sigma(37)e^{0.015(T-37)} & 0 \le T < 99^{\circ}C \\ \sigma(37) 2.5345 & 99^{\circ}C \le T \le 100^{\circ}C \\ \sigma(37) 2.5345 - 0.50183(T-100) & 100^{\circ}C < T < 105^{\circ}C \\ \sigma(37) 2.5345 \times 10^{-2} & T > 105^{\circ}C \end{cases}$$
(5)

where σ (37) represents the electrical conductivity assessed at 37°C (see Table 1). Both initial and boundary temperatures were 37°C. Tissue boundary was set at electrical insulation, excepting the bottom and top surfaces which were set at V=0, since they represented the dispersive electrode. The delivering energy protocol depended on the electrode type (see Sections 2.4 and 2.5).

The model mesh was heterogeneous, with a finer mesh size at the electrode-tissue interface, where the highest electrical and thermal gradients were expected. All the mesh elements used were linear and triangular. An adaptive scheme was used for the time step, since we let the time-stepping method chose time steps freely. The size of the finer mesh and the optimal time-step was estimated by a sensitivity analysis on the same lines as that used for the outer geometry dimensions.

2.3. Tissue characteristics

The characteristics of materials used in the model are shown in Table 1, including the metallic electrode and the plastic partially covering the electrode [7,14]. We also assessed the effect of changing the tissue characteristics expected to affect the temperature distributions (and hence the size of the coagulation zone) during $t_{POST-RF}$: thermal conductivity, specific heat, and blood perfusion. The values were varied in the range reported in the literature [14]. Specifically, we selected three values for each parameter:

maximum, minimum, and intermediate. The minimum values for thermal conductivity and specific heat were those of low water content tissue (such as fat, bone marrow, and cortical bone), while the maximum values for thermal conductivity and specific heat were those of high water content tissue (e.g. vitreus humor). We also considered an intermediate value which was that of liver. Since it is known that that the blood perfusion value can be markedly affected by the target location, affections, and tumors, we varied the value according to literature data in the case of liver tissue [15,16]. Twenty-seven simulations were conducted with all the combinations of minimum, intermediate, and maximum values of thermal conductivity *k* (0.2, 0.52, and 0.6 W/m·K, respectively), specific heat *c* (1313, 3540, and 4000 J/kg·K, respectively), and blood perfusion ω_b (0.009, 0.0155, and 0.0285 s⁻¹, respectively).

2.4. Assessment of the impact of post-RF time

The impact of considering the cooling phase after RF power is switched-off ($t_{POST-RF}$) was studied by comparing the coagulation zone dimensions at this point in time (t_{ON-RF}) with those taken after a longer duration ($t_{ON-RF} + t_{POST-RF}$), which was the time until tissue temperature returned to baseline ($\approx 37^{\circ}$ C), i.e. the upper bound limit *t* in the integral of the Eq. (4) is really that of either the active ablation time (t_{ON-RF}) or total time ($t_{ON-RF} + t_{POST-}$ RF). All the simulations showed that at 400 s after RF application the difference between the highest temperature and baseline was around 0.004°C, so that we considered $t_{POST-RF} = 800$ s as reasonable. Coagulation zone dimensions were compared by the Arrhenius model. The main idea was to calculate the difference between the electrode-D99 distances computed at t_{ON-RF} and $t_{ON-RF} + t_{POST-RF}$. The measurement was taken from the region of the cylindrical part of the active electrode, without considering contact with the insulated part or the distal tip of the active electrode (see Fig. 2a). The difference in the damage was computed by subtracting the coordinates r of the contours ($r_{post} - r_{on}$) along the z axis (Fig. 2c). Both r_{post} and r_{on} are really transverse diameters of the coagulation zone created with a needle-like electrode. They are also known as minor diameters, as opposed to the major diameters, which are parallel to the electrode axis. The mean value of the difference between r_{post} and r_{on} was computed. The graph in Fig. 2c represents the isolines shown in Fig. 2b.

2.5. *Needle-type dry electrode and constant temperature protocol*

We first modeled a non-cooled (dry) electrode working in a constant temperature mode with a target temperature of 70°C using a PI controller similar to the one described in [17]. Briefly, a dynamical system was considered in which the input variable was the voltage applied to the electrode, and the output variable was the temperature measured at the tip of the electrode. There are different ways of controlling tip temperature, the most common strategy being proportional-integral-derivative (PID) control. Here, we tuned a PI controller that reached the target temperature within 20 seconds without overshooting. Using intermediate values of the physical properties of the tissue, we obtained $K_p = 1.15^{\circ}$ C/V y K_i = 0.1°C/V·s. Ablation duration (t_{ON-RF}) was 240 s, electrode length 5 mm and electrode diameter 17G (1.473 mm), parameters typically employed in RF ablation of musculoskeletal tumors.

2.6. Cooled needle-type electrode and impedance control protocol

Long internally cooled electrodes are used for RF ablation of tumors in soft organs such as the liver. The most typical of these electrodes is the Cool-Tip device (Covidien, Boulder, CO, USA), with a 17G electrode (1.473 mm diameter) and 3 cm in length. The cooling effect of the liquid circulating inside the electrode was modeled using a thermal convection coefficient h_r with a value of 3,127 W/K·m² and a coolant temperature of 5°C, following Newton's Law of Cooling. The value of h_r was calculated by considering the electrode length (3 cm) and a flow rate of 45 mL/min through an area equivalent to half the cross section of the inner diameter of the electrode.

An impedance-controlled pulsing protocol was modeled for a 240 s ablation with a current constant of 1,500 mA, which is aimed at increasing coagulation zone size, since periods of low current deposition alternate with higher peak current, allowing the tissue near the electrode to cool while the deeper tissue is heated [18]. RF power was switched off once the impedance exceeded a threshold, or roll-off, which is usually 20 Ω higher than the initial value, and was switched on again after 15 s [19].

3. RESULTS

Figure 3 shows the mean value of the difference in damage computed along the *z* axis for the needle-type dry electrode and constant temperature protocol when the tissue characteristics were changed. Overall, the most important finding was that the differences between considering and ignoring the cooling phase in the computer simulations were always lower than 0.15 mm (in terms of coagulation zone radius), which represented approximately 5% of the radius of the coagulation zone (around 3 mm). However, in most cases the difference was even lower than 0.1 mm ($\approx 3.3\%$).

In the simulations of a needle-type cooled electrode with an impedance control protocol, there were around 6 roll offs throughout the 240 s of ablation. Figure 4 shows the mean value of the difference in damage computed along the *z* axis for the needle-type cooled electrode with an impedance control protocol when tissue characteristics were varied. In this case, the differences between considering and ignoring the cooling phase in the computer simulations were always lower than 1.5 mm (in terms of radius of coagulation zone), which represented around 5% of the radius of the coagulation zone (around 30 mm). However, in most cases the difference was less than 1 mm (\approx 3.3%).

Regarding the effect of changing the tissue characteristics, the findings were qualitatively similar for both electrode designs and protocols. Higher values of thermal conductivity involved slightly lower $r_{post} - r_{on}$ differences (see the trend lines in Figs. 3a and 4a). However, the changes in blood perfusion and specific heat showed a more marked effect. The $r_{post} - r_{on}$ differences were considerably lower for higher values of blood perfusion, and increased for higher values of specific heat.

Finally, the effect of ablation time was also assessed by analyzing the $r_{\text{post}} - r_{\text{on}}$ difference throughout the ablation period (each 60 s). We did not find any important differences or any perceptible trend over time.

4. DISCUSSION

4.1. Review of the results

As the numerical modeling technique applied to ablative therapies based on hightemperature is intended to accurately predict thermal coagulation zones, the models should hence be built on a solid foundation of realism. One of the issues that have been traditionally ignored in numerical modeling is the role played by the *remaining* heat stored in the tissue once RF power is switched off. All the computer modeling studies to date have considered only a simulation time corresponding with the duration of the RF power application, and therefore quantified the coagulation zone dimensions at the end of this period, which clearly ignores the possible effect of the *remaining* heat to slightly increase the final size of the coagulation zone.

Our goal was to compare the coagulation zone dimensions, especially the minor radius, created by considering and ignoring this *remaining heat*. We focused on the destruction of tumors, which is one of the most important applications in RF ablation. In this context, we considered two electrode designs (dry and cooled), each one with their particular protocol to deliver RF power (constant temperature and impedance control, respectively). In both cases we observed that the impact of ignoring the cooling phase in the computer simulations meant that the minor radius of the coagulation zone was underestimated by less than 6.6% and 12% for dry and cooled electrodes, respectively. The clinical impact of this finding depends on the specific objective of each modeling study. For instance, it could be considered to be negligible in the case of comparative studies in which the absolute values of coagulation zone dimensions are not important. In contrast, a difference of 5% could be important when the numerical model is employed to accurately predict the coagulation zone dimensions.

This underestimation was dependent on the tissue characteristics in a way that can be explained in physical terms. The cases in which the underestimation in the coagulation zone dimension is higher corresponded with those in which the *remaining heat* is evacuated more slowly. This occurs when: 1) the tissue has more capacity to store heat, which is

equivalent to having a high value of specific heat and/or of density (issue not evaluated in this study, since all biological tissues have a similar density, around $1,000 \text{ kg/m}^3$); 2) heat cannot be easily conducted through the tissue, which is equivalent to having a low value of thermal conductivity; and 3) the heat evacuation mechanisms are poor, which is equivalent to having a low blood perfusion value.

4.2. Limitations of the study

This study has certain limitations which need to be emphasized: firstly, the differences in coagulation zone dimensions between considering and ignoring the cooling phase were exclusively evaluated in terms of the thermal coagulative necrosis created by the direct action of heat, i.e. using the Arrhenius damage model, which associates temperature with exposure time (Eq. 4). In a real situation, the coagulation zone finally created would include a thin rim comprised of cells which would die by apoptosis (maybe caused by thermal shock) rather than by thermal coagulative necrosis. This issue is outside the scope of this study, which only focused on the role of the *remaining heat* in cell death.

Secondly, we only considered the case of RF tumor ablation with a needle-like electrode, although other settings and electrode designs are also possible in RF ablation numerical modeling. For instance, in RF ablation of cardiac arrhythmias, a semispherical electrode embedded at the tip of an intravascular catheter is placed on the endocardium. In this case, dry and cooled electrodes can also be employed, along with different protocols for delivering RF power (constant temperature and constant power). Even though RF cardiac ablation was not considered in this study, we do not think that values higher than 5% can be observed in the underestimation of the coagulation zone dimensions. In our

opinion, *remaining heat* could be rapidly evacuated towards the circulating blood through the electrode body and through the endocardium-blood interface, which would reduce the difference between r_{post} and r_{on} .

4.3. Clinical impact of the findings

Apart from a scientific interest in the effect of the cooling phase on the computed damage produced by thermal therapy, it is also of interest to consider the real clinical impact of an estimation error of 5% (equivalent to 3 mm in the case of a 3-mm long RF electrode). This is important in the context of computer-aided planning and simulation tools and image-guided ablative procedures. In this respect, although the current electromagnetic navigation systems offer tracking errors of up to 5 mm [20], in the near future imaging systems are expected to offer an improved resolution of even less than 300 microns [21]. For instance, in the context of RF ablation of the atrial wall, the mathematical processing of a finite element model based on images from computed tomographic angiography has already made it possible to estimate the atrial wall thickness to within less than 0.2 mm [22]. It is precisely in the field of RF heating of small volume targets (such as the atrial wall) where an underestimation error of 5% could have serious consequences. In fact, it is in this field, in which optical techniques offer extremely high spatial resolution, that they are being proposed to map thermal lesions [23].

5. CONCLUSIONS

Ignoring the cooling phase that occurs just after RF power has been switched off always involves underestimating the computed coagulation zone dimensions. This underestimation

was always lower than 5% of the coagulation zone radius for both dry and cooled electrodes and was dependent on tissue characteristics: it was more marked for higher values of specific heat and blood perfusion, and less marked for higher values of thermal conductivity.

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Material	ρ (kg/m ³)	c (J/kg·K)	$k (W/m \cdot K)$	σ (S/m)
		(1) (2)	(2)	
Tissue	1079 (1)	3540 ⁽¹⁾ [1313–4000] ⁽³⁾	0.52 [0.2–0.6] (3)	0.148
	(2)	(2)		
	370 (2)	2156 ⁽²⁾		
Electrode	8×10 ³	480	15	7.4×10°
Plastic	70	1045	0.026	1×10-5
Blood	1000	3639		

Table 1. Characteristics of the materials used in the theoretical model [7,8].

 ρ , density; *c*, specific heat; *k*, thermal conductivity; σ , electric conductivity;

⁽¹⁾At temperature below 100°C. ⁽²⁾At temperature above 100°C.

⁽³⁾ Minimum and maximum values for the sensitivity analysis (see section 2.3 for more details).

⁽⁴⁾ Assessed at 37°C.



Figure 1 A: Geometry of the model used in the first set of simulations (out of scale, dimensions in mm).The outer dimensions (60 mm wide and 150 mm long) were calculated by means of a sensitivity analysis. B: Detail of dry electrode geometry. C: Detail of internally cooled electrode geometry.



Figure 2 Procedure to compare the coagulation zones created with and without considering $t_{POST-RF}$ in the simulations. (a) Scheme of the variables under study. (b) Isolines D99 at t_{ON-RF} (dashed line) and at $t_{ON-RF} + t_{POST-RF}$ (continuous line). (c) Difference of damage $(r_{post} - r_{on})$ along the electrode (coordinate z). The dashed line shows the mean value, which is considered in the subsequent analyses.



Figure 3 Mean values of the difference of thermal damage $(r_{post}-r_{on})$ along the electrode axis (coordinate *z*) between considering (r_{post}) and ignoring (r_{on}) the cooling phase in the simulations. The results correspond to the case of needle-type dry electrode and constant temperature protocol, and for different values of the tissue characteristics: thermal conductivity (k), blood perfusion (ω) and specific heat (c). The value shown in Table 1 was considered if not otherwise indicated.



Figure 4 Mean values of the difference of thermal damage $(r_{post}-r_{on})$ along the electrode axis (coordinate *z*) between considering (r_{post}) and ignoring (r_{on}) the cooling phase in the simulations. The results correspond to the case of needle-type cooled electrode and impedance control protocol, and for different values of the tissue characteristics: thermal conductivity (k), blood perfusion (ω) and specific heat (c). The value shown in Table 1 was considered if not otherwise indicated.