

# Kinetic Study of the $\text{CCl}_2$ Radical Recombination Reaction by Laser-Induced Fluorescence Technique

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**ABSTRACT:** An experimental setup that coupled IR multiple-photon dissociation (IRMPD) and laser-induced fluorescence (LIF) techniques was implemented to study the kinetics of the recombination reaction of dichlorocarbene radicals,  $\text{CCl}_2$ , in an Ar bath. The  $\text{CCl}_2$  radicals were generated by IRMPD of  $\text{CDCl}_3$ . The time dependence of the  $\text{CCl}_2$  radicals' concentration in the presence of Ar was determined by LIF. The experimental conditions achieved allowed us to associate the decrease in the concentration of radicals to the self-recombination reaction to form  $\text{C}_2\text{Cl}_4$ . The rate constant for this reaction was determined in both the falloff and the high-pressure regimes at room temperature. The values obtained were  $k_0 = (2.23 \pm 0.89) \times 10^{-29} \text{ cm}^6 \text{ molecules}^{-2} \text{ s}^{-1}$  and  $k_\infty = (6.73 \pm 0.23) \times 10^{-13} \text{ cm}^3 \text{ molecules}^{-1} \text{ s}^{-1}$ , respectively. © 2013 Wiley Periodicals, Inc. *Int J Chem Kinet* 45: 306–313, 2013

## INTRODUCTION

Hazardous waste treatment is an issue of major concern. Among the various types of hazardous waste, chlorinated hydrocarbons are produced in many industries. Chlorinated compounds in commercial use include pesticides, pharmaceuticals, disinfectants, and consumer products such as plastics. Chlorinated hydro-

carbons are used as intermediates in the synthesis of other chemicals as well as to produce silicon for electronic chips and epoxy resins. Chlorinated methanes, ethanes, and ethenes solvents are used for degreasing and dry cleaning. Polychlorinated biphenyls and triphenyls are used as dielectric and coolant fluids in transformers, capacitors, and electric motors. Thus, there is a great need to destroy huge quantities of chlorinated hydrocarbons. Two processes are mainly employed for the treatment of these types of waste: incineration and pyrolysis. Alternative chemical degradation techniques such as UV/ozonation, semiconductor photocatalysis, and catalytic destruction of chlorocarbons with steam are currently under research [1–6]. The pyrolysis and

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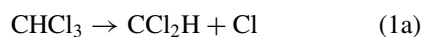
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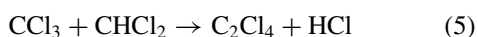
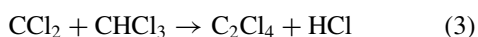
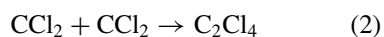
combustion of chlorinated hydrocarbons also leads to the formation of chlorinated aromatic species [7]. It is therefore of great interest to understand the elementary reaction mechanisms involved in these processes as well as to determine their rate constants in view to optimize chlorinated hydrocarbon destruction.

Particularly, one of the halogenated hydrocarbons widely used in industry is chloroform. Large amounts of this compound are produced annually since its main use is as a precursor to Teflon. It is also used as a solvent in the pharmaceutical, dye, pesticide, and dry cleaning industries. Pyrolysis of chloroform under a broad range of conditions has been extensively studied in the past few decades [7–21]. Reaction pathways to give account of the products determined at different stages of the reaction have been described. Earlier studies proposed C–Cl fission (1a) as the main primary dissociation step [8–10]. Later investigations have concluded that the three-center HCl elimination reaction (1b) is the dominant initiation step [11,14,15,17–21]:



More specifically, the infrared multiple-photon dissociation (IRMPD) of CDCl<sub>3</sub> with a TEA CO<sub>2</sub> laser occurs only through channel (1b) with no further fragmentation of the CCl<sub>2</sub> radical [11,22].

In the above-mentioned chloroform pyrolysis studies, C<sub>2</sub>Cl<sub>4</sub> has been the major intermediate chlorocarbon product detected over a wide temperature range [9]. The main pathway proposed to C<sub>2</sub>Cl<sub>4</sub> formation is the combination reaction of two dichlorocarbene radicals, CCl<sub>2</sub> [15,18,20], (2). Other pathways include the insertion reaction of CCl<sub>2</sub> radical with CHCl<sub>3</sub> [15,20], (3), and the combination reactions of CCl<sub>3</sub> (formed in the Cl abstraction of H from CHCl<sub>3</sub>) with CCl<sub>2</sub> [15,18,20] and CHCl<sub>2</sub> [15], (4) and (5), respectively:



The recombination reaction of dichlorocarbene radicals to form C<sub>2</sub>Cl<sub>4</sub> has been scarcely studied. Won and Bozzelli [15] studied the pyrolysis of CHCl<sub>3</sub> in a reaction system consisting of 1% CHCl<sub>3</sub> in 760 Torr Ar in the temperature range of 800–1073 K. Gas chromatographic analysis was used to determine the reactant and products concentrations as a function of

temperature at different reaction times. They proposed a detailed kinetic reaction mechanism consisting of 31 species and 67 elementary reactions to describe the reactant loss and the product formation, up to C<sub>2</sub> compounds. The mechanism describes the overall reaction process quite well but overestimates the C<sub>2</sub>Cl<sub>4</sub> concentrations over the temperature range of 848–953 K. The high-pressure limit rate constant for the CCl<sub>2</sub> recombination reaction to form C<sub>2</sub>Cl<sub>4</sub> calculated using unimolecular quantum version of Rice–Ramsperger–Kassel (QRRK) theory was  $1.1 \times 10^{-12} \exp(2.5 \text{ kcal/mol}/RT) \text{ cm}^3 \text{ molecules}^{-1} \text{ s}^{-1}$ .

Kumaran et al. [18] studied the pyrolysis of 1% and 4% CHCl<sub>3</sub> in Kr between 1282 and 1878 K. They used the laser schlieren technique to measure CHCl<sub>3</sub> decomposition, and time-resolved Cl formation was followed by the Cl-atom resonance absorption spectroscopic technique. A kinetic reaction mechanism consisting of six elementary reactions was used to describe the secondary CCl<sub>2</sub> chemistry. They reported a termolecular rate constant for the CCl<sub>2</sub> recombination reaction to form C<sub>2</sub>Cl<sub>4</sub> in the falloff region derived with this mechanism of  $1.6 \times 10^{-32} \exp(5.9 \text{ kcal/mol}/RT) \text{ cm}^6 \text{ molecules}^{-2} \text{ s}^{-1}$ .

In 2003, Zhu and Bozzelli [20] analyzed previous literature results of the thermal conversion of CHCl<sub>3</sub> to CCl<sub>4</sub> in the presence of Cl<sub>2</sub> in the 846–908 K temperature range using a detailed mechanism that included 38 elementary reactions. The high-pressure limit rate constant for the association reaction of two CCl<sub>2</sub> radicals to form C<sub>2</sub>Cl<sub>4</sub>, which resulted from their QRRK-master equation calculations, was  $4.5 \times 10^{30} T^{-14.2} \exp(-9.2 \text{ kcal/mol}/RT) \text{ cm}^3 \text{ molecules}^{-1} \text{ s}^{-1}$ .

In the above-mentioned works, the rate constant for the recombination reaction of two CCl<sub>2</sub> radicals to form C<sub>2</sub>Cl<sub>4</sub> was calculated from complex mechanisms used to fit products distributions from CHCl<sub>3</sub> pyrolysis, employing Rice–Ramsperger–Kassel–Marcus (RRKM) and QRRK rate theories.

In this work, we have used laser-induced fluorescence (LIF) supported by Fourier transform IR spectrometry (FTIR) to study the CCl<sub>2</sub> association reaction to form C<sub>2</sub>Cl<sub>4</sub> in an Ar bath. The LIF technique employed has allowed us to considerably simplify the reaction mechanism and to determine the rate constant in a direct manner. The CCl<sub>2</sub> radicals were generated by IRMPD of CDCl<sub>3</sub> and then excited in the transition A<sup>1</sup>B<sub>1</sub>(0,4,0) ← X<sup>1</sup>A<sub>1</sub>(0,0,0). The time-dependent radical concentration was evaluated from the intensity of the fluorescence emitted by the excited species as they returned to the electronic ground state, which was detected for different delay times between the dissociation and the excitation pulses. The particular irradiation and LIF detection system's geometry

used ensured that the LIF signal detected is related to the fraction of molecules dissociated per pulse in a region of constant fluence [22]. Since the fluence value in the intersection zone of the observation and the dissociation volumes corresponded to the dissociation saturation value, every  $\text{CDCl}_3$  molecule in this region was dissociated through the three-center  $\text{DCI}$  elimination pathway (1b) [22]. Thus, the time dependence of the  $\text{CCl}_2$  radical concentration can be directly associated to the recombination reaction to form  $\text{C}_2\text{Cl}_4$  (2). The rate constant for this reaction was determined in both the falloff and the high-pressure regimes.

## EXPERIMENTAL

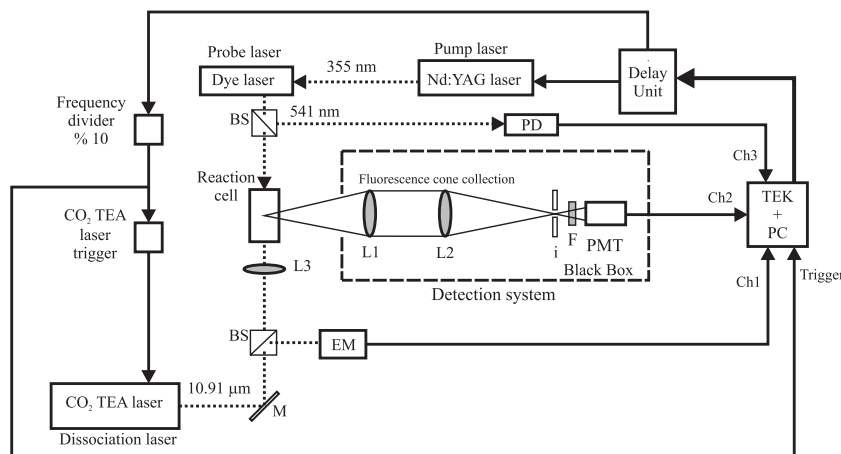
The time dependence of the  $\text{CCl}_2$  radical concentration generated in the IRMPD of  $\text{CDCl}_3$  was determined coupling IRMPD and LIF techniques. The experimental setup shown in Fig. 1 was described in detail elsewhere [22,23]. Briefly, the IRMPD of  $\text{CDCl}_3$  was implemented using a multimode, pulsed TEA  $\text{CO}_2$  laser with a repetition rate of 1 Hz tuned to the 10P(48) emission line. The laser radiation was focused into the center of the reaction cell with a 9.4 cm focal length ZnSe lens determining a beam's spot diameter at the focus and a Rayleigh range of 0.74 and 3.6 mm, respectively.

The radicals were electronically excited with a dye laser pumped by a frequency-tripled Nd:YAG laser pulsed at a repetition rate of 10 Hz. The dye laser consisted of an oscillatory stage followed by an amplification stage. The dye used was Coumarin 540 (Exciton), and the laser wavelength, controlled by an Avantes

spectrometer, Ava Spec 3648, was tuned to 541 nm resonant with the  $\text{CCl}_2$  electronic transition  $\text{A}^1\text{B}_1(0,4,0) \leftarrow \text{X}^1\text{A}_1(0,0,0)$ .

The laser beams entered the reaction cell collinearly, and the fluorescence signal was detected at  $90^\circ$  as shown in Fig. 1. A photomultiplier tube, RCA 1P28, was used for the detection of the fluorescence signal and a long-pass interference filter and an iris were disposed at the entrance of the photomultiplier tube. To increase the solid angle of detection, part of the fluorescence emitted by the  $\text{CCl}_2$  radicals was directed to the photomultiplier tube by a two-lens system (10 and 24 cm focal lengths, respectively). The iris aperture was fixed so as to obtain an observation volume coincident with the volume in which the probability of dissociation is equal to one.

The Nd:YAG and TEA  $\text{CO}_2$  lasers emissions were externally controlled with a Stanford Research Systems delay unit, DG 535. One channel of the delay unit was used to trigger the TEA  $\text{CO}_2$  laser. As shown in Fig. 1, the output of this channel was connected to the entrance of a frequency divider so as to obtain a repetition rate of 1 Hz. Another channel was used to trigger the Nd:YAG laser. The Nd:YAG laser trigger was delayed with respect to that of the TEA  $\text{CO}_2$  laser. The dye and TEA  $\text{CO}_2$  lasers' output energies, 200  $\mu\text{J}$  and 1 J, respectively, were monitored from the beam reflections on different beam-splitters and registered by two GenTec meters. In all the experiments, the TEA  $\text{CO}_2$  laser fluence at the beam waist was  $100 \text{ J cm}^{-2}$ . The fluorescence signals for different delay times as well as the pulse-to-pulse TEA  $\text{CO}_2$  and dye lasers' output energies were acquired with an oscilloscope, Tektronix DPO 7054.



**Figure 1** Experimental setup scheme. EM: energy meter, L1, L2: quartz lenses, L3: ZnSe lens, BS: beam-splitters, TEK: digital oscilloscope with incorporated PC, PMT: photomultiplier, F: filter, i: iris, PD: photodiode, M: mirror.

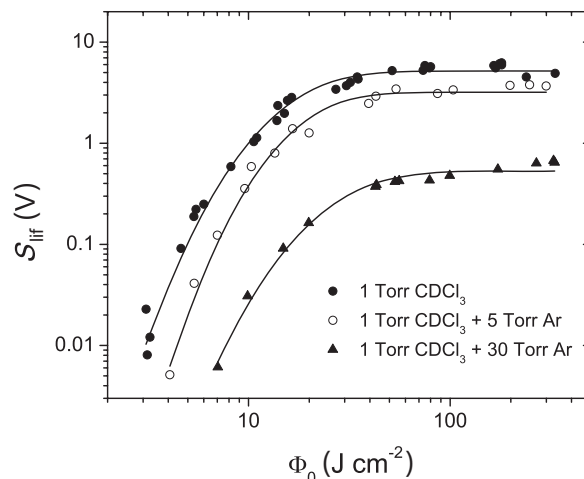
The delay times were varied in a geometric progression between 1  $\mu$ s and 10 ms so as to characterize processes that occur in different timescales with the same resolution. In addition, to eliminate history effects, the delay times of each series of experiments were chosen at random. A specific program was developed for the control of the whole experiment by a PC.

The samples of CDCl<sub>3</sub> (Merck; 99.98%) and Ar (Alphagas; 99.9995%) were prepared in a high vacuum system, consisting of a turbomolecular pump (Leybold, Turbovac) pumped by a rotary vane pump (Leybold, Trivac) and introduced in a cylindrical static Pyrex glass cell 3 cm in diameter and 10 cm long. The partial pressures of the reactant and products were controlled by FTIR spectrometry. In all experiments, the CDCl<sub>3</sub> pressure was fixed at 1 Torr and the Ar pressure was varied in the 0–30 Torr range. To consider that the CCl<sub>2</sub> concentration produced in each pulse was constant, the global dissociation was kept less than 10%.

## RESULTS

The LIF signal is characterized by a fast rise time due to the fluorescence emission produced in the electronic decay to the ground state followed by a slow decay given by the detection system response. The LIF signal peak amplitude,  $S_{\text{lif}}$ , was used as an estimator of the fluorescence intensity and, therefore, of the CCl<sub>2</sub> concentration.

To understand both the formation and the kinetic mechanisms of the CCl<sub>2</sub> radicals, we will first describe some results concerning IRMPD of CDCl<sub>3</sub> obtained in previous works [22,23] with the same experimental setup. We have implemented a LIF technique, which allowed monitoring the fluence dependence of the CCl<sub>2</sub> radicals generated in the IRMPD of CDCl<sub>3</sub> in pure samples [22] and in mixtures with Ar [23]. Figure 2 shows the results obtained for a sample of 1 Torr of pure CDCl<sub>3</sub> and in mixtures with 5 and 30 Torr of Ar. The three curves are characterized by two well-differentiated regimes. In the 1–40 J/cm<sup>2</sup> range, the LIF signal increases with increasing fluence and then it reaches a plateau corresponding to the saturation regime. The fluence value for which saturation is attained depends on the Ar pressure and ranges between 20 and 30 J cm<sup>-2</sup> for Ar pressures varying between 0 and 30 Torr. It can be observed that as the Ar pressure is increased, the LIF signal saturation value decreases due to collisional quenching of the excited CCl<sub>2</sub> radicals. Since on the one hand the dye laser is tuned resonantly with the electronic transition  $A^1B_1(0,4,0) \leftarrow X^1A_1(0,0,0)$  of the CCl<sub>2</sub> radicals and, on the other hand, at 100 J cm<sup>-2</sup> the fluorescence signal has reached a saturation regime for all Ar pressures, one can infer



**Figure 2** Fluence dependence of the LIF signal for mixtures of 1 Torr CDCl<sub>3</sub> with Ar.

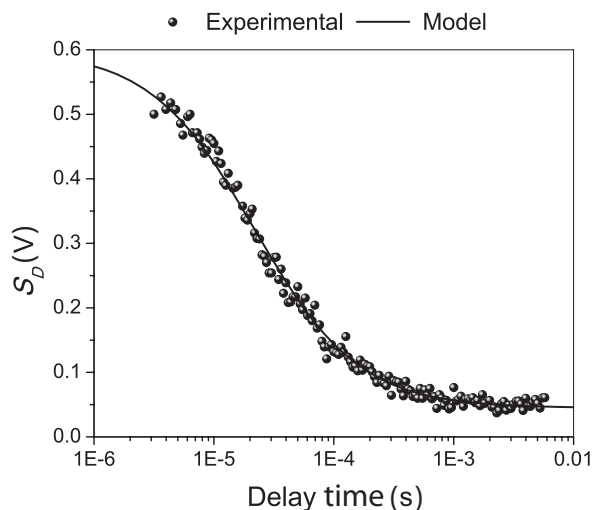
that every CDCl<sub>3</sub> molecule in the observation volume is dissociated through the three-center DCI elimination reaction (1b).

Therefore, during the acquisition time, the observation volume will only contain CCl<sub>2</sub> radicals, DCI, C<sub>2</sub>Cl<sub>4</sub>, and Ar so that the contribution of the insertion reaction of CCl<sub>2</sub> radical with CDCl<sub>3</sub> (3) to the CCl<sub>2</sub> radicals disappearance observed would be negligible. Contributions of the combination reactions of CCl<sub>3</sub> with CCl<sub>2</sub> (4) and CHCl<sub>2</sub> (5), respectively, can also be discarded since the CCl<sub>3</sub> radicals are mainly formed in the Cl abstraction of H from CHCl<sub>3</sub> [15,18,20], which does not take place in our experimental conditions. Thus, in our experimental conditions, the reaction mechanism following CDCl<sub>3</sub> IRMPD can be reduced to the occurrence of reaction (2).

In addition, we have used FTIR spectrometry to determine the final products in the irradiated samples. In accordance with the above considerations, these spectra showed that the major products obtained were DCI and C<sub>2</sub>Cl<sub>4</sub> and that CCl<sub>4</sub> was the only minor product. CCl<sub>4</sub> is assumed to be formed in secondary chemistry occurring out of the observation volume.

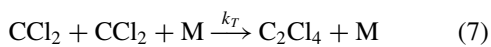
To study the CCl<sub>2</sub> radicals reactions, we have varied the delay between the dissociation and excitation lasers in the range 1  $\mu$ s–10 ms. A typical curve showing the dependence of the fluorescence signal maximum on the laser pulses' delay time for an Ar pressure of 25 Torr and a fluence of 100 J/cm<sup>2</sup> is shown in Fig. 3. Owing to the large time range used, the delay time is shown in a logarithmic scale.

Taking into account that the generation and subsequent reaction of the dichlorocarbene radicals in our



**Figure 3**  $S_D$  dependence on the delay time between the dissociation and excitation lasers' pulses. Symbols: experimental data, line: data calculated with Eq. (10).

experimental conditions can be described by



the time dependence of the  $\text{CCl}_2$  radicals after the photolysis of  $\text{CDCl}_3$  can be expressed by

$$\frac{d[\text{CCl}_2]}{dt} = -2k_T[\text{CCl}_2]^2 \quad (8)$$

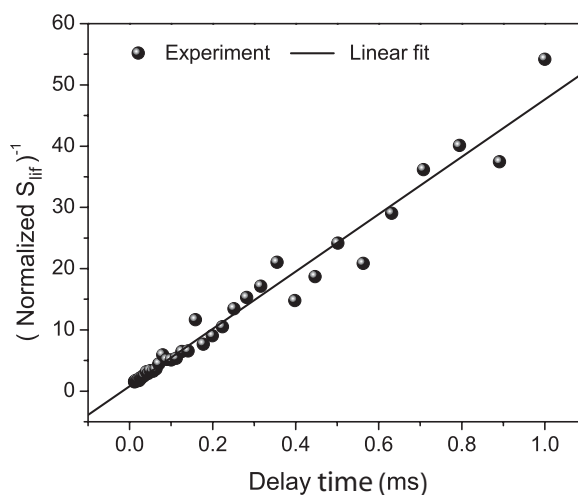
where  $k_T$  is the rate constant of the termolecular reaction. The time dependence of the  $\text{CCl}_2$  radicals concentration is obtained from the solution of Eq. (8):

$$[\text{CCl}_2](t) = \frac{[\text{CCl}_2]_0}{1 + 2k_T[\text{CCl}_2]_0 t} \quad (9)$$

Considering that the LIF signal is proportional to the radicals' concentration, the detected signal,  $S_D$ , will be given by

$$S_D(t) = S_{\text{lif}}(t) + S_b = \frac{q_i}{1 + \beta t} + S_b \quad (10)$$

where  $\beta = 2k_T[\text{CCl}_2]_0$ ,  $q_i$  is proportional to the initial concentration of radicals, and  $S_b$  is the background signal in the absence of  $\text{CCl}_2$  radicals. The coefficient  $q_i$  is a constant, which accounts for experimental factors such as the photomultiplier response, the fluorescence signal emitted in the observation solid angle, and the ratio of the number of radicals in the observation vol-



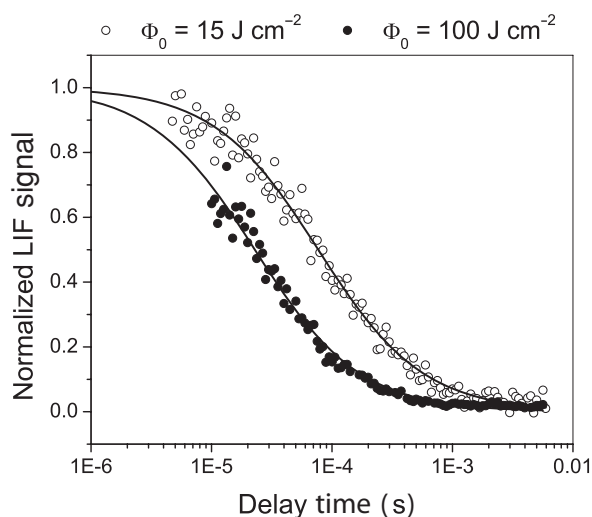
**Figure 4** Delay time dependence of the inverse of the normalized LIF signal for an Ar pressure of 6 Torr and a fluence of  $100 \text{ J cm}^{-2}$ .

ume to the total number of radicals excited by the dye laser. The inverse value of  $\beta$  represents the mean lifetime of the radicals in the observation volume,  $\tau$ , and depends on the rate constant of reaction and on the initial concentration of radicals. Equation (10) was used to fit the LIF signal maximum versus the delay time for the different Ar pressures with  $q_i$ ,  $\beta$ , and  $S_b$  as the adjusting parameters. A typical example is shown in Fig. 3.

Also, Fig. 4 shows a plot of the delay time dependence of the inverse of the normalized LIF signal for an Ar pressure of 6 Torr and fluence of  $100 \text{ J cm}^{-2}$ . The experimental data are well described by a linear fit showing that the radicals disappearance follows a second-order kinetic due to the recombination reaction (7).

The initial concentration of radicals is an essential magnitude that directly affects the recombination rate, and, as it can be observed in Fig. 2, it depends on the irradiation fluence. This fact led us to study the dependence of the fluorescence intensity estimator  $S_{\text{lif}}(t)$  of a mixture of 1 Torr  $\text{CDCl}_3$  with 5 Torr Ar on the lasers pulses' delay time for two irradiation fluences, 15 and  $100 \text{ J cm}^{-2}$ . The experimental results of the normalized fluorescence intensity (symbols) and the values calculated with the help of Eq. (10) are shown in Fig. 5.

As can be observed, the radicals' mean lifetime at  $100 \text{ J cm}^{-2}$  is shorter than that at  $15 \text{ J cm}^{-2}$ . On the other hand, from Fig. 2 it results that  $100 \text{ J cm}^{-2}$  exceeds the fluence value required for  $\text{CDCl}_3$  IRMPD saturation, whereas  $15 \text{ J cm}^{-2}$  lies below the saturation fluence level. The values of  $\beta$  for both fluences



**Figure 5** Delay time dependence of the normalized fluorescence intensity of a mixture of 1 Torr CDCl<sub>3</sub> and 5 Torr Ar for 15 and 100 J cm<sup>-2</sup>.

were obtained from the curves fitted with Eq. (10). The resulting mean lifetime values of the radicals were  $\tau(15) = (79.8 \pm 3.6) \mu\text{s}$  and  $\tau(100) = (22.9 \pm 1.8) \mu\text{s}$ , where the errors correspond to one standard deviation, and their ratio was

$$\tau(15)/\tau(100) = 3.5 \pm 0.3$$

Since the kinetic mechanism of the radicals disappearance is determined by reaction (7), this ratio must coincide with that of the radicals' concentration at 100 and 15 J cm<sup>-2</sup>. The LIF signal values obtained from Fig. 2 at these fluences were

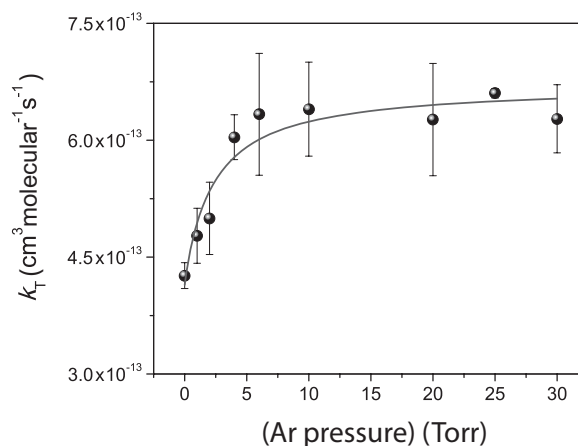
$$S_{\text{lif}}(15) = (1.1 \pm 0.2) V$$

$$S_{\text{lif}}(100) = (3.2 \pm 0.5) V$$

and therefore

$$S_{\text{lif}}(100)/S_{\text{lif}}(15) = (2.9 \pm 0.6)$$

Both results agree within experimental error supporting the statement that in our experimental conditions the CCl<sub>2</sub> radicals' disappearance is associated only to the recombination reaction to form C<sub>2</sub>Cl<sub>4</sub>, reaction (7). Therefore, for the experiments performed at 100 J cm<sup>-2</sup>, the time dependence of  $S_{\text{lif}}$  was fitted with Eq. (10) assuming an initial concentration of CCl<sub>2</sub> radicals of 1 Torr since every CDCl<sub>3</sub> molecule in the observation volume was dissociated [22]. The pressure dependence of the recombination rate constant  $k_T$  was obtained from the fits of  $S_{\text{lif}}(t)$  for the different Ar pressures.



**Figure 6** Ar pressure dependence of the rate constant of the CCl<sub>2</sub> radicals' recombination reaction. Symbols: experimental data, line: values calculated with Eq. (11).

The dependence of the CCl<sub>2</sub> radicals' recombination reaction rate constant,  $k_T$ , on Ar pressure is shown in Fig. 6. Values are the result of the average of five experiments, and error bars correspond to one standard deviation. It can be seen that  $k_T$  attains a saturation regime for Ar pressures larger than 5 Torr. This behavior is characteristic of termolecular reactions, which in the low-pressure regime behave as third order and in the high-pressure limit as second order. The termolecular reaction rate constant dependence on total pressure is given by

$$k_T = \frac{k_0([\text{Ar}] + \eta M)}{1 + \frac{k_0}{k_\infty}([\text{Ar}] + \eta M)} \quad (11)$$

where  $k_0$  and  $k_\infty$  are the low- and high-pressure limits. In obtaining Eq. (11), the termolecular reaction is considered as two successive bimolecular processes involving a reactive intermediate. The rate constant as expressed in (11) is obtained assuming a stationary state for the intermediate [24]. The collisions of the reactive intermediate, C<sub>2</sub>Cl<sub>4</sub><sup>\*</sup>, with Ar were differentiated from those with other compounds M (mainly CCl<sub>2</sub> and DCl). The empirical parameter  $\eta$  added in the pressure factor takes into account the stabilization efficiency of the reactive intermediate C<sub>2</sub>Cl<sub>4</sub><sup>\*</sup> through collisions with CCl<sub>2</sub> and DCl relative to those with Ar.

The reaction rate constant values in the low- and high-pressure limit,  $k_0$  and  $k_\infty$ , respectively, resulted:

$$k_0 = (2.23 \pm 0.89) \times 10^{-29} \text{ cm}^6 \text{ molecules}^{-2} \text{ s}^{-1}$$

$$k_\infty = (6.73 \pm 0.23) \times 10^{-13} \text{ cm}^3 \text{ molecules}^{-1} \text{ s}^{-1}$$

where the errors correspond to one standard deviation.

The values of  $k_T$  shown in Fig. 6 determine a mean radical lifetime of 20–40  $\mu\text{s}$ . In Eq. (8), we have neglected dilution effects caused by diffusion. Bialkowski et al. [25] studied the mass transport rates of  $\text{CF}_2$  radicals generated by laser-induced dissociation of  $\text{CF}_2\text{HCl}$  and  $\text{C}_2\text{F}_3\text{Cl}$  diluted in Ar and obtained a diffusion coefficient value of  $D_{\text{CF}_2} = 90 \text{ cm}^2 \text{ Torr s}^{-1}$ . Assuming the same diffusion coefficient value for the  $\text{CCl}_2$  radicals and defining the characteristic time for dilution as the time required for the radicals to travel a path comparable to the IR beam's radius at the focus, the mean lifetime of the radicals ( $\approx 20 \mu\text{s}$ ) for an Ar pressure larger than 5 Torr would be less than the time required for dilution ( $> 100 \mu\text{s}$ ), so that the kinetics of the radicals disappearance would be dominated by the reaction. Although for Ar pressures lower than 5 Torr the mean lifetime of the radicals would be comparable to the estimated dilution time, we expect diffusion effects to be negligible in this pressure range since we obtained an increase in the rate constant of the reaction with pressure. Nevertheless since Cl atoms are heavier and larger than F atoms, we would expect that the diffusion coefficient of  $\text{CCl}_2$  in Ar would be lower than that of  $\text{CF}_2$  so that the conditions described above would be still less restrictive. These arguments are reinforced by the second-order decay of the radicals shown in Fig. 4.

Another important aspect of the present work is that the measurements were performed at room temperature. In the IRMPD of molecules, the radiation interacts resonantly with one of the species present in the sample whereas the other compounds act as a thermal bath. In our experiments, the bath consisted of Ar at room temperature. The final temperature of the products will be determined by the excess of energy attained above the threshold required for dissociation. Herman et al. [11] determined for the IRMPD of  $\text{CDCl}_3$  through the  $\text{DCl}$  elimination channel a threshold energy of  $63 \text{ kcal mol}^{-1}$ . The excess of energy reached above this threshold will be distributed among the rotational and vibrational modes of the fragments, as well as in their translational energy distribution. In [11], the authors determined for the IRMPD products of  $\text{CDCl}_3$  at  $35 \text{ J cm}^{-2}$  a mean translational energy of  $2.8 \text{ kcal mol}^{-1}$ , that is equivalent to the energy of an IR quanta ( $\approx 1000 \text{ cm}^{-1}$ ), showing that the  $\text{CDCl}_3$  molecules are not highly excited above the dissociation threshold. Since  $35 \text{ J cm}^{-2}$  exceeds the value required to saturate the IRMPD of  $\text{CDCl}_3$ , we expect that the translational energy distribution of the fragments obtained in the experiments presented in this work at  $100 \text{ J cm}^{-2}$ , would not differ much from that of [11]. Given that the reaction process can be estimated to occur in about 100–1000 collisions, we can consider that

the fragments obtained in the IRMPD of  $\text{CDCl}_3$  would be thermalized before the recombination takes place.

The value of the factor  $\eta$  was obtained from the comparison of the experimental data with those calculated with Eq. (11) (Fig. 6). This value,  $\eta = 1.52 \pm 0.67$ , close to unity, indicates that collisions with  $\text{CCl}_2$  and  $\text{DCl}$  would be as effective as collisions with Ar for the stabilization of the reactive intermediate  $\text{C}_2\text{Cl}_4^*$  in the low-pressure regime.

On comparison with other experiments, Won and Bozzelli [15] have determined the high-pressure limit values of the rate constant for the association reaction of two  $\text{CCl}_2$  radicals to form  $\text{C}_2\text{Cl}_4$  in the temperature range of 800–1073 K. At 800 K, this value was  $5.1 \times 10^{-12} \text{ cm}^3 \text{ molecules}^{-1} \text{ s}^{-1}$  whereas extrapolation to 300 K would result in  $7.0 \times 10^{-11} \text{ cm}^3 \text{ molecules}^{-1} \text{ s}^{-1}$ . It is important to note that the authors inform that in the lower temperature range employed (848–953 K), their model overpredicts the concentration of  $\text{C}_2\text{Cl}_4$ . Later, Zhu and Bozzelli [20] reported revised high-pressure limit values of this rate constant in the temperature range of 300–2000 K calculated from complex mechanisms used to fit products' distributions from  $\text{CHCl}_3$  pyrolysis using RRKM and QRRK rate theories. At 300 K, the value obtained was  $6.0 \times 10^{-12} \text{ cm}^3 \text{ molecules}^{-1} \text{ s}^{-1}$ . This value results an order of magnitude higher than that obtained in a direct manner in this work.

On the other hand, the high-pressure limit value of the dichlorocarbene radicals' recombination reaction rate constant obtained in this work is in good agreement with the recombination reaction rate constants of other halocarbonated methylene radicals values obtained at room temperature in different pressure ranges and for different buffer gases reported by several authors [25–33]. Particularly, in [32] the laser flash-photolysis absorption technique was used to study the self-recombination reactions of  $\text{CClF}$  radicals as well as with  $\text{CF}_2$  at 297 K. The derived high-pressure rate coefficients obtained  $(1.2 \pm 0.2) \times 10^{-12} \text{ cm}^3 \text{ molecules}^{-1} \text{ s}^{-1}$  and  $(9 \pm 2) \times 10^{-13} \text{ cm}^3 \text{ molecules}^{-1} \text{ s}^{-1}$ , respectively, are in good agreement within experimental error with the value obtained in this work.

To our knowledge, the low-pressure limit value obtained in this work would be the first value reported at 298 K. The only reported value in the literature is that of Kumaran et al. [18], which is valid in the temperature range 1280–1880 K. Extrapolation to 298 K results in  $3.6 \times 10^{-28} \text{ cm}^6 \text{ molecules}^{-2} \text{ s}^{-1}$ , an order of magnitude higher than the value reported in this work. However, this extrapolation is not straightforward due to the temperature range in which the measurements were performed.

## CONCLUSIONS

The LIF technique supported by FTIR spectrometry was used to determine the recombination reaction rate constant of CCl<sub>2</sub> radicals to form C<sub>2</sub>Cl<sub>4</sub> in an Ar bath.

IRMPD of CDCl<sub>3</sub> with a TEA CO<sub>2</sub> laser was used to generate the CCl<sub>2</sub> radicals. The only dissociation pathway observed was the three-center DCl elimination consistent with previous studies.

The major products obtained determined by FTIR spectrometry were DCl and C<sub>2</sub>Cl<sub>4</sub>. The kinetics observed by the LIF technique is consistent with that of the recombination reaction of CCl<sub>2</sub> to form C<sub>2</sub>Cl<sub>4</sub>.

The low- and high-pressure limits of the rate constant of the CCl<sub>2</sub> radicals association reaction at 298 K were obtained in a direct manner, and the values resulted:

$$k_0 = (2.23 \pm 0.89) \times 10^{-29} \text{ cm}^6 \text{ molecules}^{-2} \text{ s}^{-1}$$

$$k_\infty = (6.73 \pm 0.23) \times 10^{-13} \text{ cm}^3 \text{ molecules}^{-1} \text{ s}^{-1}$$

To our knowledge, the low-pressure limit value obtained would be the first value reported at 298 K. The high-pressure limit value obtained results an order of magnitude lower than that calculated from complex mechanisms used to fit product distributions in a previous work of Zhu and Bozzelli [20]. It is consistent with the values obtained for other halocarbonated methylene radicals [25–33].

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