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A quasilinearization method for elliptic problems with a nonlinear boundary condition

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Abstract

We study a nonlinear elliptic second order problem with a nonlinear boundary condition. Assuming the existence of an ordered couple of a supersolution and a subsolution, we develop a quasilinearization method in order to construct an iterative scheme that converges to a solution. Furthermore, under an extra assumption we prove that the convergence is quadratic.

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1. Introduction

In this work, we study the following nonlinear elliptic boundary problem:

$$\begin{cases} \Delta u = f(x, u) & \text{in } \Omega \\ \frac{\partial u}{\partial \eta} = g(x, u) & \text{on } \partial \Omega. \end{cases}$$
 (1.1)

Here Ω is a bounded smooth domain of \mathbb{R}^n , and $f : \overline{\Omega} \times \mathbb{R} \to \mathbb{R}$, $g : \partial \Omega \times \mathbb{R} \to \mathbb{R}$ are continuous and twice continuously differentiable with respect to u.

Nonlinear boundary conditions of this kind appear for example when one considers the problem of finding extremals for the best constant in the Sobolev trace inequality (see e.g. [5]

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and [11]). On the other hand, for n = 1, the problem can be regarded as a mathematical model for the axial deformation of a nonlinear elastic beam, with two nonlinear elastic springs acting at the extremities according to the law u'(0) = -g(u(0)), u'(T) = g(u(T)), and the total force exerted by the nonlinear spring undergoing the displacement u given by f(t, u) [6,14].

The aim of this paper is to develop a quasilinearization technique for problem (1.1) assuming the existence of an ordered couple of a subsolution and a supersolution. More precisely, we construct an iterative scheme that converges to a solution. Furthermore, under an extra assumption we prove that the convergence is quadratic.

The method of supersolutions and subsolutions (definitions will be given in Section 2 below) is one of the most extensively used tools in nonlinear analysis, both for ODE and PDE problems. There exists a vast literature on this subject; see e.g. [4] for a survey. In particular, for elliptic problems with nonlinear boundary conditions such as (1.1), this method has been applied to obtain existence results for example in [7,12].

The quasilinearization method has been developed by Bellman and Kalaba [3], and generalized by Lakshmikantham [9,10]. It has been applied to different nonlinear problems in the presence of an ordered couple of a subsolution and a supersolution. In a recent work [8] it has been successfully applied for a second order ODE Neumann problem for the case in which the supersolution β and the subsolution α present the reversed order, namely $\beta \leq \alpha$.

Our main results read as follows.

Theorem 1.1. Let $\alpha, \beta \in H^1(\Omega) \cap C(\overline{\Omega})$ be respectively a subsolution and a supersolution of (1.1) such that $\alpha \leq \beta$. Furthermore, assume that

$$\frac{\partial^2 f}{\partial u^2}(x, u) \le 0$$

for $x \in \overline{\Omega}$ and $\alpha(x) \le u \le \beta(x)$, and

$$\frac{\partial^2 g}{\partial u^2}(x, u) \ge 0$$

for $x \in \partial \Omega$ and $\alpha(x) \le u \le \beta(x)$. Then the sequence defined below by (3.1) and (3.2) converges in $H^1(\Omega) \cap C(\overline{\Omega})$ to a solution of (1.1).

The proof relies on an associated maximum principle and the unique solvability of the associated linear Robin problem.

Moreover, under a monotonicity condition on f and g, we prove the quadratic convergence of the method.

Theorem 1.2. Under the hypotheses of the previous theorem, assume furthermore that

$$\frac{\partial f}{\partial u}(x, u) > 0$$

for $x \in \overline{\Omega}$ and $\alpha(x) \le u \le \beta(x)$, and

$$\frac{\partial g}{\partial u}(x,u) < 0$$

for $x \in \partial \Omega$ and $\alpha(x) \leq u \leq \beta(x)$. Then the convergence of the sequence defined by (3.1) and (3.2) is quadratic for the $C(\overline{\Omega})$ -norm.

Remark 1.1. An analogous result may be obtained by a similar argument under the following non-local boundary condition:

$$u = c$$
 on $\partial \Omega$, $\int_{\partial \Omega} \frac{\partial u}{\partial \eta} = h(c)$

where c is a constant with unknown value (see e.g. [2]).

Remark 1.2. Our proof of Theorem 1.1 cannot be extended to $f = f(x, u, \nabla u)$, although the existence of solutions can be proved by fixed point arguments (see e.g. [7,12]).

This paper is organized as follows.

In Section 2 we give some preliminary results and definitions. In Section 3, we describe the method of quasilinearization. More precisely, we define iteratively a nondecreasing sequence of subsolutions that converges to a solution of the problem.

Finally, in Section 4 we prove the quadratic convergence of the scheme under a monotonicity condition on f and g.

2. Definitions and preliminary results

We shall make use of the following maximum principle associated with our problem:

Lemma 2.1. Let $\lambda > 0$, $\mu > 0$ and assume that $u \in H^2(\Omega)$ satisfies

$$\begin{cases} \Delta u - \lambda u \ge 0 & \text{in } \Omega \\ \frac{\partial u}{\partial \eta} + \mu u \le 0 & \text{on } \partial \Omega. \end{cases}$$

Then $u \leq 0$.

Proof. We multiply by $u^+ := \max\{u, 0\}$, and integrate by parts:

$$0 \le \int_{\Omega} (\Delta u - \lambda u) u^{+} = -\int_{\Omega^{+}} |\nabla u|^{2} - \lambda \int_{\Omega^{+}} u^{2} + \int_{\partial \Omega} u^{+} \frac{\partial u}{\partial \eta}$$

where $\Omega^+ = \{x \in \Omega : u(x) > 0\}$. Moreover, as $\frac{\partial u}{\partial \eta} \le -\mu u$ on $\partial \Omega$, then

$$\int_{\partial \Omega} u^{+} \frac{\partial u}{\partial \eta} \leq -\mu \int_{\Omega} u^{+} u \leq 0.$$

Hence, we conclude that $u^+ = 0$, and the proof is complete. \square

We shall need the following basic existence and uniqueness result for the Robin problem:

Lemma 2.2. Let $\varphi \in L^2(\Omega)$ and $\psi \in L^2(\partial \Omega)$. Then, for any $\lambda, \mu > 0$ the Robin problem:

$$\begin{cases} \Delta u - \lambda u = \varphi & \text{in } \Omega \\ \frac{\partial u}{\partial \eta} + \mu u = \psi & \text{on } \partial \Omega \end{cases}$$

admits a unique solution $u \in H^2(\Omega)$.

Proof. In order to show existence, we consider the functional $J: H^1(\Omega) \to \mathbb{R}$ given by

$$J(u) = \int_{\Omega} \left[\frac{|\nabla u|^2}{2} + \lambda \frac{u^2}{2} - \varphi u \right] + \int_{\partial \Omega} \left[\mu \frac{u^2}{2} - \psi u \right].$$

It is easy to see that J is coercive and weakly lower semicontinuous; hence it achieves a minimum $u \in H^1(\Omega)$ (see e.g. [13], section 1.1), which is a weak solution. As $\|\Delta u\|_{L^2} = \|\lambda u + \varphi\|_{L^2} < \infty$, we conclude that $u \in H^2(\Omega)$. Uniqueness follows from the maximum principle. \square

Lemma 2.3. Let $\lambda, \mu > 0$, $\varphi \in L^{\infty}(\Omega)$ and $\psi \in L^{\infty}(\partial \Omega)$. Then there exists a constant C such that if u is a weak solution of

$$\begin{cases} \Delta u - \lambda u = \varphi & \text{in } \Omega \\ \frac{\partial u}{\partial \eta} + \mu u = \psi & \text{on } \partial \Omega \end{cases}$$

then

$$||u||_{L^{\infty}(\Omega)} \le C \left[||\varphi||_{L^{\infty}(\Omega)} + ||\psi||_{L^{\infty}(\partial\Omega)} \right].$$

Proof. Multiplying the equation by u and integrating, it follows that

$$\|\varphi\|_{L^{2}}\|u\|_{L^{2}} \geq -\int_{\Omega} \varphi u = \|\nabla u\|_{L^{2}}^{2} + \lambda \|u\|_{L^{2}}^{2} - \int_{\partial\Omega} u \frac{\partial u}{\partial \eta}.$$

Hence

$$\|\nabla u\|_{L^{2}}^{2}+\lambda\|u\|_{L^{2}}^{2}\leq\|\varphi\|_{L^{2}}\|u\|_{L^{2}}+\int_{\partial\Omega}u\psi-\mu\int_{\partial\Omega}u^{2}.$$

From the trace imbedding $H^1(\Omega) \hookrightarrow L^2(\partial\Omega)$ (see e.g. [1], Theorem 5.22) it follows that

$$||u||_{H^1(\Omega)} \le c \left[||\varphi||_{L^2(\Omega)} + ||\psi||_{L^2(\partial\Omega)} \right]$$

for some constant c.

Note that if n = 1, the result follows trivially. For n > 1, fix any p > n. As $W^{1,p}(\Omega) \hookrightarrow L^{\infty}(\Omega) \hookrightarrow L^{p}(\Omega)$, it suffices to show that

$$||u||_{W^{1,p}(\Omega)} \le c \left[||\varphi||_{L^p(\Omega)} + ||\psi||_{L^p(\partial\Omega)} \right]$$

for some constant c. By contradiction, suppose that there exists a sequence $\{u_k\}_{k\in\mathbb{N}}$ such that $\|u_k\|_{W^{1,p}(\Omega)}=1$ and

$$\|\Delta u_k - \lambda u_k\|_{L^p(\Omega)} + \left\|\frac{\partial u_k}{\partial \eta} + \mu u_k\right\|_{L^p(\partial\Omega)} \to 0.$$

As $W^{m,p}(\Omega) \hookrightarrow H^m(\Omega)$, we get from the previous computations that $||u_k||_{H^1(\Omega)} \to 0$, and as $\Delta u_k - \lambda u_k \to 0$ for the L^2 -norm it follows that $||u_k||_{H^2(\Omega)} \to 0$.

Then $\|u_k\|_{W^{1,2^*}(\Omega)} \to 0$, and repeating the previous argument we deduce that $\|u_k\|_{W^{2,2^*}(\Omega)} \to 0$. If $2^* \geq n$, then $W^{2,2^*}(\Omega) \hookrightarrow W^{1,p}(\Omega)$ and $\|u_k\|_{W^{1,p}(\Omega)} \to 0$, a contradiction. Otherwise, repeating the argument a certain number of times we obtain that $\|u_k\|_{W^{2,q}(\Omega)} \to 0$ for some $q \geq n$, and the proof follows. \square

Next, we recall the definition of the concept of supersolution and subsolution. We say that $\alpha \in H^1(\Omega) \cap C(\overline{\Omega})$ is a subsolution of problem (1.1) if it satisfies

$$\begin{cases} \Delta \alpha \geq f(x, \alpha) & \text{in } \Omega \\ \frac{\partial \alpha}{\partial \eta} \leq g(x, \alpha) & \text{on } \partial \Omega. \end{cases}$$

In the same way, $\beta \in H^1(\Omega) \cap C(\overline{\Omega})$ is a supersolution if it satisfies

$$\begin{cases} \Delta \beta \leq f(x,\beta) & \text{in } \Omega \\ \frac{\partial \beta}{\partial n} \geq g(x,\beta) & \text{on } \partial \Omega. \end{cases}$$

3. The quasilinearization method

In this section we define the quasilinearization method and give a proof of Theorem 1.1.

Assume that α and β are respectively a subsolution and a supersolution of (1.1) with $\alpha \leq \beta$, and fix λ , $\mu > 0$ such that

$$\lambda > \max_{x \in \overline{\Omega}, \alpha(x) \le u \le \beta(x)} \frac{\partial f}{\partial u}(x, u)$$

$$\mu > -\min_{x \in \partial \Omega, \alpha(x) \le u \le \beta(x)} \frac{\partial g}{\partial u}(x, u).$$

We define recursively a sequence $\{u_n\}$ in the following way. Set $u_0 = \alpha$, and assuming that u_n is known, define u_{n+1} as a solution to the following quasilinear Robin problem:

$$\Delta u_{n+1} - \lambda u_{n+1} = f(x, u_n) + \frac{\partial f}{\partial u}(x, u_n) [P_n(x, u_{n+1}) - u_n] - \lambda P_n(x, u_{n+1})$$
(3.1)

in the domain Ω , with the boundary condition

$$\frac{\partial u_{n+1}}{\partial \eta} + \mu u_{n+1} = g(x, u_n) + \frac{\partial g}{\partial u}(x, u_n)[P_n(x, u_{n+1}) - u_n] + \mu P_n(x, u_{n+1})$$
(3.2)

on $\partial \Omega$, where

$$P_n(x, u) = \begin{cases} u_n(x) & \text{if } u < u_n(x) \\ u & \text{if } u_n(x) \le u \le \beta(x) \\ \beta(x) & \text{if } u > \beta(x). \end{cases}$$

We shall see in the proof of Theorem 1.1 that $u_n \le \beta$ and consequently that P_n is well defined. On the other hand, as $u_n(x) \le P_n(x, u) \le \beta(x)$, a straightforward application of the Schauder Theorem shows that (3.1) and (3.2) admits at least one solution and therefore u_{n+1} is well defined.

Furthermore, from the fact that $\alpha \le u_n \le \beta$ it will follow that the sequence defined by (3.1) and (3.2) is indeed a Newton scheme (see (3.3) below). However, it is not possible to apply the Newton method directly since the linearized problem (3.3) might fail to have a unique solution.

Proof of Theorem 1.1. We will prove by induction that u_n is a subsolution, and that $\alpha \le u_n \le u_{n+1} \le \beta$.

First, observe that

$$\Delta u_{n+1} - \lambda u_{n+1} = f(x, u_n) + \left[\frac{\partial f}{\partial u}(x, u_n) - \lambda \right] [P_n(x, u_{n+1}) - u_n] - \lambda u_n.$$

From our choice of λ , $\frac{\partial f}{\partial u}(x, u_n) - \lambda \leq 0$, and by the inductive hypothesis $\Delta u_n \geq f(x, u_n)$. As $P_n(x, u_{n+1}) - u_n \geq 0$, it follows that

$$\Delta u_{n+1} - \lambda u_{n+1} \le \Delta u_n - \lambda u_n.$$

In a similar way, using the fact that $\frac{\partial u_n}{\partial \eta} + \mu u_n \leq g(x, u_n)$, we conclude that

$$\frac{\partial u_{n+1}}{\partial \eta} + \mu u_{n+1} = g(x, u_n) + \left[\frac{\partial g}{\partial u}(x, u_n) + \mu \right] [P_n(x, u_{n+1}) - u_n] + \mu u_n$$

$$\geq \frac{\partial u_n}{\partial n} + \mu u_n.$$

From the maximum principle, it follows that $u_{n+1} \ge u_n$.

In order to show that $u_{n+1} \leq \beta$, we use a Taylor expansion:

$$f(x,v) = f(x,u_n) + \frac{\partial f}{\partial u}(x,u_n)(v-u_n) + \frac{1}{2}\frac{\partial^2 f}{\partial u^2}(x,\xi)(v-u_n)^2$$

for some ξ between $u_n(x)$ and v.

Choosing $v = P_n(x, u_{n+1})$, as $\frac{\partial^2 f}{\partial u^2}(x, \xi) \le 0$ we deduce that

$$f(x, P_n(x, u_{n+1})) \le f(x, u_n) + \frac{\partial f}{\partial u}(x, u_n)(P_n(x, u_{n+1}) - u_n).$$

Hence, from the definition of u_{n+1} ,

$$\Delta u_{n+1} - \lambda u_{n+1} \ge f(x, P_n(x, u_{n+1})) - \lambda P_n(x, u_{n+1}).$$

Moreover, from the choice of λ , the mapping $u \mapsto f(x, u) - \lambda u$ is nonincreasing and, therefore,

$$\Delta u_{n+1} - \lambda u_{n+1} \ge f(x, \beta) - \lambda \beta \ge \Delta \beta - \lambda \beta.$$

In a similar way,

$$\frac{\partial u_{n+1}}{\partial \eta} + \mu u_{n+1} \le g(x, P_n(x, u_{n+1})) + \mu P_n(x, u_{n+1})$$
$$\le g(x, \beta) + \mu \beta \le \frac{\partial \beta}{\partial n} + \mu \beta.$$

By the maximum principle, we conclude that $u_{n+1} \leq \beta$.

Next we observe that, as it has been already proved that $\alpha \le u_{n+1} \le \beta$, the definition of u_{n+1} reduces to a Newton iteration:

$$\begin{cases} \Delta u_{n+1} = f(x, u_n) + \frac{\partial f}{\partial u}(x, u_n)[u_{n+1} - u_n] & \text{in } \Omega \\ \frac{\partial u_{n+1}}{\partial \eta} = g(x, u_n) + \frac{\partial g}{\partial u}(x, u_n)[u_{n+1} - u_n] & \text{on } \partial \Omega. \end{cases}$$
(3.3)

Moreover, using the Taylor expansion as before we obtain

$$\Delta u_{n+1} = f(x, u_{n+1}) - \frac{1}{2} \frac{\partial^2 f}{\partial u^2}(x, \xi) (u_{n+1} - u_n)^2 \ge f(x, u_{n+1}) \quad \text{in } \Omega,$$

and similarly

$$\frac{\partial u_{n+1}}{\partial \eta} \le g(x, u_{n+1}) \quad \text{on } \partial \Omega.$$

Hence, u_{n+1} is a subsolution of the problem.

As $\{u_n\}$ is monotone nondecreasing, it converges pointwise to some function u, with $\alpha(x) \le u(x) \le \beta(x)$. Moreover, from the proof of Lemma 2.3

$$||u_{n+1}||_{H^{1}(\Omega)} \leq C_{1} ||f(x, u_{n}) + \frac{\partial f}{\partial u}(x, u_{n})[u_{n+1} - u_{n}] - \lambda u_{n}||_{L^{2}(\Omega)} + C_{2} ||g(x, u_{n}) + \frac{\partial g}{\partial u}(x, u_{n})[u_{n+1} - u_{n}] + \mu u_{n}||_{L^{2}(\partial \Omega)}.$$

Hence, as $\alpha \leq u_n \leq \beta$, it follows that $\{u_n\}$ is bounded in $H^1(\Omega)$, and from (3.3) we deduce that it is bounded in $H^2(\Omega)$. Moreover, by the compactness of the imbedding $H^2(\Omega) \hookrightarrow H^1(\Omega)$ there exists a subsequence u_{n_k} such that $u_{n_k} \to u$ in $H^1(\Omega)$. For any test function $\varphi \in H^1_0(\Omega)$,

$$-\int_{\Omega} \nabla u_{n_k} \nabla \varphi = \int_{\Omega} \left[f(x, u_{n_k-1}) + \frac{\partial f}{\partial u}(x, u_{n_{k-1}})(u_{n_k} - u_{n_k-1}) \right] \varphi.$$

By dominated convergence, we conclude that $-\int_{\Omega} \nabla u \nabla \varphi = \int_{\Omega} f(x, u) \varphi$, and hence u is a weak solution of

$$\Delta u = f(x, u).$$

Therefore $u \in H^2(\Omega)$. Then, if we consider again a test function $\varphi \in H^1(\Omega)$,

$$-\int_{\Omega} \nabla u_{n_k} \nabla \varphi + \int_{\partial \Omega} \left[g(x, u_{n_k-1}) + \frac{\partial g}{\partial u}(x, u_{n_k-1})(u_{n_k} - u_{n_k-1}) \right] \varphi$$
$$= \int_{\Omega} \left[f(x, u_{n_k-1}) + \frac{\partial f}{\partial u}(x, u_{n_k-1})(u_{n_k} - u_{n_k-1}) \right] \varphi.$$

By dominated convergence,

$$-\int_{\Omega} \nabla u \nabla \varphi + \int_{\partial \Omega} g(x, u) \varphi = \int_{\Omega} f(x, u) \varphi.$$

Using the divergence theorem and the fact that the range of the trace operator $H^1(\Omega) \hookrightarrow L^2(\partial\Omega)$ is $H^{1/2}(\partial\Omega)$ (see e.g. [1], 7.56), which is dense in $L^2(\partial\Omega)$, we conclude that $\frac{\partial u}{\partial\eta} = g(x,u)$ on $\partial\Omega$. \square

4. Quadratic convergence

In this section we give a proof of Theorem 1.2.

Proof of Theorem 1.2. Let us define the error $\mathcal{E}_n = u - u_n$. Using the Taylor expansion around u_n , it follows from (3.3) that

$$\begin{cases} \Delta \mathcal{E}_{n+1} = f(x,u) - f(x,u_{n+1}) + \frac{1}{2} \frac{\partial^2 f}{\partial u^2}(x,\xi) [u_{n+1} - u_n]^2 & \text{in } \Omega \\ \frac{\partial \mathcal{E}_{n+1}}{\partial n} = g(x,u) - g(x,u_{n+1}) + \frac{1}{2} \frac{\partial^2 g}{\partial u^2}(x,\theta) [u_{n+1} - u_n]^2 & \text{on } \partial \Omega \end{cases}$$

$$(4.1)$$

for some $\xi = \xi(x) \in [u_n(x), u_{n+1}(x)]$ for $x \in \overline{\Omega}$, and $\theta = \theta(x) \in [u_n(x), u_{n+1}(x)]$ for $x \in \partial \Omega$. Set

$$\lambda = \min_{x \in \overline{\Omega}, \alpha(x) \le u \le \beta(x)} \frac{\partial f}{\partial u}(x, u) > 0,$$

$$\mu = -\max_{x \in \partial \Omega, \alpha(x) \le u \le \beta(x)} \frac{\partial g}{\partial u}(x, u) > 0.$$

As $\mathcal{E}_{n+1} > 0$, it follows that

$$f(x, u) - f(x, u_{n+1}) \ge \lambda \mathcal{E}_{n+1}$$

and

$$g(x, u) - g(x, u_{n+1}) \le -\mu \mathcal{E}_{n+1}$$
.

Thus, we may define ϕ as the unique solution of the linear Robin problem

$$\begin{cases} \Delta \phi - \lambda \phi = \frac{1}{2} \frac{\partial^2 f}{\partial u^2}(x, \xi) \mathcal{E}_n^2 & \text{in } \Omega \\ \frac{\partial \phi}{\partial n} + \mu \phi = \frac{1}{2} \frac{\partial^2 g}{\partial u^2}(x, \theta) \mathcal{E}_n^2 & \text{on } \partial \Omega. \end{cases}$$

Using the fact that $0 \le u_{n+1} - u_n \le \mathcal{E}_n$ and that $\frac{\partial^2 f}{\partial u^2}(x, \xi) \le 0 \le \frac{\partial^2 g}{\partial u^2}(x, \theta)$ we conclude

$$\begin{split} &\Delta \mathcal{E}_{n+1} - \lambda \mathcal{E}_{n+1} \geq \Delta \phi - \lambda \phi & \text{in } \Omega, \\ &\frac{\partial \mathcal{E}_{n+1}}{\partial \eta} + \mu \mathcal{E}_{n+1} \leq \frac{\partial \phi}{\partial \eta} + \mu \phi & \text{on } \partial \Omega. \end{split}$$

By the maximum principle, $\mathcal{E}_{n+1} \leq \phi$, and from Lemma 2.3 we deduce that

$$\|\phi\|_{L^{\infty}(\varOmega)} \leq c \left\{ \left\| \frac{\partial^2 f}{\partial u^2}(x,\xi)\mathcal{E}_n^2 \right\|_{L^{\infty}(\varOmega)} + \left\| \frac{\partial^2 g}{\partial u^2}(x,\theta)\mathcal{E}_n^2 \right\|_{L^{\infty}(\partial \varOmega)} \right\}.$$

Hence.

$$0 \le \mathcal{E}_{n+1} \le c \|\mathcal{E}_n\|_{L^{\infty}(\Omega)}^2$$

for some constant c independent of n, and the proof is complete. \square

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References

- [1] R. Adams, Sobolev Spaces, Academic Press, NY, 1975.
- [2] P. Amster, P. De Nápoli, M.C. Mariani, Existence of solutions to n-dimensional pendulum-like equations, Electron. J. Differential Equations 2004 (125) (2004) 1–8.
- [3] R. Bellman, R. Kalaba, Quasilinearisation and Nonlinear Boundary Value Problems, American Elsevier, New York, 1965.
- [4] C. De Coster, P. Habets, An overview of the method of lower and upper solutions for ODE, in: M.R. Grossinho, M. Ramos, C. Rebelo, L. Sanchez (Eds.), Nonlinear analysis and its Applications to Differential Equations, in: Progress in Nonlinear Differential Equations and their Applications, vol. 43, Birkhauser, Boston, 2001, pp. 3–22.

- [5] J. Fernández Bonder, J.D. Rossi, Existence results for the *p*-Laplacian with nonlinear boundary conditions, J. Math. Anal. Appl. 263 (2001) 195–223.
- [6] M. Grossinho, T.F. Ma, Symmetric equilibria for a beam with a nonlinear elastic foundation, Port. Math. 51 (1994) 375–393.
- [7] F. Inkmann, Existence and multiplicity theorems for semilinear elliptic equations with nonlinear boundary conditions, Indiana Univ. Math. J. 31 (2) (1982) 213–221.
- [8] R. Khan, Existence and approximation of solutions of second order nonlinear Neumann problems, Electron. J. Differential Equations 2005 (3) (2005) 1–10.
- [9] V. Lakshmikantham, An extension of the method of quasilinearization, J. Optim. Theory Appl. 82 (1994) 315–321.
- [10] V. Lakshmikantham, Further improvement of generalized quasilinearization, Nonlinear Anal. 27 (2) (1996) 223–227.
- [11] S. Martínez, J. Rossi, Weak solutions for the *p*-laplacian with a nonlinear boundary condition at resonance, Electron. J. Differential Equations 2003 (27) (2003) 1–14.
- [12] J. Mawhin, K. Schmitt, Upper and lower solutions and semilinear second order elliptic equations with nonlinear boundary conditions, Proc. Roy. Soc. Edinburgh 97 A (1984) 199–207.
- [13] J. Mawhin, M. Willem, Critical Point Theory and Hamiltonian Systems, Springer-Verlag, NY, Berlin, Heidelberg, 1989.
- [14] C. Rebelo, L. Sanchez, Existence and multiplicity for an O.D.E. with nonlinear boundary conditions, Differential Equations Dynam. Systems 3 (4) (1995) 383–396.