



Vortex pinning: A probe for nanoscale disorder in iron-based superconductors

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ABSTRACT

The pinning of quantized flux lines, or vortices, in the mixed state is used to quantify the effect of impurities in iron-based superconductors (IBS). Disorder at two length scales is relevant in these materials. Strong flux pinning resulting from nm-scale heterogeneity of the superconducting properties leads to the very disordered vortex ensembles observed in the IBS, and to the pronounced maximum in the critical current density j_c at low magnetic fields. Disorder at the atomic scale, most likely induced by the dopant atoms, leads to “weak collective pinning” and a magnetic field-independent contribution j_c^{coll} . The latter allows one to estimate quasiparticle scattering rates.

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1. Introduction

Quantifying the effects of material disorder is important for the understanding of the superconducting ground state of iron-based superconductors (IBS). On the one hand, the proximity of the superconducting state to anti-ferromagnetism may lead to phase segregation that will be heavily influenced, for example, by macroscopic heterogeneity of the chemical composition. On the other hand, the superconducting ground state in these materials was proposed to have a so-called s_{\pm} symmetry [1,2], in which the superconducting gap not only has a different value on different

Fermi surface sheets, but may also change sign from one sheet to another. Then, superconductivity is thought to be exquisitely sensitive to interband scattering [3,4].

The effect of impurities is usually characterized by the scattering rate such as this can be extracted from resistivity measurements in the normal state, or from the surface resistance in the superconducting state. Another well-known but little exploited probe of microscopic disorder is the pinning of vortex lines in the superconducting mixed state. The radius of their core, of the order of the coherence length $\xi \sim 1.5$ nm, makes vortices ideal probes for impurities. The fact that the vortex density can be easily varied by orders of magnitude simply by adjusting the value of the applied magnetic field H_a makes “vortex matter” sensitive to local variations of material properties on different length scales. The bulk pinning force F_p exerted by the material disorder on the vortex ensemble inhibits the latter's motion for currents smaller than the critical current density $j_c = F_p/B$. In the following, we analyze vortex distributions and j_c -data in IBS, show how these can be used to characterize the type of disorder, and to bracket impurity scattering rates.

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2. Critical currents

Fig. 1 shows hysteresis loops of the irreversible magnetization $M(H_a)$ versus the applied magnetic field H_a , measured on a single crystal of $\text{Ba}(\text{Fe}_{0.925}\text{Co}_{0.075})_2\text{As}_2$, with critical temperature $T_c = 23.5$ K, using a commercial Superconducting Quantum Interference Device magnetometer. As in all IBS, the hysteretic magnetic moment features a pronounced low-field maximum, superposed on a nearly field-independent contribution. A double-logarithmic plot of the critical current density as this follows from the Bean model [5], $j_c \sim 3 \text{ M/a}$ (with a the crystal width), shows that the low-field maximum amounts to a plateau $j_c(0)$, followed by a decrease $j_c(B) \sim B^{-1/2}$ (Fig. 2). Such a behavior is naturally described in terms of strong pinning by sparse, extrinsic defects of dimension greater than ξ , and density $n_i \ll \xi^{-3}$ [6,7]. The plateau-value

$$j_c(0) \sim (f_p / \sqrt{\pi} \Phi_0 \varepsilon) (U_p n_i / \varepsilon_0)^{1/2} \quad (1)$$

and the decrease $j_c(B) \sim (f_p / \Phi_0 \varepsilon) (U_p n_i / \varepsilon_0) (\Phi_0 / B)^{1/2}$ can be parameterized in terms of the elementary pinning force f_p of a single strong defect, and the ratio of the energy U_p gained by placing a vortex on a defect to the vortex line energy ε_0 [8] ($\varepsilon \approx 0.4$ is the material anisotropy parameter [9] and $\Phi_0 = h/2e$ is the flux quantum). Elimination of $(U_p n_i / \varepsilon_0)^{1/2}$ directly yields the experimental value $f_p = 3 \times 10^{-13} \text{ N}$ at low T .

Further measurements were performed using the magneto-optical- [8,10] and Hall-array techniques [10,11]. The temperature dependence $j_c(0, T)$, and that of the field-independent contribution $j_c^{\text{coll}}(T)$ observed at fields exceeding $\sim 1 \text{ T}$ is plotted in Fig. 3. In the following, we first discuss the strong pinning contribution, before turning to j_c^{coll} in Section 4.

3. Heterogeneity and strong pinning

Fig. 4 shows the vortex ensemble in single crystalline $\text{Ba}(\text{Fe}_{0.9}\text{Co}_{0.1})_2\text{As}_2$ (with $T_c = 19.5 \text{ K}$), at $H_a = 10 \text{ G}$ (1 mT). Vortex positions are revealed using Bitter decoration at 4.2 K, after field-cooling through the superconducting transition [8]. The featureless Fourier

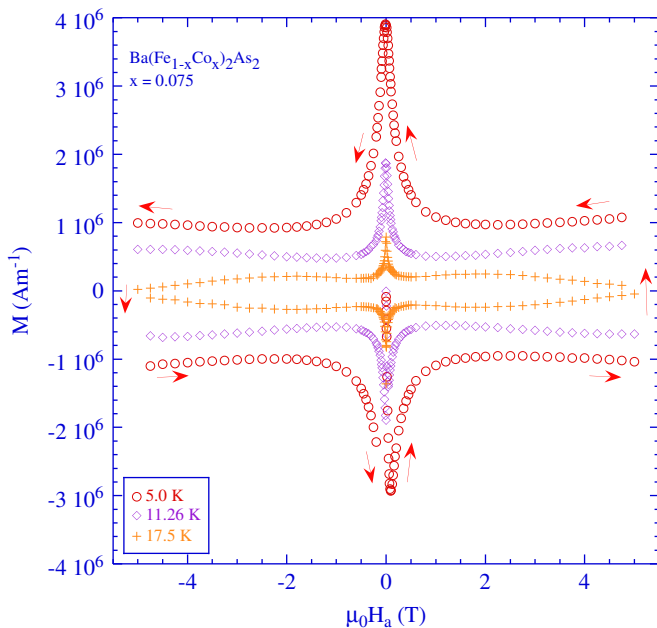


Fig. 1. Loops of the hysteretic magnetization $M(H_a)$ versus the magnetic field H_a , measured on a $\text{Ba}(\text{Fe}_{0.925}\text{Co}_{0.075})_2\text{As}_2$ single crystal, at different indicated temperatures.

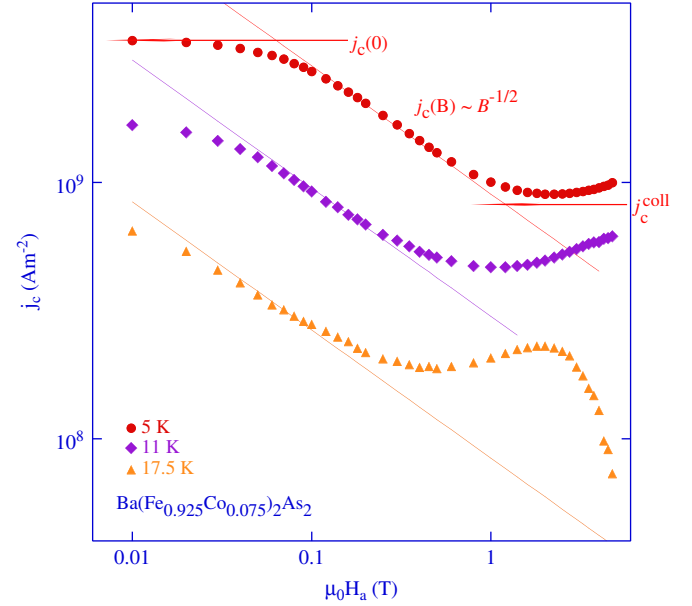


Fig. 2. Field-dependence of the critical current density $j_c \sim 3 \text{ M/a}$ (with a the crystal width), as extracted from Fig. 1.

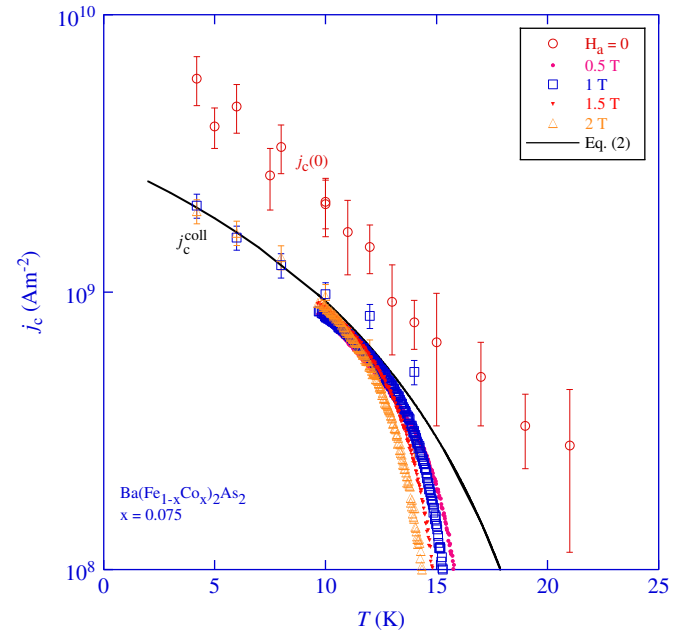


Fig. 3. Temperature dependence of j_c at different applied fields. Values for $\mu_0 H_a \gtrsim 1 \text{ T}$ are representative of $j_c^{\text{coll}}(T)$. The drawn line is a fit to Eq. (2), with $\sigma_{\text{tr}} = 3 \text{ \AA}^2$.

transform of the set of vortex positions indicates the absence of long-range positional- or orientational order, and the presence of large fluctuations in the nearest-neighbor distance. The very disordered vortex structure is the combined result of the narrow temperature trajectory over which it is frozen in during field-cooling and the importance of flux pinning at high T [8].

To characterize pinning, the vortex interaction energies, $\mathcal{E}_{\text{int}}^i = \sum_j (\Phi_0^2 / 2\pi\mu_0\lambda_{ab}^2) K_0(r_{ij}/\lambda_{ab})$, and the force to which each vortex is subjected, $\mathbf{f}_i = \sum_j (\Phi_0^2 / 2\pi\mu_0\lambda_{ab}^2) (\mathbf{r}_{ij}/|\mathbf{r}_{ij}|) K_1(r_{ij}/\lambda_{ab})$, were determined from the inter-vortex distance ensemble $\{r_{ij}\}$ (λ_{ab} is the ab -plane penetration depth [12], and K_0 and K_1 are modified Bessel functions). The result is mapped out in Fig. 5. The probability distribution of $\mathcal{E}_{\text{int}}^i$, depicted in Fig. 6a, shows that the

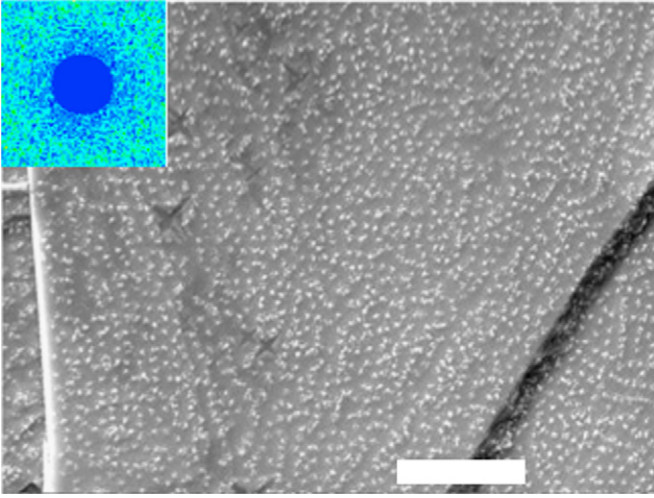


Fig. 4. Bitter decoration of vortices in a $\text{Ba}(\text{Fe}_{0.9}\text{Co}_{0.1})_2\text{As}_2$ single crystal, for a field $\mu_0 H_a = 1$ mT applied parallel to the crystalline c -axis. The scale bar represents 10 μm . Inset: Fourier transform of vortex positions.

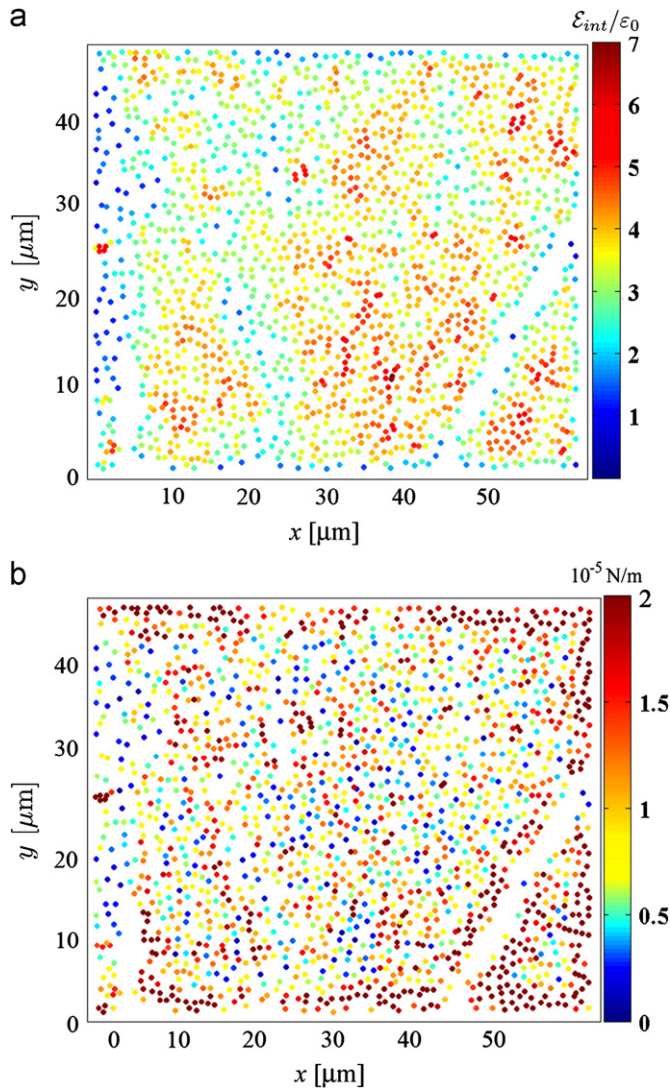


Fig. 5. Maps of (a) the vortex interaction energy and (b) the pinning force, obtained from the image in Fig. 4.

mean interaction energy $\langle \mathcal{E}_{int}^i \rangle$ exceeds the value \mathcal{E}_A for vortices arranged in a triangular vortex lattice of the same (average) density. Also, the distribution is considerably broadened with respect to the regular lattice δ -function. We interpret the \mathcal{E}_{int} -distribution as being determined by the contribution from the vortex pinning energy, with mean $\langle \mathcal{E}_p \rangle = \langle \mathcal{E}_{int} \rangle - \mathcal{E}_A \approx 0.5\mathcal{E}_0$. This large value of $\langle \mathcal{E}_p \rangle$ signifies that pinning cannot originate from well-defined defects in a homogeneous superconducting matrix, but must be interpreted as arising from the heterogeneity of the superconducting properties of the material [8].

In order to verify whether such heterogeneity can account for j_c , we turn to the distribution of $|\mathbf{f}_i|$. Since the vortex system is at rest, Newton's third law implies that the force map (Fig. 5b) and the histogram of $|\mathbf{f}_i|$ (Fig. 6b) are to be interpreted as representing the local pinning forces \mathbf{f}_i acting on each vortex. The critical current density will be determined by the average force $\langle \mathbf{f}_i \rangle \sim 5 \times 10^{-6} \text{ N m}^{-1}$, rather than by the most strongly pinned vortices.

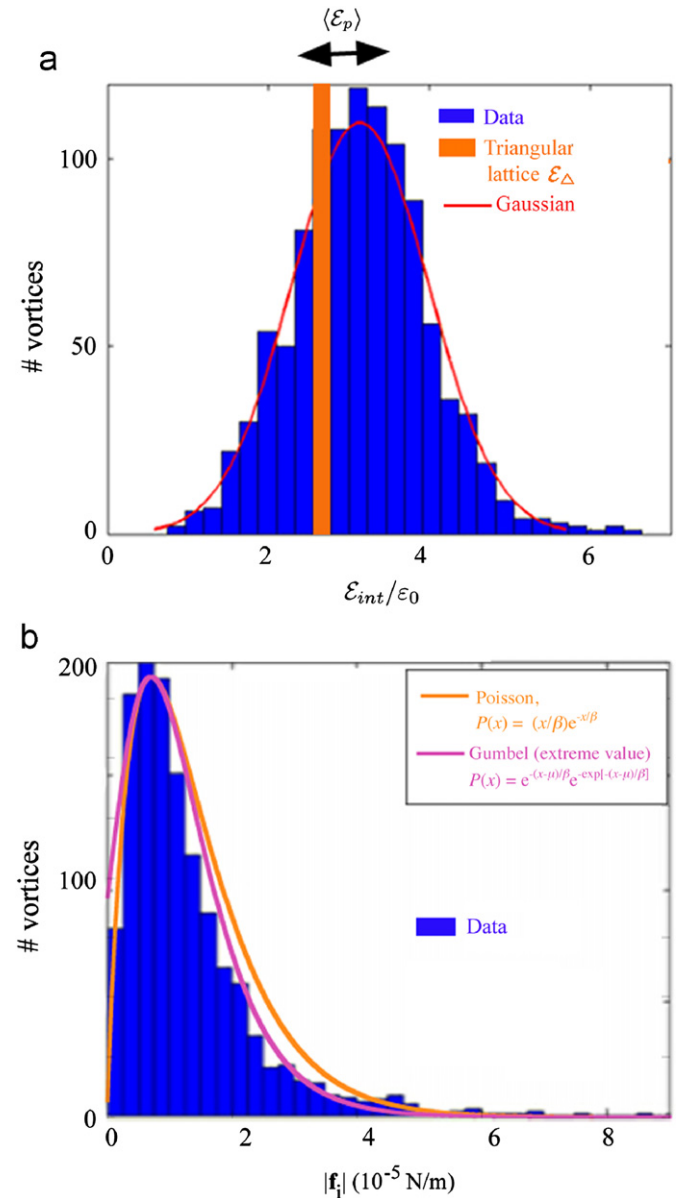


Fig. 6. Histograms of the vortex interaction energy (a) and of the pinning force (b) distribution of Fig. 5. The first is well approximated by a Gaussian, while the force distribution (b) resembles a Poisson- or extreme value distribution. The energy distribution in (a) is compared to the δ -peak distribution for a (ordered) triangular vortex lattice of the same density.

The effective pin density can now be estimated from the ratio of the pinning force per vortex, $\langle \mathbf{f}_i \rangle$, and the elementary force per pin, $f_p = 3 \times 10^{-13}$ N determined above. This yields an average distance between effective pins of 60 nm, and a pin density $n_i \sim 5 \times 10^{21} \text{ m}^{-3} \ll \xi^{-3}$. Strong pinning in $\text{Ba}(\text{Fe}_{1-x}\text{Co}_x)_2\text{As}_2$ must therefore be due to heterogeneity on the nm scale. Substituting the obtained values in Eq. (1) reproduces the magnitude of j_c , for a spatial variance of the line energy $\Delta\epsilon_0/\epsilon_0 \approx 0.05$.

4. Quasiparticle scattering and weak pinning

The intermediate field ($\mu_0 H_a \gtrsim 1$ T) plateau of constant $j_c = j_c^{\text{coll}}$ is attributed to weak collective pinning [13] by fluctuations of the dopant atom density $n_d \gg \xi^{-3}$ at length scales much less than the coherence length [15]. This contribution to pinning is found in all charge-doped IBS, as well as in Ru-doped BaFe_2As_2 and $\text{FeSe}_{1-x}\text{Te}_x$ [14], but not in the P-doped materials [15]. The magnitude and temperature dependence of j_c^{coll} indicate [15] that dopant atom density fluctuations are effective through the variation of the quasiparticle mean free-path they entail [13,15,16]. In the field-regime where vortex lines are pinned independently from their neighbors (i.e. the “single-vortex regime” of collective pinning [13]), the resulting critical current contribution reads [15]

$$j_c^{\text{coll}} = \frac{\Phi_0}{\sqrt{3}\mu_0\lambda_{ab}^2\xi} \left[\frac{0.01n_d\sigma_{tr}^2}{\epsilon\xi} \left(\frac{\xi_0}{\xi} \right) \right]^{2/3} \quad (2)$$

(with $\xi_0 = 1.35\xi$ the T -independent Bardeen–Cooper–Schrieffer coherence length). Fig. 3 shows that, after insertion of a quasiparticle transport scattering cross-section $\sigma_{tr} \approx 3 \text{ \AA}^2$ per Co ion, and of the $\lambda_{ab}(T)$ -dependence of Ref. [12], Eq. (2) also satisfactorily describes both the magnitude and the T -dependence of the intermediate-field critical current density of our $\text{Ba}(\text{Fe}_{1-x}\text{Co}_x)_2\text{As}_2$ single crystals. The extracted value of the transport cross-section can be related to the elastic quasiparticle mean-free path as $l = (n_d\sigma_{tr})^{-1} \approx 26 \text{ nm}$, to the scattering phase angle as $\sin \delta_0 = (k_F^2\sigma_{tr}/2\pi)^{1/2} \approx 0.2$, and to

the quasiparticle scattering rate $\Gamma = n_d[\pi N(0)]^{-1} \sin^2 \delta_0 \approx 3 \text{ meV}$ (here $N(0) \sim mk_F/\pi^2 \hbar^2$ is the normal state Density of States and m is the electronic mass) [15].

5. Summary and conclusions

The bulk vortex pinning force F_p and the critical current density j_c in single crystalline $\text{Ba}(\text{Fe}_{1-x}\text{Co}_x)_2\text{As}_2$ are representative of that found in other iron-based superconductors and can be consistently described in terms of two additive contributions. At low magnetic fields, strong pinning by nm-scale variations of the vortex line energy (of the order of 5%) is the most relevant. Such variations may arise from the inhomogeneity of the gap [17] or from that of the superfluid density [12]. At higher fields, the effect of quasiparticle mean-free path variations due to spatial fluctuations of the dopant atom density is dominant.

References

- [1] I.I. Mazin, D.J. Singh, M.D. Johannes, M.H. Du, Phys. Rev. Lett. 101 (2008) 057003.
- [2] K. Kuroki, et al., Phys. Rev. Lett. 101 (2008) 087004; K. Kuroki, et al., Phys. Rev. Lett. 102 (2009) 109902.
- [3] K. Kuroki, H. Usui, S. Onari, R. Arita, H. Aoki, Phys. Rev. B 79 (2009) 224511.
- [4] S. Onari, H. Kontani, Phys. Rev. Lett. 103 (2009) 177001.
- [5] E.H. Brandt, Phys. Rev. B 58 (1998) 6506.
- [6] Y.N. Ovchinnikov, B.I. Ivlev, Phys. Rev. B 43 (1991) 8024.
- [7] C.J. van der Beek, et al., Phys. Rev. B 66 (2002) 024523.
- [8] S. Demirdiř, et al., 2011 arXiv:1106.6065.
- [9] J. Hänisch, et al., IEEE Trans. Appl. Superconductivity 21 (2010) 2887.
- [10] C.J. van der Beek, et al., Phys. Rev. B 81 (2010) 174517.
- [11] M. Konczykowski, et al., 2011 arXiv:1108.5583.
- [12] L. Luan, et al., Phys. Rev. Lett. 106 (2011) 067001.
- [13] G. Blatter, et al., Rev. Mod. Phys. 66 (1994) 1125.
- [14] M. Konczykowski, et al., Bull. Am. Phys. Soc. 56 (2011) J26-00002.
- [15] C.J. van der Beek, et al., Phys. Rev. Lett. 105 (2011) 267002.
- [16] E.V. Thuneberg, J. Kurkijärvi, D. Rainer, Phys. Rev. Lett. 48 (1982) 1853; E.V. Thuneberg, J. Kurkijärvi, D. Rainer, Phys. Rev. B 29 (1984) 3913.
- [17] F. Massee, et al., Phys. Rev. B 79 (2009) 220517.