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One-dimensional tungsten oxide nanostructures on Cu(110) surface

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KEYWORDS Tungsten oxide, model system, one-dimensional structures

ABSTRACT Thin epitaxial layers of tungsten oxide on metal substrates are suitable as model systems for investigation of chemical reactivity and catalytic properties. However, the ability to prepare epitaxial tungsten oxide model system in situ is quite rare. Here we present a method to prepare highly ordered tungsten oxide thin film on a Cu(110) single crystal substrate using physical vapor deposition in a reactive atmosphere of atomic oxygen. The oxygen induced reconstruction of the copper substrate gives rise to unique self-organized one-dimensional structures of tungsten oxide parallel with the Cu[1-10] crystallographic direction. Utilizing a combination of photoemission spectroscopy and density functional theory calculations we reveal emergent physicochemical properties related to the low-dimensionality of the system.

Specifically, we observe a support mediated charge redistribution at the interface and a momentum dependent modulation of the valence-band electronic structure. The unusual character of the one-dimensional oxide nanostructures on Cu(110) surface opens up a unique avenue for preparation of tungsten oxide-based functionalized nanostructures and provides options for further investigation of the fundamental properties of tungsten oxide.

Introduction

One-dimensional nanostructures with high surface to volume ratio can exhibit unique physicochemical properties compared with their bulk counterparts [1] and provide an intriguing field of research [2,3]. In particular, low-dimensional oxides on metal substrates have been extensively studied, revealing novel structural properties, electronic and magnetic behavior and also enhanced chemical reactivity [4,5]. Apart from the reduced dimensionality, another important feature of oxide nanostructures supported on metallic substrates is the proximity of the oxide-metal interface, which can significantly influence their electronic structure and catalytic performance. Accordingly, low-dimensional oxide systems are valuable for innovative model studies in the fields of heterogeneous catalysis and gas sensing.

Tungsten oxide is widely used in industrial processes as a catalyst and a sensor both in the pure form and doped by different metals [6–8]. The microstructure of reducible tungsten oxide can have a significant influence on its physicochemical properties. In non-stoichiometric tungsten, oxide W^{6+} ions are reduced to W^{5+} or W^{4+} ions through the formation of oxygen vacancies. These may serve as active sites for adsorption and determine selective catalytic processes taking place on the surface [9]. Vacancies alone or species adsorbed from reaction atmosphere also influence conductivity of the oxide [10,11].

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3 Formation of tungsten oxide nanowires was reported on a mica substrate due to interaction
4 with potassium atoms in the surface layer [12,13]. However, this system is not suitable for model
5 studies due to mica's insulating properties severely restricting applicable surface science methods.
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7 Other forms of low-dimensional tungsten oxide nanostructures have been prepared [14,15],
8 including a highly ordered monolayer of ternary tungsten oxide CuWO_4 grown on the (110)
9 plane of a copper single crystal [16]. The low-index Cu(110) surface is known for its ability to
10 reorganize easily upon adsorption of different species, minimizing its energy in the presence of
11 the adsorbates. This process is typically related to sufficient mobility of surface atoms, which is
12 thermally activated. The reorganization of the surface can result in the formation of new
13 crystallographic facets [17,18] or monoatomic steps, possibly leading to periodic modulation of
14 the copper-adsorbate interface [19–21].
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28 Theoretical methods provide powerful complementary tools for understanding physical and
29 chemical properties of nanostructured materials at an atomic level of detail. The few existing
30 density functional theory studies concerning tungsten oxide nanowires point to a strong influence
31 of the surface structure on the electronic density of states, the band structure and the band gap
32 width [22,23]. Theoretical studies also confirm the generally expected effect of oxygen vacancies
33 in tungsten oxide nanowires on the interaction with adsorbates, specifically NO_2 [24].
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42 **Experimental details and computational methods**

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44 The experiments took place in the Surface Science Laboratory in Prague. The tungsten oxide
45 layers were prepared on the surface of a Cu(110) single crystal supplied by MaTeck, with a
46 declared miscut less than 0.1° . The crystal was cleaned by cycles of Ar^+ ion bombardment and
47 annealing up to 600°C under ultra-high vacuum (UHV) conditions. A thin tungsten oxide layer
48 was deposited onto the clean Cu (110) surface at 400°C by evaporation of tungsten trioxide
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3 powder from an electron-heated evaporation source. The evaporation source allowed for
4 deposition of very clean material at low deposition rate ($2 \text{ \AA}/\text{min}$) necessary for the formation of
5 epitaxial layers. The deposition was carried out in a flux of reactive atomic oxygen produced by
6 thermal gas cracker (TC50) in order to prevent excessive reduction of the vaporized tungsten
7 trioxide. Oxygen partial pressure was maintained at $5 \cdot 10^{-5} \text{ Pa}$ during the deposition. The
8 thickness of the layer was 0.74 nm . It was determined via the attenuation of the Cu 2p signal of
9 the substrate measured by X-ray photoelectron spectroscopy (XPS), assuming a homogeneous
10 layer and disregarding elastic scattering. The inelastic mean free path of electrons in tungsten
11 trioxide used in the calculation was $d = 11.41 \text{ \AA}$, as obtained from the TPP-2M formula [25].
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24 Reflection high energy electron diffraction (RHEED), low energy electron diffraction (LEED)
25 and scanning tunneling microscopy (STM) were used to investigate the surface structure. The
26 chemical state and the electronic structure of the system were probed by means of X-ray
27 photoelectron spectroscopy and angle-resolved photoelectron spectroscopy (ARPES). The
28 samples were prepared and characterized in situ in the apparatus combining RHEED and XPS
29 methods. For characterization by STM, LEED and ARPES the samples were transferred through
30 the ambient air and recovered by annealing to $400 \text{ }^\circ\text{C}$ in $5 \cdot 10^{-5} \text{ Pa}$ of oxygen. This procedure
31 was carefully tested in order to make sure that it does not influence the structural and electronic
32 properties of the samples. Theoretical models were proposed based on density functional theory
33 (DFT) calculations addressing the geometry, the electronic structure and the electric charges of
34 the system.
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49 X-ray photoelectron spectroscopy of W 4f, O 1s and Cu 2p core levels was performed using X-
50 ray radiation with energy of 1486.6 eV (Al K_α). ARPES measurements were conducted with He
51 lamp providing radiation with energy of 21.2 eV (He I). Deconvolution of obtained data was
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3 performed using W 4f 5/2 and W 4f 7/2 doublet with a spin-orbit splitting of 2.15 eV. Spectral
4 lines were represented by pseudo-Voigt functions and a Shirley-type background was subtracted
5 from the spectra. STM of the tungsten oxide layers was performed using chemically etched
6 tungsten tips annealed in vacuum. The images were obtained probing occupied states of the
7 sample. DFT calculations were performed using the Vienna Ab-Initio Simulation Package
8 (VASP) [26–28], which employs a plane-wave basis set and a periodic supercell method.
9 Potentials within the projector augmented wave (PAW) method [29] and the gradient-corrected
10 functionals in the form of the generalized-gradient approximation (GGA) with the Perdew-
11 Burke-Ernzerhof (PBE) functional [30,31] were applied. We used 6-8 ML thick slabs of
12 sufficient lateral dimensions to accommodate the 6x2 reconstruction of the Cu surface and the
13 deposited tungsten oxide species. During optimization, the first two surface layers and deposited
14 species were allowed to relax, until a convergence of 1.0 meV for the total energy was reached.
15 In all cases, a set of 3×3×1 Monkhorst-Pack k-points was used to sample the Brillouin zone.
16 Density of states (DOS) curves were used to analyze the electronic structure of the system and
17 the electronic charges of atoms were computed using Bader analysis [32]. Simulated STM
18 images based on the DFT calculated charge density were used for evaluation of the models'
19 accuracy.
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42 **Results and discussion**

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44 Well-defined thin layers of tungsten oxide were prepared by thermal evaporation of WO₃
45 powder in reactive atmosphere of atomic oxygen. The Cu(110) single crystal substrate was kept
46 at 400 °C during the growth. Due to the combination of the elevated temperature and the
47 exposure to atomic oxygen the copper substrate exhibited an oxygen induced Cu(110)-c(6×2)-O
48 surface reconstruction [33].
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Electron diffraction indicates one-dimensional character of the obtained tungsten oxide structures. We observe a set of uniformly spaced parallel planes in reciprocal space. According to the kinematic diffraction theory [34] such diffraction pattern is characteristic for a set of parallel one-dimensional wires in the real space, which are perpendicular to the reciprocal planes and, in our case, parallel to the Cu[1-10] crystallographic direction. The distance between the reciprocal planes is given by a regular spacing of the scatterers along the wires. The observed RHEED and LEED patterns are shown in Fig. 1. We recognize two sets of diffraction maxima at different intensity levels belonging to the tungsten oxide structures. No spots or lines of the copper surface are visible. The more intense maxima reflect the distance between the scatterers of $(3.83 \pm 0.01) \text{ \AA}$. This distance is equal to the lattice parameter of tungsten oxide in the pseudo-cubic phase [35]. The less intense lines in between the stronger ones refer to a superstructure with double periodicity.

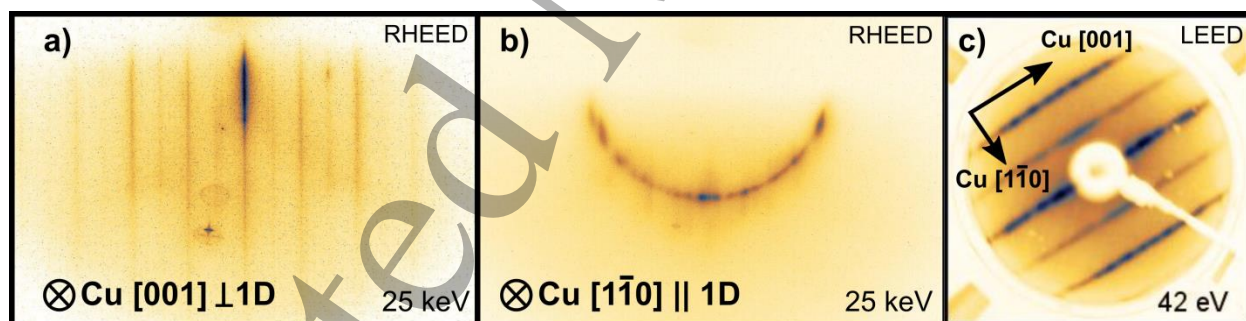


Figure 1. Diffraction patterns of the tungsten oxide thin film on Cu(110) surface. All visible patterns originate from the tungsten oxide overlayer, no diffraction spots or lines of the copper substrate are visible: (a) RHEED pattern in the direction perpendicular to 1D structures and (b) RHEED pattern in the parallel direction, both taken at primary energy $E_p = 25 \text{ keV}$; (c) LEED pattern ($E_p = 42 \text{ eV}$).

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3 STM results confirm the appearance of parallel one-dimensional structures on the surface. Fig.
4 2(a) shows an overview STM image revealing structures prolonged in the Cu[1-10] direction and
5 high steps perpendicular to the structures. The height of the steps is up to 3 nm according to the
6 height profile in Fig. 2(b). The step height corresponds neither to the thickness of the deposited
7 layer (0.74 nm) nor to the height of a monolayer step on the Cu(110) surface (1.28 Å). We
8 presume that the copper surface rearranged itself due to interaction with the tungsten oxide
9 overlayer. The schematic of the rearrangement is shown in Fig. 2(c). Considering the geometry
10 of the problem we can determine that the newly established high steps are formed by Cu(1-10)
11 facets, which are equivalent to the Cu(110) surface. In a high-resolution STM image (Fig. 2(d))
12 we can identify two different types of one-dimensional structures that are present on the Cu(110)
13 terraces between the high steps. The difference between the structures becomes apparent when
14 looking at averaged height profiles taken along the Cu[1-10] direction (Fig. 2(g)). Both profiles
15 are modulated with the main period equal to the distance of 3.83 Å, which corresponds to the
16 more intense lines in the diffraction patterns. However, one of the structures exhibits a
17 significantly deeper minimum every second period, which corresponds to the characteristic
18 length of 7.66 Å and therefore to the less intense maxima in the diffraction patterns. Structural
19 elements of the LEED pattern of the sample (Fig. 1(c)) are fully replicated in the 2D fast Fourier
20 transform of the STM image (Fig. 2(e)). To examine both structures further we also took
21 averaged height profiles along Cu[001] direction perpendicular to the one-dimensional
22 structures. Profiles taken in the areas with double periodicity (Fig. 2(h)) incorporate comparable
23 features. Also height profiles from areas with single periodicity (Fig. 2(i)) contain similar
24 sections in a particular order. However, without further analysis we cannot unambiguously
25 determine the actual lateral size of the tungsten oxide one-dimensional structures.
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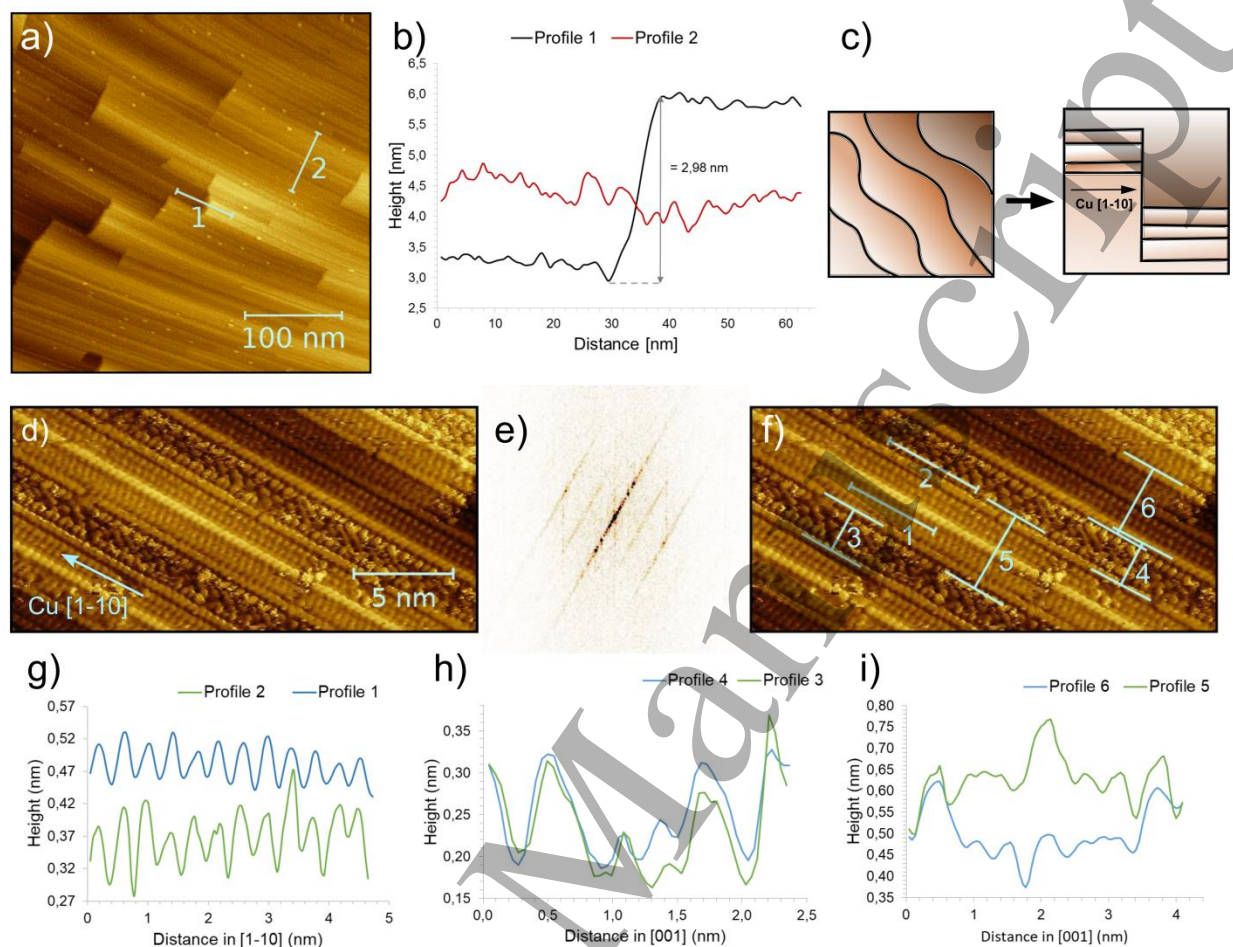
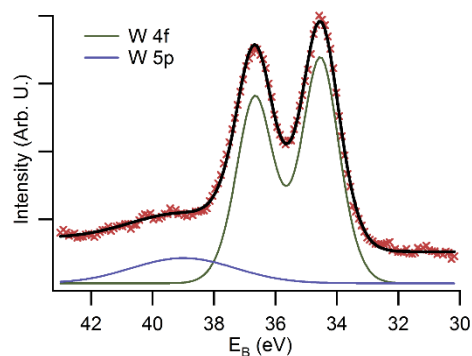


Figure 2. (a) Overview STM image of the $\text{WO}_3/\text{Cu}(110)$ system showing reorganization of the copper surface steps. (b) Height profiles taken from (a). (c) Schematic representation of the reorganization of the copper surface due to the presence of tungsten oxide thin film. (d) Detailed STM image of the tungsten oxide nanowires. (e) 2D Fast Fourier Transform (FFT) of the panel (d), reproducing the LEED pattern in Figure 1(c). (f) Areas of origin of height profiles shown in the following panels. (g) Averaged profiles (width 10 px) taken along Cu [1-10] direction. (h) and (i) Averaged profiles (width 74 px) taken along Cu [001] direction.

According to the X-ray photoelectron spectrum shown in Fig. 3 the film stoichiometry corresponds to tungsten trioxide. The W 4f doublet is positioned at 34.9 eV with spin-orbit splitting of 2.15 eV and intensity ratio of 0.75. As is typical for oxides, it is located at higher

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3 binding energy compared with the original metallic state, but its position is slightly shifted to a
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5 lower binding energy than what is usually reported for the W^{6+} state of the bulk material [36].
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20 **Figure 3.** X-ray photoelectron spectrum of W 4f core emission line.
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24 Reduced dimensionality of the prepared tungsten oxide structures is expected to give rise to
25 perturbations of their electronic structure. In order to confirm the effect we have followed the
26 valence band dispersion in directions parallel and perpendicular to the Cu [1-10] by ARPES. The
27 results are shown in Fig. 4. We observed a disparity in the character of the valence band in the
28 two directions. While a clear dispersing feature can be identified in the direction parallel to Cu[1-
29 10], all bands are completely flat in the perpendicular direction. The lack of dispersion in the
30 direction perpendicular to the Cu[1-10] is a footprint of the 1D character of the electronic states
31 of the tungsten oxide [3], allowing us to refer to the observed structure as nanowires both
32 structurally and electronically. Consequently, the reduced dimensionality of the electronic
33 structure implies that the interwire electronic coupling is so weak that it is essentially masked by
34 the intrinsic line width.
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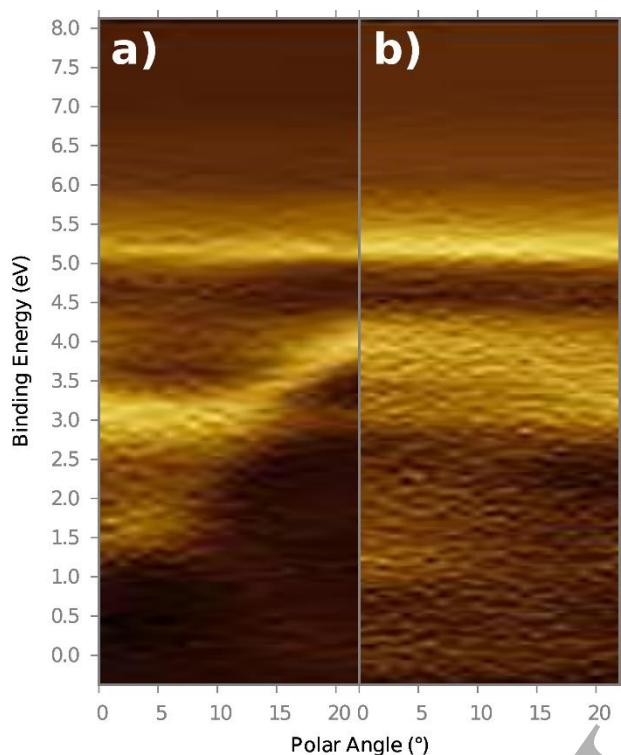


Figure 4. Angle resolved photoemission spectroscopy showing band structure of tungsten oxide nanowires measured along Cu[1-10] direction (a) and Cu[001] direction (b), which are directions parallel and perpendicular with respect to the nanowires, respectively. Measurements were performed at the photon energy of 21.2 eV. Visibility of the bands is enhanced by the first derivative over energy.

In order to learn more about the structures and epitaxy of the tungsten oxide nanowires on Cu(110) surface two theoretical models were considered - a nanowire composed of one line of tungsten oxide octahedrons (one line model) and a nanowire composed of two lines of tungsten oxide octahedrons (two line model), both supported on a Cu(110)-c(6x2)-O reconstructed surface [33] and oriented along Cu[1-10] direction. For the one line model, the copper surface was modelled by an eight layers slab with a lateral size of 6x2 the Cu(110) surface unit cell. It resulted in a supercell of about 7.3 Å by 15.4 Å and around 10 Å thick, including the eight oxygen atoms and two copper adatoms on the surface. The vacuum spacing between two

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3 repeated slabs was 10 Å. For the two lines model we used a six layers slab with a lateral size of
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5 6x4 the Cu(110) surface unit cell, resulting in a supercell of about 14.6 Å by 15.4 Å with
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7 thickness of 7.2 Å, including 16 oxygen atoms and 4 copper adatoms belonging to the Cu(110)-
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9 c(6×2)-O reconstruction.

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12 The resulting equilibrium configurations for both models are shown in Fig. 5. Tungsten atoms
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14 are bonded on top of the oxygen atoms belonging to the Cu(110)-c(6×2)-O reconstruction. Each
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16 tungsten atom is coordinated by six oxygen atoms, forming a tungsten oxide octahedron - the
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18 standard building block of tungsten oxide [37,38]. Oxygen atoms of the Cu-(110)-c(6×2)-O
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20 reconstruction are part of the octahedrons, forming their bottom apex (below W atom, marked
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22 O_{below} in Fig 5). Further oxygen atoms on the apices of the octahedrons are marked O_{aligned} -
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24 forming the W-O-W-O chain along Cu[1-10] direction, O_{lateral} - bonded to the sides of W atoms
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26 and O_{above} - bonded on top of the W atoms (Fig. 5). The W-O-W-O chains along Cu[1-10]
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28 direction exhibit a zigzag lateral distortion, more pronounced for the one line model than for the
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30 two lines model. Thus, the distance between the tungsten atoms is on average 3.84 Å for the two
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32 lines model and 3.91 Å for the one line model. The lateral displacement of tungsten atoms
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34 perpendicular to the Cu[1-10] direction is approximately 0.17 Å for the one line model versus
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36 0.07 Å for the two lines model. The average width of the wire is 3.39 Å and 7.04 Å for the one
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38 and the two lines models, respectively, while their heights are 4.32 Å and 4.54 Å.

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41 The distance between the W atoms along Cu[1-10] direction corresponds to the main
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43 diffraction lines in the RHEED and LEED patterns (Fig. 1). The lateral distortions of the W-O-
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45 W-O chains represent a plausible source of the observed diffraction maxima belonging to the
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47 surface structure with double periodicity. Considering that the lateral distortions are more
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3 prominent in the case of one line model, we propose that the structure with double periodicity
4 represents an early stage of the nanowire's growth, which corresponds to the one line model.
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8 Electronic charges of the atoms were computed by Bader analysis and they are comparable for
9 both models. In the case of the one line model the tungsten atoms present an average charge of
10 4.427 e with an average fluctuation of 0.057 e, for the two lines model the average charge is
11 4.430 e and the fluctuations are smaller - 0.024 e. We observed that electronic charge transfer
12 from the copper surface to the nanowires is lower and electronic charge transfer to oxygen atoms
13 bonded with tungsten atoms is higher for the two lines model. Generally, tungsten atoms in the
14 two lines model show a tendency to have a decrease in their electron density compared with the
15 one line model, although the difference is negligible. To consider the possible effect of these
16 results on the position of photoemission lines we can compare them with values of charge
17 calculated for a tungsten atom in bulk oxides, which is 4.654 e for monoclinic WO_3 and 2.971 e
18 for rutile WO_2 . The charge of the tungsten in our models is slightly lower but still closer to the
19 value for the bulk W^{6+} than bulk W^{4+} ions. Thus we can assume that tungsten atoms in the
20 observed tungsten oxide nanowires possess the W^{6+} oxidation state and the measured shift in the
21 position of the W 4f core line originates from the reduced dimensionality of the structure.
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40 In order to analyze the electronic structure, projected density of states (PDOS) curves were
41 computed. PDOS of a tungsten atom is shown in Fig. 5(e)). It exhibits two bands below the
42 Fermi level and these are shifted to lower energies in the case of the two lines model. The shape
43 of the higher energy band for the two lines model resembles the shape of the bulk WO_3 PDOS
44 curve more closely than the one line model. In case of PDOS of oxygen atoms shown in Fig. 5(f)
45 two kinds of oxygen atoms were considered: laterally bonded oxygen $\text{O}_{\text{lateral}}$ and oxygen atom
46 above the tungsten atom O_{above} . There are two oxygen bands below the Fermi level and again for
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the two lines model these bands are shifted to lower energies. Due to these shifts we can conclude that the two lines model corresponds to a configuration that is more stable.

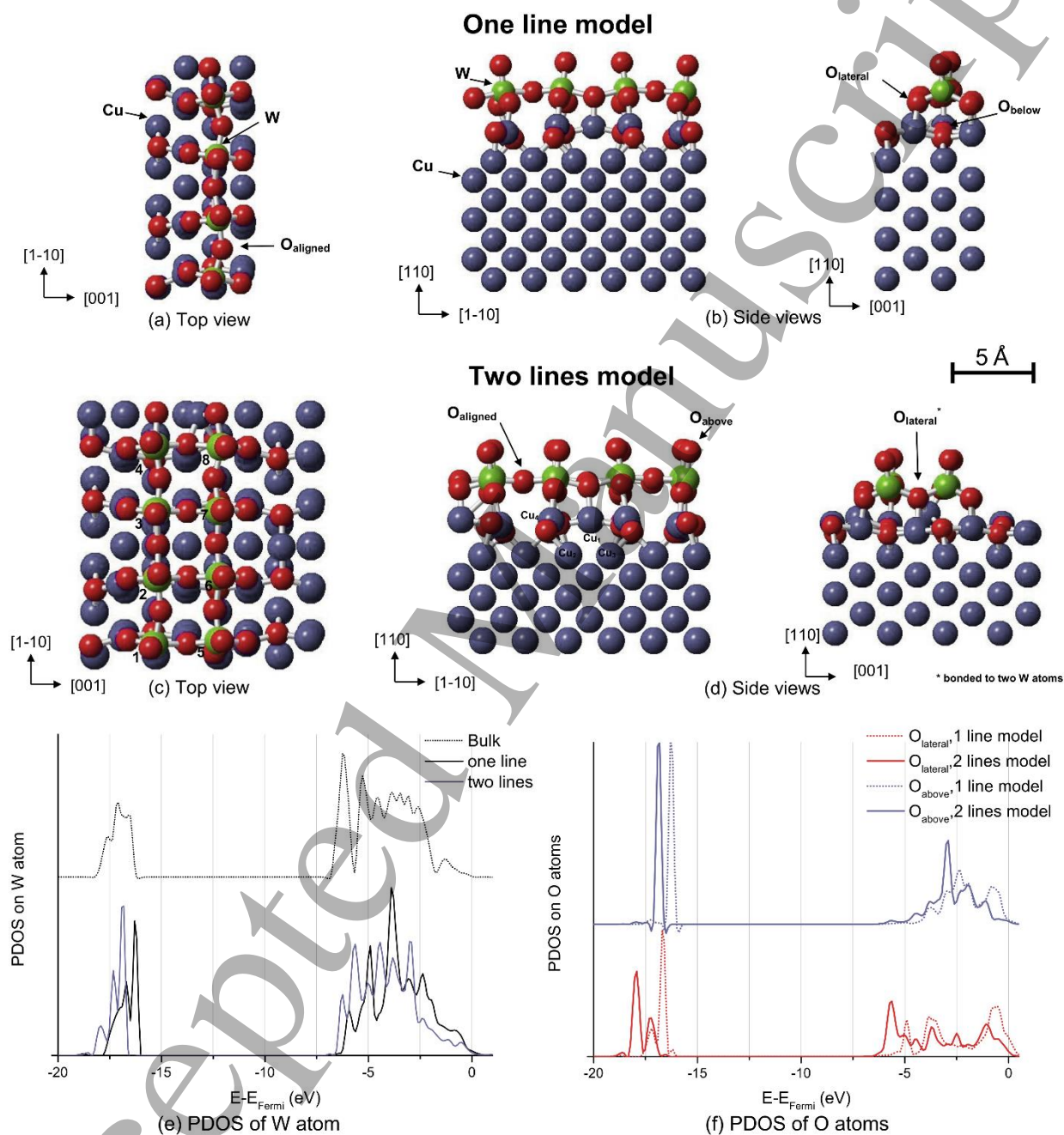
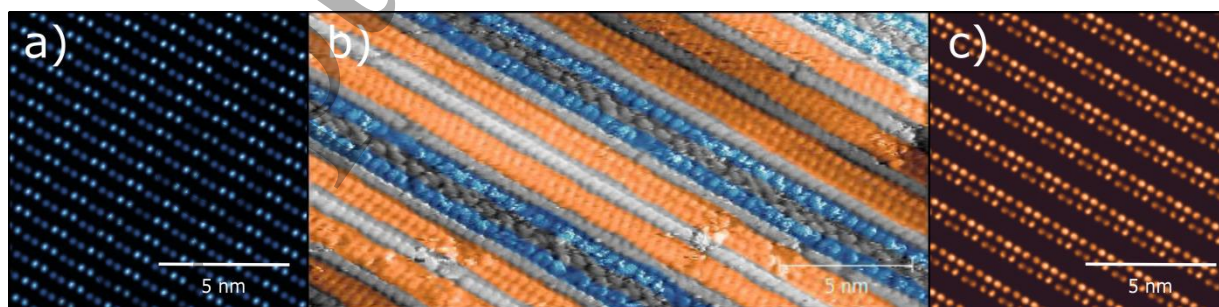


Figure 5. The lowest energy structures calculated by DFT: one line model (a), (b) and two lines model (c), (d) (Cu - blue, O - red, W - green). PDOS curves for different W (e) and O atoms (f).

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3 Simulated STM images (Fig. 6(a) and (c)) obtained for both models represent electron states
4 from -1 eV up to the Fermi energy. A value of the constant charge density corresponding to a
5 height of 6 Å above the surface was used. Comparing the calculated STM images for the one line
6 and two lines models with experimental STM image (Fig. 6(b)) we can propose that both the one
7 line and two lines features are present on the surface, alternating with oxidized Cu(110) surface
8 exhibiting the $c(6 \times 2)$ -O reconstruction. One line features and $c(6 \times 2)$ -O like regions in between
9 them exhibit the double periodicity of 7.66 Å along Cu[1-10] direction (Fig. 2(g), bottom
10 profile). This interpretation supports our hypothesis based on the shape of calculated structures
11 that the one line model represents the areas with double periodicity. Two lines features exhibit
12 the single periodicity of 3.83 Å along Cu[1-10] direction (Fig. 2(g), top profile). One line
13 features are appearing in pairs bordering the $c(6 \times 2)$ -O like regions which is clearly reflected by
14 height profiles taken in Cu[001] direction in Fig. 2(h). Two lines features are also appearing in
15 pairs approximately 7 Å apart (Fig. 2(i)). Linear protrusions or depressions along Cu[1-10]
16 direction that separate the one line and two lines regions or two lines and two lines regions
17 remain unidentified. These can be related to the rearrangement of the Cu(110) substrate (Fig.
18 2(a)) and represent Cu adatom rows on the Cu(110) substrate separating the tungsten oxide
19 nanowires.
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53 **Figure 6.** Simulated STM images of one line model (a) and two lines model (c) compared with
54 experimental results (b) where the corresponding areas are marked by matching colors.
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Origin of the formation of one-dimensional structures presumably lies in the Cu(110)-c(6x2)-O reconstruction of the surface since it is the most prominent difference between our experiment and the one carried out by Denk et al. [16]. They deposited WO₃ on the Cu(110)-(2x1)-O surface at low temperature (< 100 K) and after heating up to 600 K they obtained two-dimensional copper tungstenate. Closer inspection of our DFT models also supports this proposal because the nanowires grow directly on top of the oxygen atoms of the Cu(110)-c(6x2)-O reconstruction.

Conclusions

We have demonstrated the ability to prepare tungsten oxide model system by thermal evaporation of tungsten trioxide powder in highly reactive atmosphere of atomic oxygen. The obtained epitaxial tungsten oxide nanowires were comprehensively characterized by the means of experimental surface science and DFT calculations. One-dimensionality of the system was confirmed by RHEED, LEED, STM and DFT study of the structure. We identified unique electronic effects related to the system's one-dimensionality – loss of dispersive valence band features in the direction perpendicular to the structures and chemical shift of the W 4f core emission line, which was clarified by the results of DFT calculations.

The characterized system can be used as an exemplary inverse model system for investigations of chemical properties and surface reactivity, e.g. in the NO₂ sensing applications of copper-doped tungsten oxide [39], or it can serve as a template for self-organized growth of novel one-dimensional systems.

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