

Continuous dynamic recrystallization in a Zn-Cu-Ti sheet subjected to bilinear tensile strain

M. Leonard^{1,2,*}, C. Moussa³, A. Roatta^{1,2}, A. Seret⁴, J. W. Signorelli^{1,2}

¹ Instituto de Física Rosario (IFIR-CONICET/UNR), Ocampo y Esmeralda, Rosario (2000), Santa Fe, Argentina.

² Facultad de Cs. Exactas, Ingeniería y Agrimensura (FCEIA-UNR), Av. Pellegrini 250, (2000) Rosario, Argentina.

³ MINES ParisTech, PSL - Research University, CEMEF, CNRS UMR 7635, CS 10207 1 rue Claude Daunesse, 06904 Sophia Antipolis Cedex, France.

⁴ Technical University of Denmark, Department of Mechanical Engineering, Section of Manufacturing Engineering, 2800 Kgs. Lyngby, Denmark

*Corresponding author: leonard@ifir-conicet.gov.ar

Abstract

Research on zinc sheet formability, relevant for its technological applications, requires deeper understanding of the microstructural features that govern the plastic response of the material. In this work, a microstructural analysis by means of electron back-scattered diffraction (EBSD) technique was conducted on a commercial Zn-Cu-Ti sheet subjected to a bilinear tensile test at room temperature and low strain rate ($5 \times 10^{-4} \text{ s}^{-1}$). The refinement of the granular structure is analyzed in terms of the development of subgrains within initially large grains, which eventually evolve into high angle boundary grains. This continuous dynamic recrystallization (CDRX) mechanism appears as a key factor in order to explain the grain fragmentation process and the weakening of the texture observed during straining of this alloy.

Keywords: zinc sheet; continuous dynamic recrystallization; formability; texture.

1. Introduction

Zinc based alloys in the form of sheet have interesting applications in the architectural and building industry due to their good malleability and weldability, and their excellent corrosion resistance and surface aspect [1]. Early development of rolled zinc products included several alloys with small additions of copper, magnesium and/or aluminum to impart higher strength by means of solute strengthening or grain refining. The addition of titanium, in particular, remarkably enhances creep resistance to wrought zinc or Zn-Cu alloys, so the Zn-Cu-Ti alloy is now well established and it is the most widely used in sheet form for building applications such as facades, roofs or gutters [1,2].

31 Formability of zinc sheet is regularly good but it has some limitations. The most relevant drawback is
32 the high anisotropy exhibited by the material due to the sharp crystallographic texture inherited from the
33 rolling process. This strong texture is developed during deformation processing because of the reduced
34 number of slip systems available in zinc, which has a hexagonal close-packed (HCP) crystal structure with c/a
35 ratio of 1.856 [3]. Philippe et al. [4] and Funderberger et al. [5] studied and simulated by means of the Taylor
36 model [6] the texture evolution of several HCP materials, including a Zn-Cu-Ti alloy, relating the texture to the
37 basic mechanical properties of the sheets (i.e., tensile strength and plastic anisotropy). Recent investigations
38 dealing with experimental and modeling of zinc alloys formability have focused mostly on the anisotropic
39 plastic behavior and its relationship with the forming limit strains. When a phenomenological approach has
40 been used [7,8] the macroscopic behavior was captured by a Hill 90 yield criterion [9] in conjunction with a
41 hardening power law which accounts for the temperature and strain-rate sensibility of the material. On a
42 crystal plasticity approach, Schwindt et al. [10] improved the combination of the Marciniak-Kuczynski (MK)
43 technique with the visco-plastic self-consistent (VPSC) model proposed originally by Signorelli et al. [11] to
44 predict the forming limits of a Zn-Cu-Ti sheet with relatively good results. However, the model
45 underestimates the forming limit curve (FLC) when the major strains are parallel to the rolling direction
46 (RD). Cauvin et al. [12] also applied the VPSC model to analyze the mechanical anisotropy of a Zn-Cu-Ti alloy,
47 focusing mainly on the tensile response and the texture evolution. Even though the crystal plasticity
48 approaches are based on the deformation mechanisms operative at the grain scale, the works mentioned are
49 based mostly on the texture evolution but not on topological aspects of the microstructure or dislocation-
50 based hardening laws. Borodachenkova et al. [13] applied a microstructure-based hardening model
51 embedded in the VPSC scheme to simulate texture evolution and mechanical response of a Zn sheet subjected
52 to forward-reverse simple shear deformation. More recently, Schlosser et al. [14] experimentally analyzed the
53 influence of bilinear strain paths (i.e., balanced biaxial expansion followed by uniaxial tension) on the
54 formability of a Zn-Cu-Ti sheet. The extended, somewhat unexpected ductility encountered lead to the
55 hypothesis that for an accurate modeling of zinc sheet formability it would be necessary to take into account
56 not only texture but also the microstructure evolution in terms of grain fragmentation. In fact, microscopical
57 evidence revealed a grain-refining process which cannot be entirely explained by slip activity and may be due
58 to dynamic recrystallization [14].

59 Recrystallization processes which develop during deformation are generally categorized as dynamic
60 recrystallization (DRX). These phenomena consist on the creation of new grains in the microstructure of a
61 polycrystalline material subjected to plastic straining, reflecting a combination of thermally-activated and

62 strain-induced mechanisms different from the typical static recrystallization process (SRX) which occurs
63 during annealing. Depending on material properties and deformation conditions, DRX may take place as a
64 discontinuous nucleation and growth process, which is known as discontinuous dynamic recrystallization
65 (DDRX); or it may occur in a continuous fashion as a gradual creation of new grains, in which case it is termed
66 continuous dynamic recrystallization (CDRX) [15,16]. Furthermore, it is possible that during straining a
67 dynamic recovery process takes place (DRV), either as a preceding step to CDRX or preventing DDRX [16].

68 Mechanical response and microstructure evolution during warm or hot working of HCP materials are
69 not as widely studied as in the case of cubic metals, especially regarding sheet forming. Most of the past and
70 current literature deals with magnesium alloys, for which formability can be enhanced by DRX during
71 warm/hot forming operations [17–19]. Mg and its alloys are known to undergo CDRX at low to intermediate
72 temperature deformation through a mechanism of gradual lattice rotation near the original grain boundaries
73 [20,21]. This mechanism is associated to plastic inhomogeneity due to the lack of sufficient slip systems in the
74 low-symmetry, hexagonal crystal structure. Thanks to the advances in the electron back-scattered diffraction
75 (EBSD) technique much research has been devoted during the last decades to understand the mechanisms
76 responsible for CDRX in Mg alloys by means of orientation imaging microscopy (e.g., refs. [22–24]).
77 Regarding zinc, it is extensively reported in the technical literature [1,2] that, owing to its low melting point
78 ($T_m = 693$ K), recrystallization can occur at regular processing temperatures (room temperature
79 corresponding to 293 K, i.e., $\sim 0.42T_m$). Little work hardening can be obtained for this reason, even for alloys
80 containing solute additions which tend to elevate the recrystallization temperature [2]. However, only a few
81 systematic studies can be found (most of them previous to the development of the EBSD technique) and the
82 results concerning deformation mechanisms are not always concurrent. For instance, Neumeier and Risbeck
83 [25] studied the creep behavior of Zn-Cu and Zn-Cu-Ti rolled slabs with varying composition and processing
84 temperatures. They found a very fine-grained microstructure for the Ti-containing alloys rolled at room
85 temperature and this was attributed to a non-conventional recrystallization mechanism. A Zn-0.4%Al alloy
86 was found to exhibit room temperature superplasticity by Naziri and Pearce [26], a behavior which was
87 explained by a combination of DRV and grain boundary sliding (GBS). Malin et al. [27] reported DRX in pure
88 zinc rolled at room temperature up to moderately high strains. During the early stages of rolling, deformation
89 was accommodated mostly by twin and slip, but at a strain of ~ 0.5 the deformation localized in shear bands
90 consisting of recrystallized grains. With increasing strain, these DRX bands covered the whole volume of the
91 samples, allowing further deformation by means of renewed slip and twinning activity. Solas et al. [28]
92 developed an N -site VPSC model accounting for local interactions with the aim of simulating plastic

93 deformation and recrystallization. The model was applied to commercial purity zinc (CP-Zn) deformed by
94 channel-die compression at 398 K ($\sim 0.57T_m$); at these conditions the material showed neither DRX nor
95 twinning. Besides these rather divergent observations, most of the works mentioned focus mainly on the
96 metallurgical aspects of the granular evolution of low-alloyed zinc and its relationship to the mechanical
97 response, but not to their influence on anisotropy of the sheet products and their formability.

98 The aim of the present work is to investigate the microstructure evolution of a Zn-Cu-Ti sheet
99 subjected to a bilinear tensile test by means of orientation imaging microscopy. The main features of this
100 evolution are analyzed in terms of the grain size and morphology, as well as the development of
101 substructures which account for a grain fragmentation process that shall be classified as CDRX. These results
102 are complemented with an analysis of the orientation influence on the fragmentation behavior and a study on
103 the boundary character of the subgrains, together with an evaluation of the effect of a higher strain rate.

105 2. Material and methods

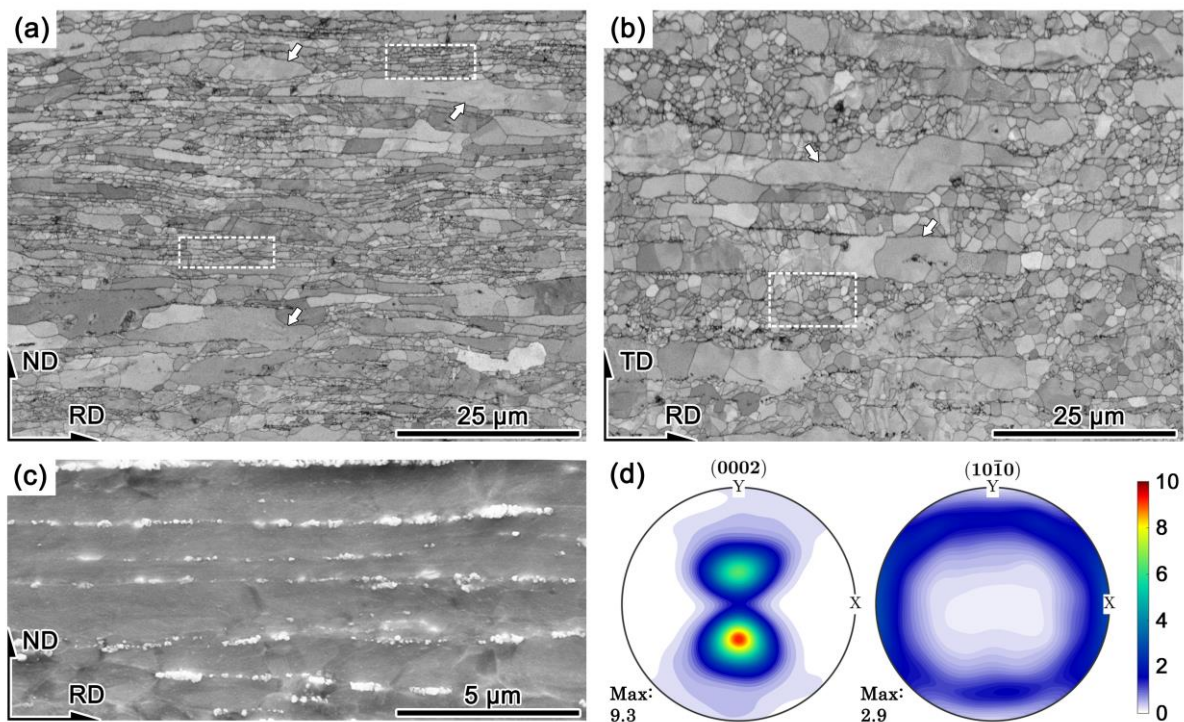
106 The material chosen for this study is a commercial Zn-Cu-Ti sheet of 0.65 mm thick. The manufacturing
107 process consists of continuous casting followed by rolling in two steps, the first one comprising two passes at
108 high temperature. Then the sheet is coiled and the final rolling step consists on several passes near room
109 temperature ($\sim 50\text{--}80\text{ }^\circ\text{C}$) until the final thickness is achieved. The chemical composition contains 0.094 Cu
110 and 0.047 Ti (wt. %).

111 Orientation imaging analysis was carried out by the EBSD technique, using a Bruker Quantax EBSD
112 system (comprising an e-Flash^{HR} detector and Esprit software) mounted on a Zeiss Supra40 field emission
113 gun scanning electron microscope (FEG-SEM) operated at 20 keV. Maps having a size of 1200×900 data
114 points were acquired over a square grid with a step of 94 nm (total map size: 122.8×84.6 μm). Post-
115 processing of the EBSD data was done with the MTEX toolbox [29]. The misorientation angle to define a high-
116 angle grain boundary was set as $\theta > 12^\circ$ (i.e., HABs). Boundaries with misorientation angles between 2° and
117 12° were considered as low-angle boundaries (LABs). Further details regarding the post-processing of the
118 EBSD data are stated below. Pole figures for texture analysis were derived from the orientation distribution
119 function (ODF) calculated from the raw EBSD data. This texture representation, for the size of the EBSD maps
120 chosen, was found to be statistically reliable when compared to macroscopic textures measured on the as-
121 received material by means of the X ray diffraction (XRD) technique. Moreover, the as-received texture shown
122 below is in good agreement with the texture obtained by XRD for a different sheet of the same alloy [30]. The

123 samples were mounted on a cold-setting epoxy resin, polished with SiC abrasive paper down to P2400 grit
124 size (FEPA standard), then with 3 μm and 1 μm diamond suspension and finished with colloidal silica of
125 $\sim 0.05 \mu\text{m}$ average particle size.

126 The microstructure of the as-received Zn-Cu-Ti sheet revealed by EBSD is presented in Fig. 1, showing
127 two band-contrast (BC) maps corresponding to the longitudinal (RD-ND) and in-plane (RD-TD) sections of
128 the sheet (ND and TD refer to the normal and transverse directions with respect to the rolling direction, RD,
129 respectively).

130



131

132 **Figure 1.** Microstructure and crystallographic texture of the Zn-Cu-Ti sheet in the as-received condition. (a)
133 and (b): band contrast of the EBSD maps taken from the RD-ND and RD-TD sections, respectively. Some large,
134 elongated grains are indicated by arrows and clusters of small, equiaxed grains are bounded in dashed boxes.
135 (c) Secondary electron micrograph revealing the TiZn₁₆ phase. (d) Basal and prismatic pole figures
136 recalculated from the EBSD data (X and Y labels correspond to TD and RD, respectively).

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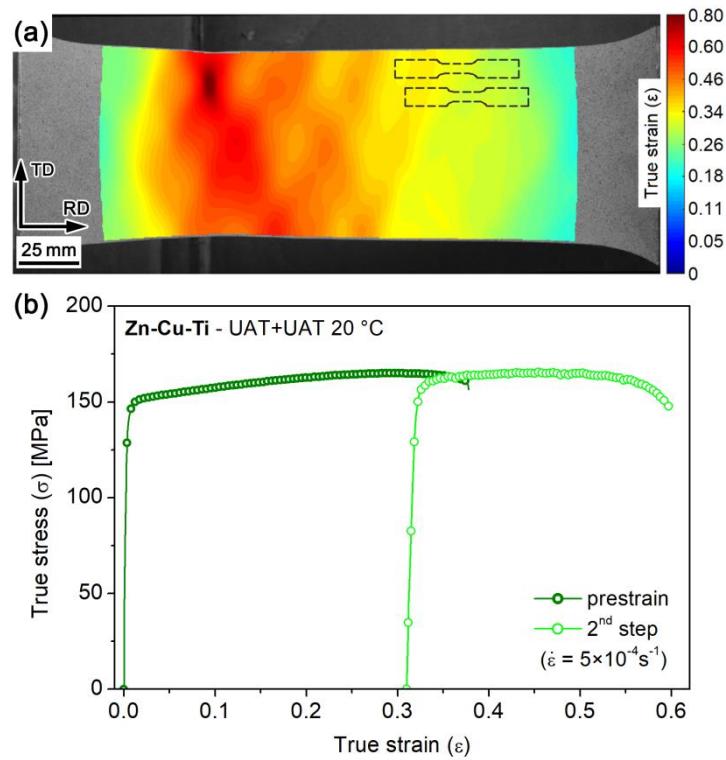
138 The morphology revealed by the BC maps presents a nearly bimodal distribution of grains, with a
139 population of large, approximately pancaked grains, and a different population of smaller, more equiaxed
140 grains (grain reconstruction is presented in section 3.2). Cu remains in solid solution for the composition
141 reported [2]. The low solubility of Ti in Zn results in the formation of the intermetallic compound TiZn₁₆ [31],
142 which is hardly seen in the BC maps because of the small size of their particles. These are shown in more

143 detail in Fig. 1c by means of a secondary electron scan with higher magnification. The weight fraction of the
144 intermetallic phase in the present alloy has been estimated at ~3% maximum [30]. It can be seen that the
145 TiZn_{16} particles are sub-micron sized and they form thin bands along the RD. This is the sheet microstructure
146 inherited from the rolling process, which also results in the crystallographic texture shown by the basal
147 (0002) and prismatic $\{10\bar{1}0\}$ pole figures in Fig. 1d. The main texture component consists of a majority of the
148 grains having their c -axis pointing at $\sim 23^\circ$ of the normal direction in the ND-RD plane (the angle estimation
149 has been done by calculating the orientations of the maxima of the ODF). It can also be seen that there is a
150 subtle preferred orientation of the prismatic poles towards the TD. Such texture is typical for a rolled HCP
151 aggregate with $c/a > 1.633$, resultant of $\langle a \rangle$ basal slip $((0001)\langle 11\bar{2}0 \rangle)$ accompanied by minor $\langle c+a \rangle$ gliding on
152 the 2nd order pyramidal plane (i.e., $\{11\bar{2}2\}\langle 11\bar{2}3 \rangle$ system) [3,4].

153 The analysis of the strain effect on the microstructure will be limited to uniaxial tensile (UAT)
154 deformation along the RD in two steps. A non-standard sample (140×85 mm) was subjected to tensile loading
155 along RD up to an early stage of localized necking, which corresponds to an average major true strain of
156 ~ 0.35 (Fig. 2a). Then a smaller tensile sample (11×4 mm) was cut from a region which locally attained a
157 prestrain of ~ 0.31 , with its longitudinal axis also parallel to RD. The second test was conducted until fracture,
158 giving a total accumulated major strain in the EBSD-analyzed zone of ~ 0.66 . This two-step, bilinear tensile
159 test (referred to as UAT+UAT) allowed for achieving higher deformation than with a regular tensile test, since
160 the geometric effect of the diffuse neck is eliminated when cutting the smaller sample. Figure 2b shows the
161 flow curves of the UAT+UAT test. It is worth noting that repeated tests gave similar results, with a maximum
162 deviation in the second step of $\sim 10\%$ in flow stress and $\sim 25\%$ in total elongation (the relatively high spread
163 in the latter is expected for this material [32]).

164 All tests were carried out with an Instron 3382 universal testing machine at room temperature (~ 20
165 $^\circ\text{C}$), applying a constant engineering strain rate of $5 \times 10^{-4} \text{ s}^{-1}$. A second-step sample was deformed at $5 \times 10^{-1} \text{ s}^{-1}$
166 to test the effect of a high strain rate; the results are shown and discussed in Sec. 3.6. The large tensile sample
167 was cut from the as-received sheet by a laser cutting machine; the smaller samples were extracted by means
168 of electro-discharge machining (EDM). Both cutting processes were accomplished without significant heating
169 of the material. The strains developed during the tests were determined by the digital image correlation (DIC)
170 technique, as detailed in Schlosser et al. [14].

171



172

173 **Figure 2.** Bilinear tensile test (UAT+UAT) of the Zn-Cu-Ti sheet. (a) Sample pre-strained along RD with the
 174 major true strain field measured by digital image correlation. Geometry of the second UAT samples are
 175 shown in dashed lines. (b) Corresponding flow curves. The prestrain value of the second-step samples is
 176 ~ 0.31 .

177

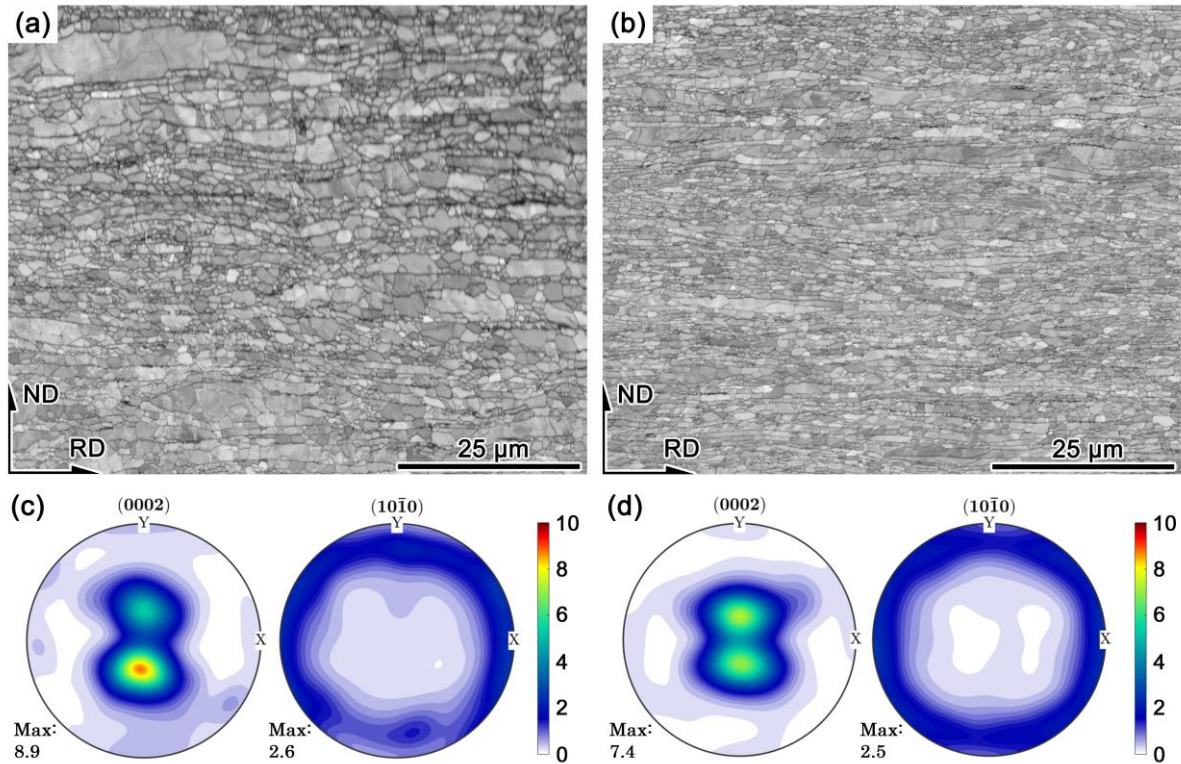
178 The EBSD scans for the deformed material were done over the RD-ND section of the samples, in a
 179 region between the middle plane of the sheet and the surface.

180

181 3. Results and discussion

182 3.1. Basic grain morphology and texture evolution

183 We take the BC data from the EBSD maps as a first approach to the grain morphology evolution (Fig.
 184 3). It is apparent from these maps that the grain size tends to decrease with increasing strain level. Compared
 185 to the as-received condition (Fig. 1a), the UAT prestrain sample (Fig. 3a) seems to show fewer large grains.
 186 This trend is even more evident for the final stage of tensile loading ($\epsilon \sim 0.66$, Fig. 3b), in which there appears
 187 to be no grain with a size in the order of the initially large ones.



188

189 **Figure 3.** Band contrast maps and pole figures of the deformed Zn-Cu-Ti samples. (a) and (c): UAT prestrain
 190 ($\epsilon \sim 0.31$). (b) and (d): second UAT loading ($\epsilon \sim 0.66$). X and Y labels correspond to TD and RD, respectively.
 191

192

193 Crystallographic texture of the deformed material, represented by the basal and prismatic pole figures
 194 (Figs. 3c and 3d), shows a modest decrease in intensity together with a rotation of the *c*-axis lying on the RD-
 195 ND plane towards the ND. The angular deviation between the maximum intensities and the ND changes from
 196 the initial 23° to 20° and 17° for the two successive tensile tests, respectively. The described evolution can be
 197 interpreted as a gradual rotation of the majority of the basal planes towards the plane of the sheet as a
 198 consequence of the basal slip activity, which is dominant in zinc [4,30]. However, the weakening of the main
 199 texture component should also be taken into consideration, as gliding alone on the basal system would not be
 200 enough to explain it. Moreover, there is a slight spread of the prismatic poles around ND at the end of the
 201 deformation process. The spreading of the main component and the rotation of the prismatic planes around
 202 ND can be related to CDRX, as it will be discussed below.

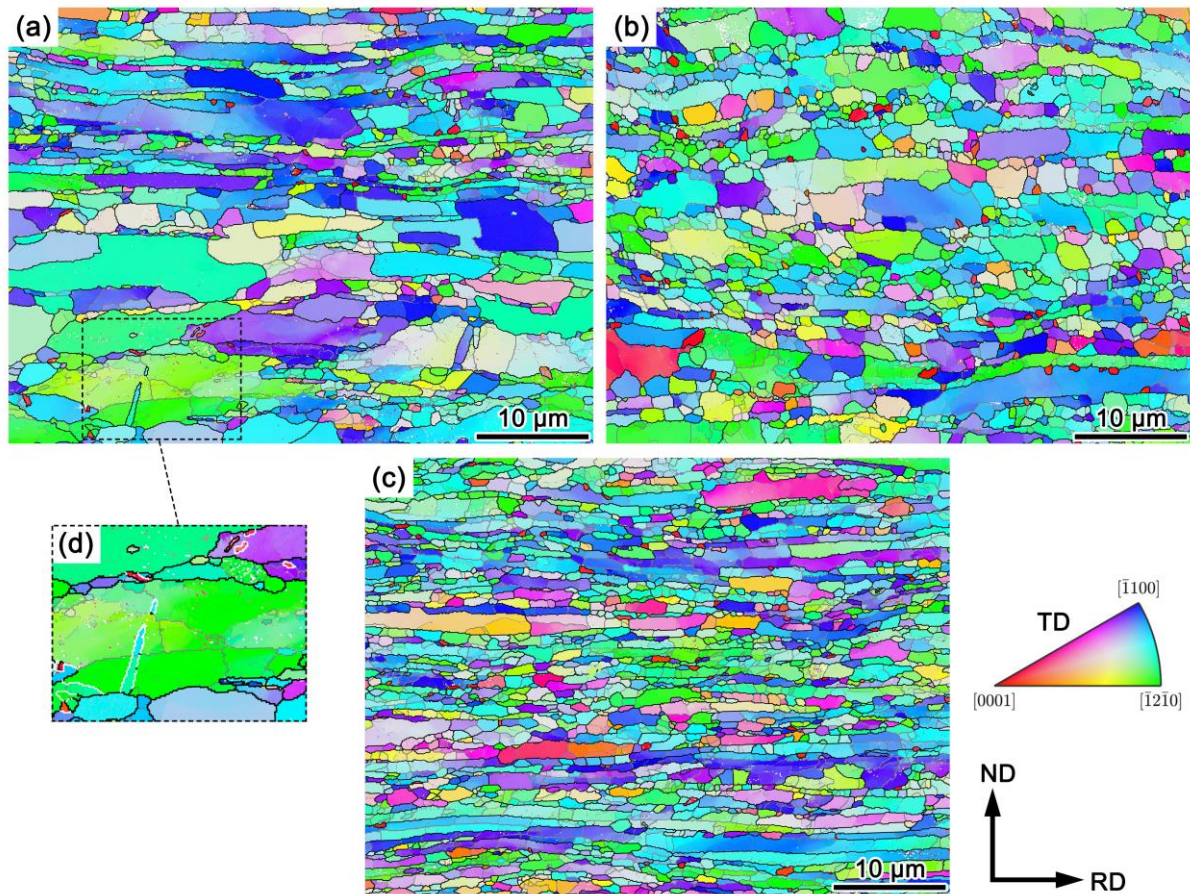
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203 3.2. Grain size and shape evolution

204 Figure 4 shows the reconstructed grain boundaries from the EBSD data plotted over the inverse pole
 205 figure (IPF) maps. For reasons of visual clarity only a detail of each map is shown, representative of the
 206 sample as a whole. Black lines correspond to HABs ($\theta > 12^\circ$) and gray lines to LABs ($2^\circ < \theta \leq 12^\circ$). Only few twins

207 can be seen in the microstructure. Considering zinc's compressive $\{10\bar{1}2\}\langle 10\bar{1}\bar{1}\rangle$ twinning system [33], twin
 208 boundaries were detected by the corresponding rotation angle and axis with a tolerance of 5° (ie.,
 209 $86.03^\circ\langle 11\bar{2}0\rangle \pm 5^\circ$). This computation gave a total twin boundary density of 0.060, 0.059 and 0.049 $\mu\text{m}/\mu\text{m}^2$
 210 for each strain level, which corresponds to only 4.5%, 4.1% and 2.5% of the total grain boundary density,
 211 respectively. A detail of the as-received IPF map is given in Fig. 4d in order to illustrate the detection of twin
 212 boundaries. Although twinning is usually a relevant deformation mechanism in HCP metals, the presence of
 213 Cu and Ti in zinc diminishes twinning activity, as observed by Philippe et al. [4] and Faur and Cosmeleată [34]
 214 for similar alloys.

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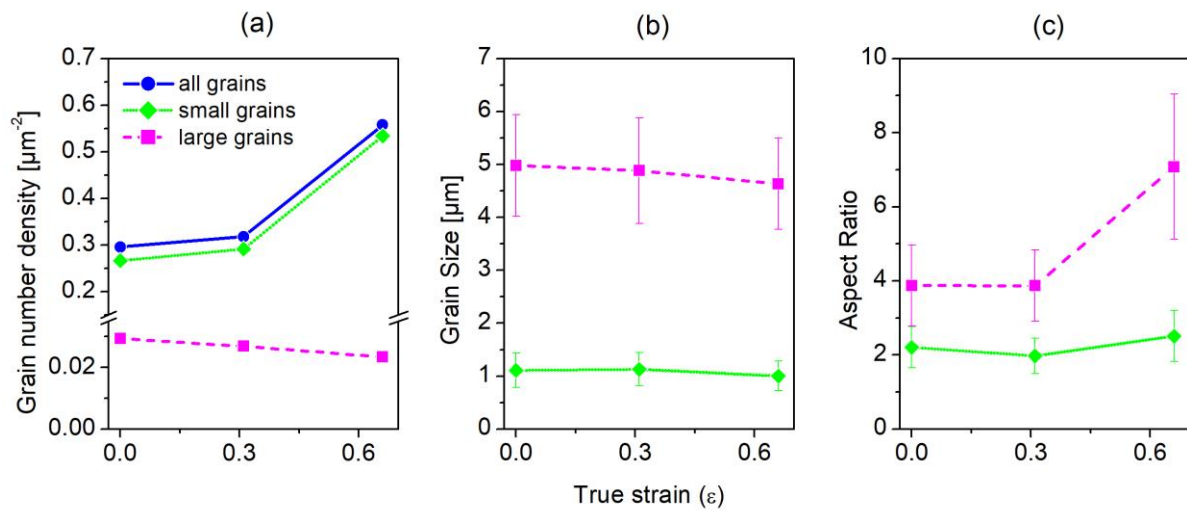
217 **Figure 4.** Grain boundaries plotted over orientation maps corresponding to the IPF with respect to the TD.
 218 (a) As-received; (b) UAT prestrain ($\varepsilon\sim 0.31$); (c) second UAT loading ($\varepsilon\sim 0.66$). Black lines: HABs ($\theta>12^\circ$); gray
 219 lines: LABs ($2^\circ<\theta\leq 12^\circ$). (d) Detail showing $\{10\bar{1}2\}\langle 10\bar{1}\bar{1}\rangle$ twin boundaries in white.

220

221 Statistics derived from the grain reconstruction are shown in Fig. 5. Grain sizes are reported as an
 222 equivalent diameter value, which corresponds to the diameter of a disc having the same area of the polygonal

223 shape of each reconstructed grain. This area-based approach allows to capture the shape effect when grains
 224 are not equiaxed, as in the present case. Grains with a diameter $\leq 0.33 \mu\text{m}$ (area ≤ 10 data points for the
 225 current step) were disregarded. For analysis purposes and taking into account the roughly bimodal
 226 distribution in the initial material, the set of grains in each sample was subdivided into two populations by a
 227 size criterion. All grains with an equivalent diameter greater than $3 \mu\text{m}$ were considered as large grains. The
 228 remainder grains belong to the set of small grains. Grain size and aspect ratio values reported in Fig. 5b and
 229 5c correspond to the arithmetic mean of each population, while the error bars correspond to the standard
 230 deviation.

231



232

233 **Figure 5.** Grain statistics of the Zn-Cu-Ti sheet microstructure for each strain level, subdividing the
 234 populations by a size criterion: (a) number density of reconstructed grains; (b) average grain sizes
 235 (calculated as equivalent circle diameter); (c) mean aspect ratios. Error bars correspond to the standard
 236 deviation of each population.

237

238 As can be seen in Fig. 5a, the increase of strain correlates with an increase of the amount of grains in
 239 the microstructure. Most of this increase is due to the fragmentation of large grains leading to the
 240 multiplication of small grains. Indeed, the decrease in the number density of large grains (from 0.029 to 0.023
 241 μm^{-2} , i.e., 25%) is accompanied by a slight reduction of their mean size (Fig. 5b). This double trend is related
 242 to the fact that the number increase of small grains accounts for a larger fraction of the map area: this fraction
 243 grows from 35% to 40% after the first UAT step, and up to 55% for the second test. It should be noted that
 244 the overall grain size of the specimens is relatively small, even for the grains considered large here.
 245 Particularly, for the set of smaller grains, their average value of $\sim 1 \mu\text{m}$ is below the range of the typical grain

246 sizes ($\sim 3\text{--}10\ \mu\text{m}$) achieved by means of standard thermomechanical processing [35]. Grain sizes of this order
247 are reported in the early literature of fine-grained zinc sheets produced by rolling near room temperature
248 [2,36]. Moreover, researchers seeking for superplasticity in Zn-0.4%Al alloy found an optimum grain size of
249 approximately $0.5\text{--}1\ \mu\text{m}$ without needing to apply severe plastic deformation processes [26,37].

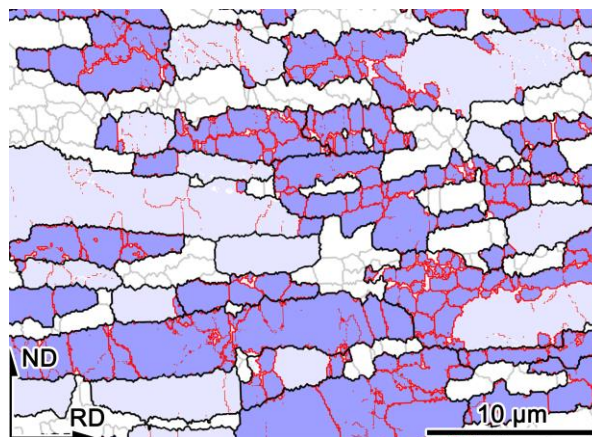
250 Regarding the shape of the grains, from Fig. 5c it can be appreciated that the smaller grains are slightly
251 elongated along the RD, but their aspect ratio stays almost constant throughout the whole deformation
252 process. On the other hand, the remaining large grains –particularly at the highest UAT strain– are much
253 more elongated along the RD. The stretching of grains along the direction of applied stress is typical for
254 polycrystals accommodating plastic deformation; however, as it will be discussed below, in the case of the Zn-
255 Cu-Ti sheet under study the elongation of the large grains is accompanied by their fragmentation.

256

257 3.3. Grain fragmentation by subgrain development

258 The LABs plotted on the maps of Fig. 4 allow visualizing the substructure development within the
259 larger grains as a consequence of the imposed deformation. The formation of boundaries of low
260 misorientation angles in the grain interiors is the typical microstructural feature of recovery or
261 recrystallization occurring during deformation processing, i.e., DRV/DRX [15,16]. In the present case,
262 observation of the microstructure reveals that the large grains suffer a subdivision or fragmentation process
263 by developing LABs (either as closed subgrains or as inner boundaries which appears to be a transient step
264 towards the formation of subgrains). A closer examination of this feature is presented in Fig. 6, in which a
265 detail of the UAT prestrain map is shown.

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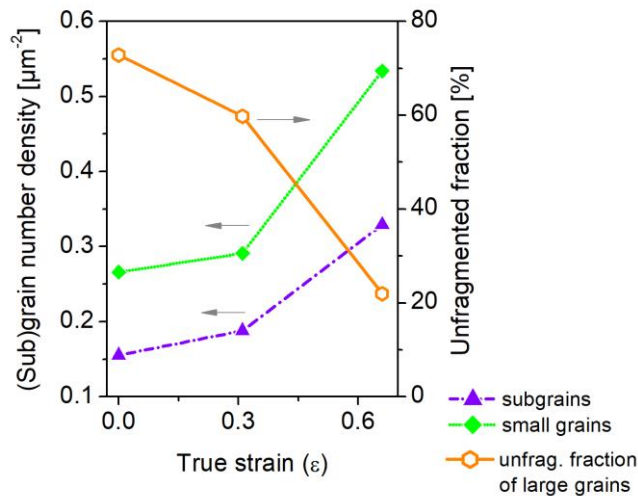


267

268 **Figure 6.** Map detail illustrating fragmentation of the larger grains ($\epsilon \sim 0.31$). Gray lines: HABs ($\theta > 12^\circ$) of
269 small grains ($\leq 3\ \mu\text{m}$); black lines: HABs of larger grains ($> 3\ \mu\text{m}$); thick red lines: LABs ($2 < \theta \leq 12^\circ$) of subgrains

270 within large grains; thin red lines: inner LABs. Blue: subgrains within large grains; light blue: large grains
 271 without subgrains or region of a large grain with an area greater than 50% of the whole grain area (i.e.,
 272 unfragmented parent grains).
 273

274 In order to analyze the general evolution of the subgrains with respect to their parent large grains,
 275 following Fig. 6 the subgrains or fragments are defined by the closed LABs ($2 < \theta \leq 12^\circ$) within each large grain.
 276 A large grain is considered unfragmented if it does not have any subgrain or if it has an unfragmented area
 277 greater than the half of the total grain area. As Fig. 7 displays, there is an important increase of the number of
 278 subgrains during straining. Also, the fraction of unfragmented large grains tends to decrease markedly with
 279 deformation, meaning that within the bigger grains there is a relative increase in subgrain formation. This
 280 observation is clearly consistent with a fragmentation process. Moreover, the shape and the average size of
 281 the subgrains (thick red boundaries in Fig. 6) throughout the deformation course are equivalent to those of
 282 the small grains (gray boundaries in Fig. 6). In addition, the number density of subgrains and small grains
 283 increase in a similar way, as reported in Fig. 7.
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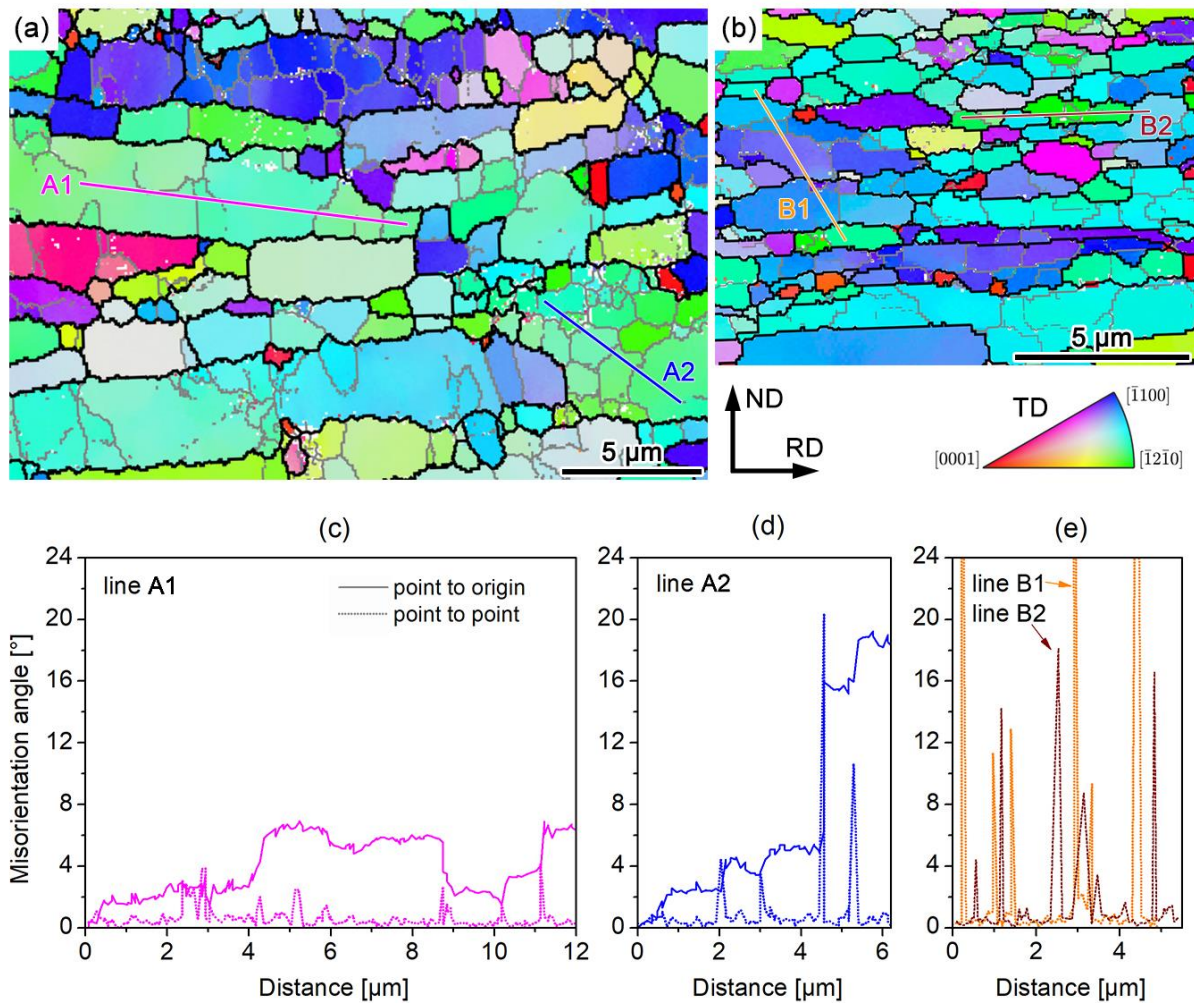


285
 286 **Figure 7.** Evolution of the populations of grains and subgrains, and of the fraction of large grains which are
 287 not fragmented (number density of small grains is repeated from Fig. 5 for reference).
 288

289 A closer inspection of the subgrains evolution can be realized by analyzing the misorientation gradient
 290 across the large grains, as shown in Fig. 8 where representative regions of the microstructure were selected.
 291 From the tensile prestrain (Fig. 8a) it is seen that many subgrains develop by progressive accumulation of
 292 misorientation from the center of the large parent grains to their boundaries, as displayed by the cumulative
 293 misorientation profiles over lines A1 and A2. The gradient over line A1 is smoother and the boundaries

294 involved have low misorientation angles, reflecting an early stage of the fragmentation process for the
 295 selected grain. Subgrains crossed by line A2 have boundaries with higher misorientation and account for a
 296 more pronounced misorientation accumulation from the center of the parent grain to its boundary. This
 297 evolution is usually reported as the mechanism by which Mg alloys undergo CDRX due to progressive lattice
 298 rotation [20,23,38].

299 The misorientation profiles over lines B1 and B2 (Fig. 8b and 8e), corresponding to the highest strain
 300 level, reflect the evolved stage of the fragmentation process, since most of the boundaries between subgrains
 301 have higher misorientation angles than the preceding case. Moreover, some other boundaries have already
 302 surpassed the 12° threshold, thus becoming the HAB of a new small recrystallized grain. Note that most of the
 303 large grains in the map detail of Fig. 8b have become completely fragmented, as LABs have developed over
 304 their entire volume.



305

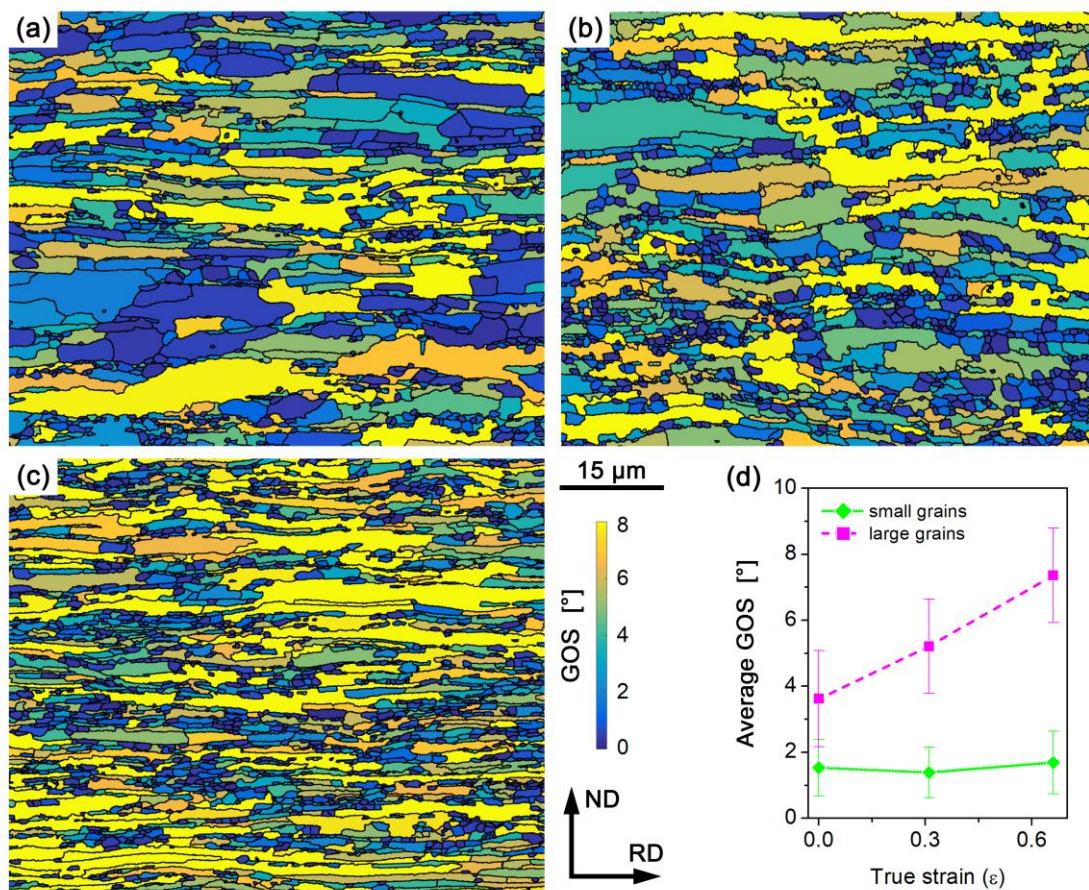
306 **Figure 8.** Misorientation gradients across fragmented grains: (a) and (b) EBSD map details of the first
 307 ($\epsilon \sim 0.31$) and the second ($\epsilon \sim 0.66$) UAT tests, respectively. (c–e) Changes of misorientation angle over the

308 lines indicated in the maps (solid lines: cumulative or point-to-origin misorientation; dotted lines: point-to-
309 point misorientation).

310
311 In summary, the analyzed evolution of the (sub)boundaries and the similarities between the size and
312 shape of the subgrains to those of the small grains, indicate a continuous fragmentation process which gives
313 place to subgrains that evolve to become small grains with HABs. Similar observations have been reported for
314 several Mg alloys which undergo CDRX at different deformation conditions [17,22,39,40]. Therefore, it is
315 concluded that the grain refinement process observed during the deformation of the present Zn-Cu-Ti alloy
316 can be classified as CDRX, with no significant influence of twinning.

317 Partitioning of the grain populations by size regarding their fragmentation behavior is consistent with
318 an analysis of their internal orientation deviations. The average values of the orientation spread by grain –i.e.,
319 grain orientation spread (GOS)– is a typical parameter used for identifying recrystallized grains in deformed
320 microstructures [24,41]. GOS maps for the studied conditions are displayed in Fig. 9a-c, while Fig. 9d shows
321 the GOS value for the subpopulations of small and large grains, averaged over each map. It can be noticed that
322 most of the small grains exhibit low GOS values and these keep almost constant during deformation. On the
323 other hand, the large grains present a wider spread of GOS, although always higher in average than the small
324 grains. Moreover, the average GOS for the set of large grains increases with strain. These values are
325 compatible with the description of the large grains accommodating most of the strain by developing internal
326 lattice distortions and subgrains, eventually leading to new, small recrystallized grains. Nevertheless, this
327 analysis should not be taken as a definitive proof of CDRX due to the fact that the GOS parameter is sensitive
328 to the grain size and shape [41,42]; this is the reason why the former subgrain analysis is preferred.

329



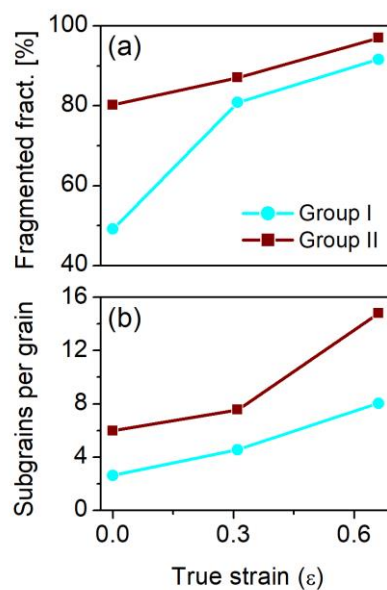
330
 331 **Figure 9.** Grain orientation spread (GOS) maps for the Zn-Cu-Ti sheet under study. (a) As-received; (b) UAT
 332 prestrain ($\epsilon \sim 0.31$); (c) second UAT loading ($\epsilon \sim 0.66$). HABs ($\theta > 12^\circ$) are displayed in black lines. (d) Average
 333 GOS values for the subpopulations of grains in each map (error bars correspond to the standard deviation).
 334

335 It is interesting to note that the fragmentation process described so far does not appear to start with
 336 the tensile straining, as the initial microstructure already reveals an important amount of subgrains and
 337 internal LABs within the large grains. Evidently, the previous warm-temperature rolling process of the sheet
 338 produces a partial fragmentation of the cast and hot-rolled structure in a way that can be inferred from the
 339 EBSD data of the as-received material. Additionally, although the present study is not intended to investigate
 340 the manufacturing process of the commercial sheet, the presence of the TiZn_{16} phase particles along some of
 341 the grain boundaries (Fig. 1c) may contribute to the fragmentation mechanism described. Early studies on
 342 low-alloyed zinc attributed the grain-refining ability of the Ti-bearing alloys to the intermetallic particles
 343 [25]; and it was also reported that impurity particles along grain boundaries in rolled CP-Zn stabilize its fine-
 344 grained structure [36].

3.4. Influence of orientation on the fragmentation behavior

Concerning the orientations of the large grains and their relationship with the main texture component, the fragmentation behavior appears to be different between the grains which have their c -axis pointing away from the ND and those whose basal planes are more closely parallel to the plane of the sheet. In this regard, we split the population of large grains into two groups according to their average orientation: (I) the large grains with their c -axis having an angular deviation to the ND greater than 45° , on one hand; and (II) those large grains whose basal planes are inclined to the plane of the sheet closely than 45° , on the other. The grains which fulfill the first condition for the three states considered represent a fraction of only 20%, 18% and 11%, respectively, of the total population of large grains. Meanwhile, group II involves the large grains that are part of the main texture component. Following Fig. 10a, the fragmented fraction of the large grains belonging to group I is smaller than the one of group II, particularly in the initial state. After the first tensile test there is a marked increase in the fragmented portion, but this is still smaller than the fraction of grains which develop subgrains near the main texture component (group II). Moreover, the average number of subgrains contained in the fragmented grains (Fig. 10b) is higher for group II; and this value increases with strain more rapidly than the amount of subgrains within group I. Average grain sizes of both groups are almost the same. In summary, the grains oriented near the main texture component are more easily fragmented than the ones which have their c -axis pointing away from the plane of the sheet, accounting for $\sim 90\%$ of the subgrains developed along the whole deformation process.

364



365

366 **Figure 10.** Differences in the fragmentation process of the Zn-Cu-Ti sheet according to the mean orientation
367 of the larger grains. Group I: large grains whose *c*-axis have an angular deviation to the ND greater than 45°;
368 group II: large grains with *c*-axis inclined to the ND closely than 45°. The fragmented fraction (a) and the
369 average amount of subgrains per grain (b) are higher for grains of group II.
370

371 Differences in grain fragmentation behavior according to orientation are common in other materials,
372 such as in FCC metals having grains with the Goss orientation [15]. In the case of hexagonal alloys, due to
373 their anisotropy it is expected that preferred orientation plays even a major role in the DRX mechanisms, and
374 that is the case in fact for Mg alloys, as reviewed by Kaibyshev [18]. Del Valle and Ruano [43] found that DRX
375 is enhanced in an AZ31 alloy subjected to tensile straining at $\sim 0.5\text{--}0.6T_m$ when the stress axis is normal to the
376 *c*-axis in samples with a strong basal fiber. They attributed this behavior to the enhancement of multiple-slip
377 instead of basal single-slip. Likewise, Kaibyshev [18] states that, since non-basal glide is mandatory for the
378 development of CDRX in Mg alloys, textures which hinder basal activity are optimal for grain-refining by
379 means of CDRX. Thus, these observations are similar to the current Zn-Cu-Ti sheet in the sense that the CDRX
380 process is enhanced for the grains having their *c*-axis more close to the ND, in which basal slip is
381 geometrically less favored than $\langle c+a \rangle$ 2nd order pyramidal glide. Furthermore, simulations by means of the
382 affine-VPSC model (assuming relative critical shear stresses of 15 and 10 for the prismatic and the 2nd order
383 pyramidal systems, respectively, with respect to the basal system; and an aspect ratio of the ellipsoidal
384 inclusion of 2.4:1.4:1, which corresponds to the average aspect ratio of the as received microstructure) show
385 that grains belonging to group II present $\sim 60\%$ more pyramidal $\langle c+a \rangle$ activity than grains in group I.
386

387 3.5. Boundary analysis

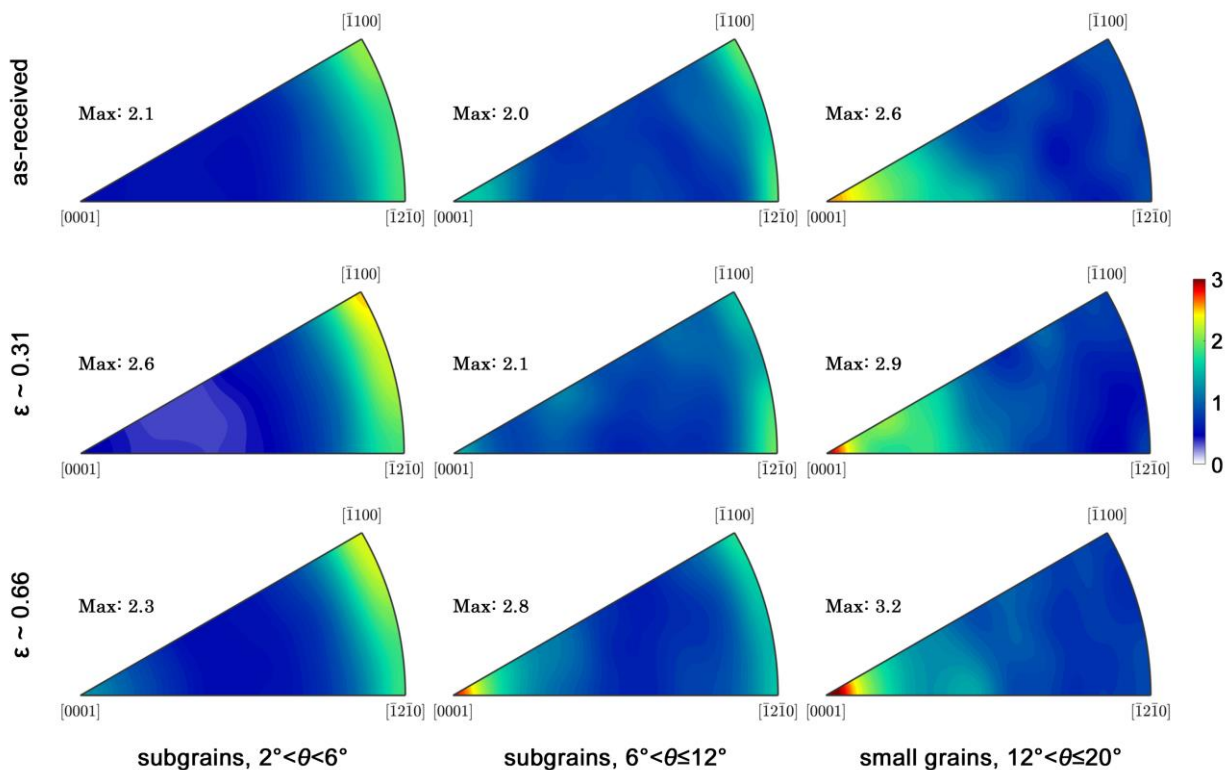
388 Regarding the fragmentation process of the grains in group II, the evolution of their subgrains'
389 boundaries can be further analyzed in terms of their angles and axes of misorientations. By counting the
390 length of the LABs with a misorientation angle range of $2^\circ < \theta \leq 6^\circ$ and $6^\circ < \theta \leq 12^\circ$, we can estimate the fraction
391 of these kind of boundaries for the three levels of strain. This counting includes the inner LABs which do not
392 form a closed subgrain, in order to capture the gradual process of formation of a fragment. It is important to
393 note that these inner boundaries account for almost half of the fraction of LABs in the initial state. Although
394 the fraction of $2^\circ < \theta \leq 6^\circ$ approximately doubles the $6^\circ < \theta \leq 12^\circ$ portion for the three strain levels, the former
395 fraction decreases with strain (from $\sim 19\%$ to $\sim 16\%$) and the latter increases from $\sim 8\%$ to $\sim 11\%$. Also, the

396 inner LABs which still do not form a subgrain belong mostly to the $2^\circ < \theta \leq 6^\circ$ range and their total length
 397 reduces notably between the as-received and the final condition.

398 Figure 11 shows the distribution of misorientation axes of the LABs, separated into the two angular
 399 ranges previously mentioned; and also of the boundaries between small grains having a misorientation angle
 400 $12^\circ < \theta \leq 20^\circ$ -i.e., those subgrains which we assume that have evolved from LABs to HABs. The axes
 401 distributions are shown in Fig. 11 as probability density plots in the unit triangle of the stereographic
 402 projection, reflecting the relative importance of certain axes of misorientation across boundaries. The
 403 $2^\circ < \theta \leq 6^\circ$ fraction of LABs presents axes distributions predominantly based on rotations around directions
 404 lying on the basal planes ($\langle hki0 \rangle$ directions), particularly for the lower levels of deformation. On the other
 405 hand, for the LABs in the range of $6^\circ < \theta \leq 12^\circ$, the axes distributions show an increasing tendency of rotations
 406 around the basal plane's normal. Moreover, it is interesting to note that the boundaries between the newly
 407 formed small grains have also an axis distribution concentrated in the basal pole, even more pronounced than
 408 the former LABs.

409 Applying the same analysis to the LABs developed within the large grains oriented away from the
 410 plane of the sheet (group I), it was found that the axis distribution corresponds mostly to rotations around
 411 the $\langle hki0 \rangle$ directions. This tendency is the same for the whole range of LABs.

412



413

414 **Figure 11.** Distribution of misorientation axes across boundaries (within the group of grains with *c*-axis
415 deviation from the ND <45°) for the conditions indicated. Left and middle columns: LABs within large grains
416 in the angular ranges specified. Right column: HABs of small grains with misorientation angles ranging from
417 12° to 20°.
418

419 From a crystal plasticity point of view, differences of orientation between (sub)grains by rotations
420 around an axis parallel to the basal plane would be explained by slip in the basal <a> or pyramidal <c+a>
421 systems. Rotations about the *c*-axis, on the other hand, could be due to prismatic <a> slip by a Taylor-axis
422 analysis [24] or by multiple glide of the pyramidal <c+a> systems. Since prismatic <a> activity is rarely found
423 in Zn at room temperature [44,45], the rotations around the *c*-axis observed for the intermediate
424 misorientation ranges (i.e., between 6° and 20°) could be attributed to multiple slip of 2nd order pyramidal
425 system. This is consistent with the above discussion about the role of <c+a> activity in CDRX for the grains in
426 group II.

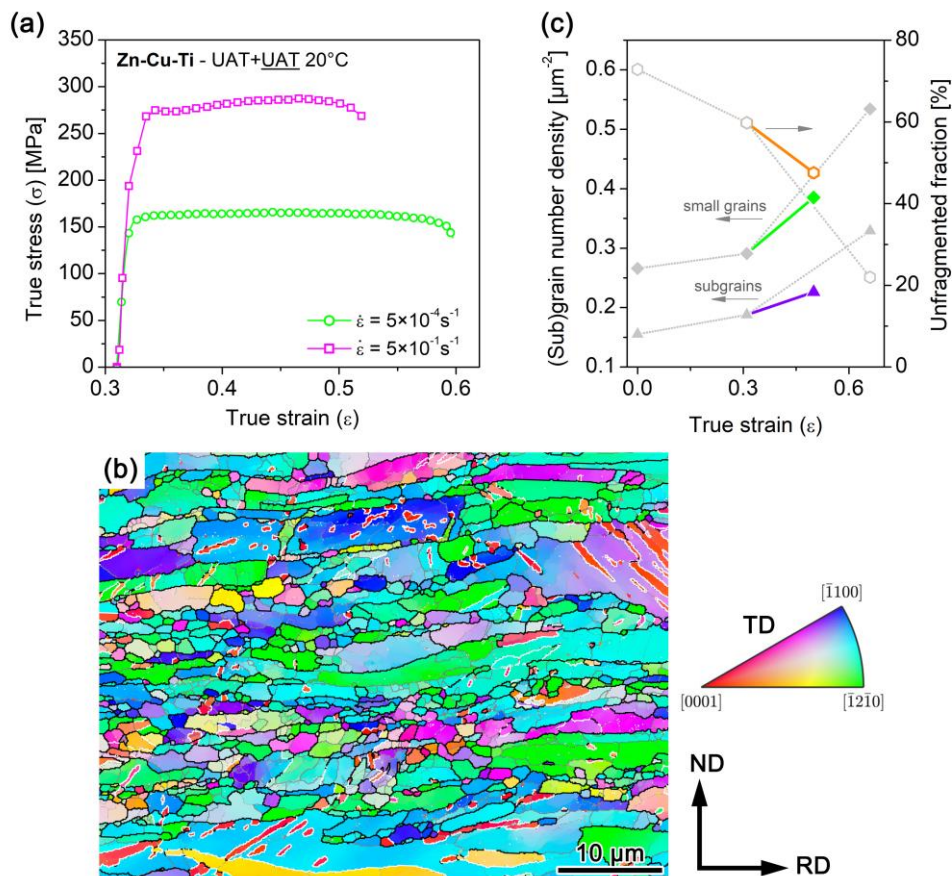
427 Even though this analysis would require a more thorough examination of the boundary character and
428 the local geometrically-necessary dislocations, the mechanisms involved can explain the stability and the
429 moderate weakening of the crystallographic texture observed during straining. Indeed, the grains originally
430 inclined away from the sheet plane tend to rotate –mainly by basal slip– toward the ND, stabilizing the
431 principal texture component. Meanwhile, the large grains oriented near the main texture component
432 accommodate deformation mostly by developing subgrains which have similar orientations to that of their
433 parent grains, though with moderate rotations around both the <*hki*0> directions and the *c*-axis due to a
434 combined activity of basal <a> and pyramidal <c+a> dislocations. These rotations make up the slight
435 dispersion of the texture around its main component. Similar findings regarding texture weakening by
436 subgrain rotation during CDRX have been reported in the literature on Mg alloys [39,40].
437

438 **3.6. Effect of a higher strain rate**

439 Since CDRX is a strain and thermally activated phenomenon [15,21], and considering that for Zn room
440 temperature corresponds to $0.42T_m$, it would be expected that a change in strain rate may influence the
441 deformation mechanisms in the material under study. In this regard, a tensile test of a second-step sample
442 was carried out at a strain rate of $5 \times 10^{-1} \text{ s}^{-1}$ (i.e., three orders of magnitude faster than the previous tests) and
443 the obtained microstructure was analyzed. The corresponding flow curve and EBSD map are included in Fig.
444 12.

445 Tensile response at this high strain rate (Fig. 12a) exhibits ~73% higher maximum true stress and
 446 ~33% lower total engineering elongation, reflecting the relatively high strain-rate sensitivity of this alloy.
 447 Also, the shape of the flow curve shows that at higher strain rate the work-hardening effect is increased. The
 448 differences briefly outlined may reflect the changes noticeable in the microstructure. In fact, the EBSD map
 449 (Fig. 12b) reveals more twins and larger grains for this microstructure when compared to the corresponding
 450 low strain-rate map (Fig. 4c). Twin boundary density is $0.209 \mu\text{m}/\mu\text{m}^2$, approximately 4.2 times higher than
 451 the value for the slow strain rate test. Applying the same fragmentation criterion described above, as shown
 452 in Fig. 12c, after the present test there is a 48% of remaining parent large grains, implying an important
 453 difference with respect to the 23% corresponding to the slow tensile test. Also, the amount of subgrains and
 454 small grains is considerably smaller, so it is evident that the fragmentation mechanism is not as easily
 455 operative at this higher strain rate.

456



457

458 **Figure 12.** Effect of a higher strain rate ($5 \times 10^{-1} \text{ s}^{-1}$) on the second UAT test of the Zn-Cu-Ti sheet. (a) Flow curve (the second low-strain rate test is repeated for reference). (b) EBSD map for the high strain-rate test (total strain level ~ 0.50): grain boundaries plotted over IPF-TD map; black lines: HABs ($\theta > 12^\circ$), gray lines:
 459
 460
 461 LABs ($2^\circ < \theta \leq 12^\circ$), white lines: $\{10\bar{1}2\}\langle 10\bar{1}\bar{1}\rangle$ twin boundaries. (c) Comparison of (sub)grain numbers and

462 fragmented fractions of large grains between low (gray dotted lines) and high strain rates (colored solid
463 lines).
464

465 It seems clear now that, for the lower strain rate, CDRX is a non-negligible deformation mechanism
466 which allows the microstructure to accommodate macroscopic strain by formation and rotation of subgrains
467 and of the newly developed small grains. However, since CDRX is a thermally activated process, if the strain
468 rate is increased at the amount applied, there is not enough time for the microstructure to develop profuse
469 CDRX and slip is no longer sufficient for accommodating the imposed deformation. Twinning is evidently
470 activated at these conditions, but it is not enough for strain accommodation so the macroscopic failure is
471 achieved at a lower strain level.

472 Additionally, from the preceding observations, it can be inferred that the extended ductility found by
473 Schlosser et al. [14] during bilinear strain paths of the current Zn-Cu-Ti alloy at the same conditions (i.e.,
474 room temperature and strain rate of $5 \times 10^{-4} \text{ s}^{-1}$) may be explained by the CDRX process, since a similar grain
475 fragmentation behavior was observed in the corresponding microstructures, as it was preliminarily shown in
476 that work.

477

478 **4. Conclusions**

479 The analyzed Zn-Cu-Ti sheet subjected to tensile straining at room temperature and slow strain rate
480 shows a grain fragmentation process which can be categorized as continuous dynamic recrystallization. This
481 process exhibits a growing development of subgrains of low misorientation (LABs) within initially large
482 grains, which gradually increase their number and become new grains by the continuous increase of their
483 misorientation.

484 The texture evolution reflects the combination of plastic slip with CDRX, since the *c*-axes tend to keep
485 around the main texture component but with increasing deviation towards the ND and slight rotations of the
486 prismatic poles due to the subgrain development. Furthermore, the grain fragmentation process due to CDRX
487 is orientation dependent, since the large grains which belong to the main texture component exhibit more
488 abundant subgrain formation.

489 The misorientation analysis applied to the LABs and to the low-angle fraction of small grains suggests
490 that CDRX may be associated with enhanced activity of 2nd order pyramidal $\langle c+a \rangle$ slip. Increasing the strain

491 rate results in a noticeable change of the deformation mechanisms, as the microstructure develops twinning
492 in detriment of CDRX. This in turn is related to an important raise of the flow stress and a loss of ductility.

493

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499

500 **Data availability**

501 The raw/processed data required to reproduce these findings cannot be shared at this time as the data
502 also forms part of an ongoing study.

503

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622 **Figure 1.** Microstructure and crystallographic texture of the Zn-Cu-Ti sheet in the as-received condition. (a)
623 and (b): band contrast of the EBSD maps taken from the RD-ND and RD-TD sections, respectively. Some large,
624 elongated grains are indicated by arrows and clusters of small, equiaxed grains are bounded in dashed boxes.
625 (c) Secondary electron micrograph revealing the TiZn_{16} phase. (d) Basal and prismatic pole figures
626 recalculated from the EBSD data (X and Y labels correspond to TD and RD, respectively).

627 **Figure 2.** Bilinear tensile test (UAT+UAT) of the Zn-Cu-Ti sheet. (a) Sample pre-strained along RD with the
628 major true strain field measured by digital image correlation. Geometry of the second UAT samples are
629 shown in dashed lines. (b) Corresponding flow curves. The prestrain value of the second-step samples is
630 ~ 0.31 .

631 **Figure 3.** Band contrast maps and pole figures of the deformed Zn-Cu-Ti samples. (a) and (c): UAT prestrain
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633 **Figure 4.** Grain boundaries plotted over orientation maps corresponding to the IPF with respect to the TD.
634 (a) As-received; (b) UAT prestrain ($\epsilon \sim 0.31$); (c) second UAT loading ($\epsilon \sim 0.66$). Black lines: HABs ($\theta > 12^\circ$); gray
635 lines: LABs ($2^\circ < \theta \leq 12^\circ$). (d) Detail showing $\{10\bar{1}2\}\{10\bar{1}\bar{1}\}$ twin boundaries in white.

636 **Figure 5.** Grain statistics of the Zn-Cu-Ti sheet microstructure for each strain level, subdividing the
637 populations by a size criterion: (a) number density of reconstructed grains; (b) average grain sizes
638 (calculated as equivalent circle diameter); (c) mean aspect ratios. Error bars correspond to the standard
639 deviation of each population.

640 **Figure 6.** Map detail illustrating fragmentation of the larger grains ($\epsilon \sim 0.31$). Gray lines: HABs ($\theta > 12^\circ$) of
641 small grains ($\leq 3 \mu\text{m}$); black lines: HABs of larger grains ($> 3 \mu\text{m}$); thick red lines: LABs ($2^\circ < \theta \leq 12^\circ$) of subgrains
642 within large grains; thin red lines: inner LABs. Blue: subgrains within large grains; light blue: large grains
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644 unfragmented parent grains).

645 **Figure 7.** Evolution of the populations of grains and subgrains, and of the fraction of large grains which are
646 not fragmented (number density of small grains is repeated from Fig. 5 for reference).

647 **Figure 8.** Misorientation gradients across fragmented grains: (a) and (b) EBSD map details of the first
648 ($\epsilon \sim 0.31$) and the second ($\epsilon \sim 0.66$) UAT tests, respectively. (c-e) Changes of misorientation angle over the
649 lines indicated in the maps (solid lines: cumulative or point-to-origin misorientation; dotted lines: point-to-
650 point misorientation).

651 **Figure 9.** Grain orientation spread (GOS) maps for the Zn-Cu-Ti sheet under study. (a) As-received; (b) UAT
652 prestrain ($\epsilon \sim 0.31$); (c) second UAT loading ($\epsilon \sim 0.66$). HABs ($\theta > 12^\circ$) are displayed in black lines. (d) Average
653 GOS values for the subpopulations of grains in each map (error bars correspond to the standard deviation).

654 **Figure 10.** Differences in the fragmentation process of the Zn-Cu-Ti sheet according to the mean orientation
655 of the larger grains. Group I: large grains whose *c*-axis have an angular deviation to the ND greater than 45° ;
656 group II: large grains with *c*-axis inclined to the ND closely than 45° . The fragmented fraction (a) and the
657 average amount of subgrains per grain (b) are higher for grains of group II.

658 **Figure 11.** Distribution of misorientation axes across boundaries (within the group of grains with *c*-axis
659 deviation from the ND $< 45^\circ$) for the conditions indicated. Left and middle columns: LABs within large grains
660 in the angular ranges specified. Right column: HABs of small grains with misorientation angles ranging from
661 12° to 20° .

662 **Figure 12.** Effect of a higher strain rate ($5 \times 10^{-1} \text{ s}^{-1}$) on the second UAT test of the Zn-Cu-Ti sheet. (a) Flow
663 curve (the second low-strain rate test is repeated for reference). (b) EBSD map for the high strain-rate test
664 (total strain level ~ 0.50): grain boundaries plotted over IPF-TD map; black lines: HABs ($\theta > 12^\circ$), gray lines:
665 LABs ($2^\circ < \theta \leq 12^\circ$), white lines: $\{10\bar{1}2\}\{10\bar{1}\bar{1}\}$ twin boundaries. (c) Comparison of (sub)grain numbers and
666 fragmented fractions of large grains between low (gray dotted lines) and high strain rates (colored solid
667 lines).