



# Low-metallicity stellar halo populations as tracers of dark matter haloes

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## ABSTRACT

We analyse the density profiles of the stellar halo populations in eight Milky Way mass galaxies, simulated within the  $\Lambda$  cold dark matter scenario. We find that accreted stars can be well fitted by an Einasto profile, as well as any subsample defined according to metallicity. We detect a clear correlation between the Einasto fitting parameters of the low-metallicity stellar populations and those of the dark matter (DM) haloes. The correlations for stars with  $[\text{Fe}/\text{H}] < -3$  allow us to predict the shape of the dark matter profiles within residuals of  $\sim 10$  per cent, in case the contribution from *in situ* stars remains small. Using Einasto parameters estimated for the stellar halo of the Milky Way and assuming the later formed with significant contributions from accreted low-mass satellite, our simulations predict  $\alpha \sim 0.15$  and  $r_2 \sim 15$  kpc for its dark matter profile. These values, combined with observed estimations of the local DM density, yield an enclosed DM mass at  $\sim 8$  kpc in the range  $3.9\text{--}6.7 \times 10^{10} M_\odot$ , in agreement with recent observational results. These findings suggest that low-metallicity stellar haloes could store relevant information on the DM haloes. Forthcoming observations would help us to further constrain our models and predictions.

**Key words:** Galaxy structure – galaxies evolution – galaxies formation – galaxies haloes.

## 1 INTRODUCTION

The stellar halo populations are important components of galaxies which extend around them out to few hundreds of kiloparsecs. Observations show that they are complex structures where debris from stripped satellites, stars and surviving satellites can be identified within diffuse stellar envelopes. Stellar halo populations have been detected out to several kiloparsecs from the galaxy centre in nearby galaxies, regardless of their morphology (e.g. Rejkuba et al. 2011; Sarajedini et al. 2012; Crnojević et al. 2013; Ibata et al. 2014). There are also pieces of evidence of stellar haloes in galaxies at  $z \sim 1$  (Trujillo & Bakos 2013). The Milky Way (MW) is the best observed stellar halo where detailed kinematics and chemical abundances for individual stars are being obtained from different surveys, allowing the identification of chemical–dynamical patterns (e.g. Carollo et al. 2007; Nissen & Schuster 2010; Deason, Belokurov & Evans 2011a; Sheffield et al. 2012; Beers et al. 2012; Schönrich, Asplund & Casagrande 2014). Such a detailed analysis is now being extended to nearby galaxies as in the case of M31 and M32 (e.g. Sarajedini et al. 2012; Ibata et al. 2014). These observational results provide invaluable information to constrain cosmological models.

Simulations within a hierarchical clustering scenario show that most of the stars in the stellar haloes come from accreted satellites in wet or dry mergers (e.g. Bullock & Johnston 2005; Purcell, Kazantzidis & Bullock 2009; Zolotov et al. 2009; Cooper et al. 2010; Font et al. 2011; Tissera, White & Scannapieco 2012; Tissera et al. 2013; Pillepich et al. 2014). Hydrodynamical simulations, including chemical evolution, produce stellar halo populations that are old and highly  $\alpha$ -enriched in agreement with observations. Tissera et al. (2014) show how the metallicity distribution of stellar haloes varies with the radius depending on the accretion histories. Simulations also detect the existence of a fraction of stars formed *in situ*, which could explain the presence of stars with lower  $\alpha$ -enhancement and significant rotation velocities reported in observations of the MW (e.g. Sheffield et al. 2012).

Stellar haloes contain an important fraction of old and metal-poor stars, which are a window towards the first epochs of the formation of galaxies (e.g. Placco et al. 2014, and references therein) and hence might store information relevant to galaxy formation. Using cosmological simulations of MW mass galaxies, Tissera et al. (2014) reported the fraction of low-metallicity stars to increase with radius. They also found about  $\sim 60\text{--}90$  per cent of these stars formed in accreted low-mass satellites ( $M < 10^9 M_\odot$ ). Low-metallicity stars populate the whole stellar haloes, mapping the potential well more uniformly out to the virial radius.

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The stellar halo populations inhabit the potential well determined by the dark matter (DM) mass. As satellites fall in, they are disrupted, contributing with stars and gas to the haloes. Unlike the gas that experiences significant dissipation during infall, the stars are expected to store information for a longer time (Eggen, Lynden-Bell & Sandage 1962). In particular, the low-metallicity stars might be the best tracers of the potential well, as they come in small-mass satellites which are easily disrupted and are mainly subject to the gravitational potential well dominated by the DM.

The aim of this Letter is to explore if the stellar halo populations, and particularly the low-metallicity stars, can trace the DM distributions. For this purpose, we use a suite of MW mass haloes from the Aquarius project which were modified to include gas dynamics. The chemodynamical properties and history of assembly of the stellar halo populations of these haloes have been extensively studied by Tissera et al. (2012, 2013, 2014). The DM profiles of six of them were analysed by Tissera et al. (2010).

This Letter is organized as follow. Section 2 describes the numerical simulations and models. In Section 3, we analyse and discuss the density profiles of stellar halo populations and correlate them with those of their DM haloes. In Section 4, we summarize our main results.

## 2 THE SIMULATED HALOES

We analysed the suite of eight MW mass galaxy systems of the Aquarius project which were modified to include baryons by Scannapieco et al. (2009). The systems were selected from a cosmological volume of  $100 \text{ Mpc } h^{-1}$  box, with only a mild isolation criterion imposed at  $z = 0$  and the requirement to have similar total mass to the MW [see Springel et al. (2008) for further details on the generation of the initial conditions]. The simulations are consistent with a  $\Lambda$  cold dark matter cosmogony with  $\Omega_m = 0.25$ ,  $\Omega_\Lambda = 0.75$ ,  $\sigma_8 = 0.9$ ,  $n_s = 1$  and  $H_0 = 100 h \text{ km s}^{-1} \text{ Mpc}^{-1}$  with  $h = 0.73$ . The simulated haloes have  $\approx 1$  million total particles within the virial radius and virial mass in the range  $\approx 5\text{--}11 \times 10^{11} M_\odot h^{-1}$ . DM particles have masses of the order of  $\approx 10^6 M_\odot h^{-1}$  while initially gas particles have  $\approx 2 \times 10^5 M_\odot h^{-1}$ . More details on the simulated galaxies can be found in Scannapieco et al. (2009) and Tissera et al. (2012). We used a version of GADGET-3 which includes stochastic star formation, a physically motivated supernova (SN) feedback model and chemical enrichment by Type II and Ia SNe. Details on the SN model can be found in Scannapieco et al. (2005, 2006). Scannapieco et al. (2009, 2010, 2011) studied the formation of the disc and the spheroidal components of these systems and their dynamical and photometric properties. The main galaxies have a variety of morphologies with different bulge-to-disc ratios.

We follow Tissera et al. (2012, 2013, 2014), and apply a criterion based on angular momentum and binding energy to separate the different dynamical components. Stellar halo stars are a subset of the spheroid stars, those for which the binding energy is lower than the minimum one at the half-mass radii of the galaxies. Note that we do not consider in the stellar haloes particles belonging to friends-of-friends groups of more than 20 particles. Hence, our stellar haloes do not contain satellites, but they do include streams and small substructures from disrupted satellites.

Tissera et al. (2013, 2014) analysed the formation of the stellar haloes and in particular how to use the chemical distributions to set constraints on the history of mass assembly. In our simulations, the stellar diffuse haloes are mainly made by debris stars which formed in accreted satellites (outside or inside the virial radius) and disc-heated stars (so-called *in situ*). The latter are found to be

more centrally concentrated. Tissera et al. (2014) showed that stellar haloes formed by the accretion of massive subgalactic systems have steeper metallicity gradients. They also reported that low-metallicity stars are mainly contributed by low-mass satellites and are more frequent in the outskirts of haloes.

## 3 THE RELATION BETWEEN THE STELLAR AND DM HALOES

We construct the density profiles of the stellar halo populations by using spherical shells from  $\sim 2 \text{ kpc } h^{-1}$  to the virial radius, centred at the galaxy centre. The density profiles are constructed for the *in situ* and accreted stellar components, separately. Within the accreted stars, we include debris (i.e. stars formed in satellites before coming into the virial radius of the main galaxy) and endo-debris (i.e. those stars formed in gas-rich satellites within the virial radius of the main galaxy). We define two population subsamples according to their metallicity: very metal poor (VMP) and extremely metal poor (EMP) by requiring them to satisfy  $[\text{Fe}/\text{H}] < -2$  and  $[\text{Fe}/\text{H}] < -3$ , respectively.

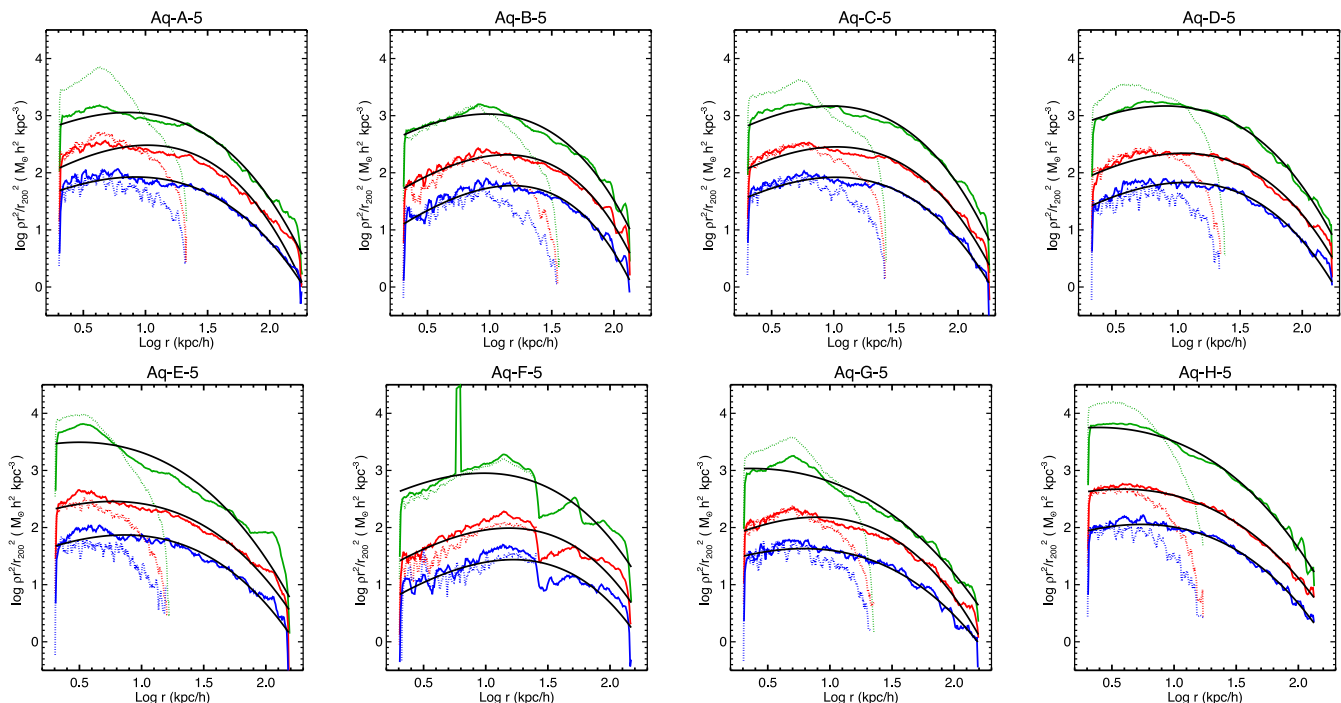
We estimate the density profiles for each metallicity subsample and for both types of stellar populations. The profiles for the *in situ* (green dotted lines) and accreted (solid green lines) stellar components, without metallicity segregation, are shown in Fig. 1. These distributions show how *in situ* stars are more concentrated in the central regions while accreted ones extend up to the virial radius. There are two haloes with larger relative *in situ* contributions: Aq-A-5 and Aq-H-5. In the case of Aq-B-5 and Aq-F-5, the contributions of *in situ* are very small.

The lower subsequent set of profiles included in Fig. 1 are those obtained for the VMP (red lines) and EMP (violet lines) subsamples. The contributions from *in situ* stars diminish systematically as the metallicity threshold gets more stringent, since these stars have a large probability to have been formed from recycled material, polluted by previous stellar generations. Disc-heated stars are younger than accreted stars and have larger  $\alpha$ -enhancements, on average (Tissera et al. 2012, 2013). Hence, as we move to lower metallicity subsamples, the density profiles are mainly determined by accreted stars. As discussed by Tissera et al. (2014), accreted stars are mainly contributed by small galaxy satellites (few times  $10^9 M_\odot$ ) and have a frequency which increases outwards.

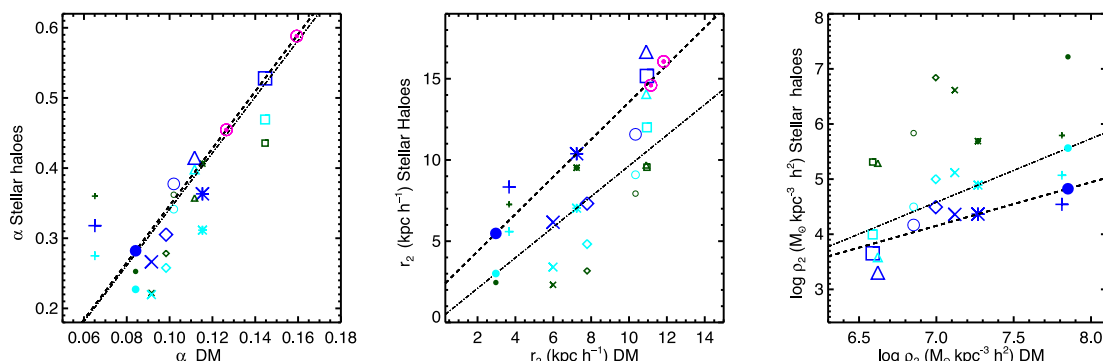
We fit Einasto profiles to the density distributions from  $\sim 2 \text{ kpc } h^{-1}$  to the virial radius. The free parameters are  $n = 1/\alpha$ ,  $r_2$  and  $\rho_2$  which indicate the sharpness of the profiles, and the radius and density where their logarithmic slope takes the isothermal value, respectively. The accreted stellar haloes are resolved with more than 100 000 star particles while these numbers decrease as more stringent metallicity thresholds are imposed. The accreted EMPs are resolved by  $\sim 5000\text{--}10\,000$  star particles.

In these simulated haloes, the density profiles determined by the accreted stars are well fitted by the Einasto profile, regardless of the metallicity threshold adopted (Fig. 1; solid black lines). This is not the case for the whole stellar haloes because of the contribution of *in situ* stars which are more concentrated than accreted ones. However, the contribution of *in situ* stars to the EMPs is between 10 and 20 per cent of the total EMP mass. And in fact, the Einasto profile also yields a good fit for the whole EMPs.

The DM profiles of these haloes were studied by Tissera et al. (2010), who found that they are well fitted by Einasto profiles (in the same radial range). The final properties of the DM profiles are the outcome of the joint evolution with their main galaxy and, consequently, they have been already modified by its formation.



**Figure 1.** Density profiles of the *in situ* (dashed line) and accreted (solid line) stellar halo components (green lines) and the corresponding density profiles obtained by applying metallicity thresholds:  $[\text{Fe}/\text{H}] < -2$  (VMP; red) and  $[\text{Fe}/\text{H}] < -3$  (EMP; blue). The black solid lines depict the Einasto profile fit for each subsample.

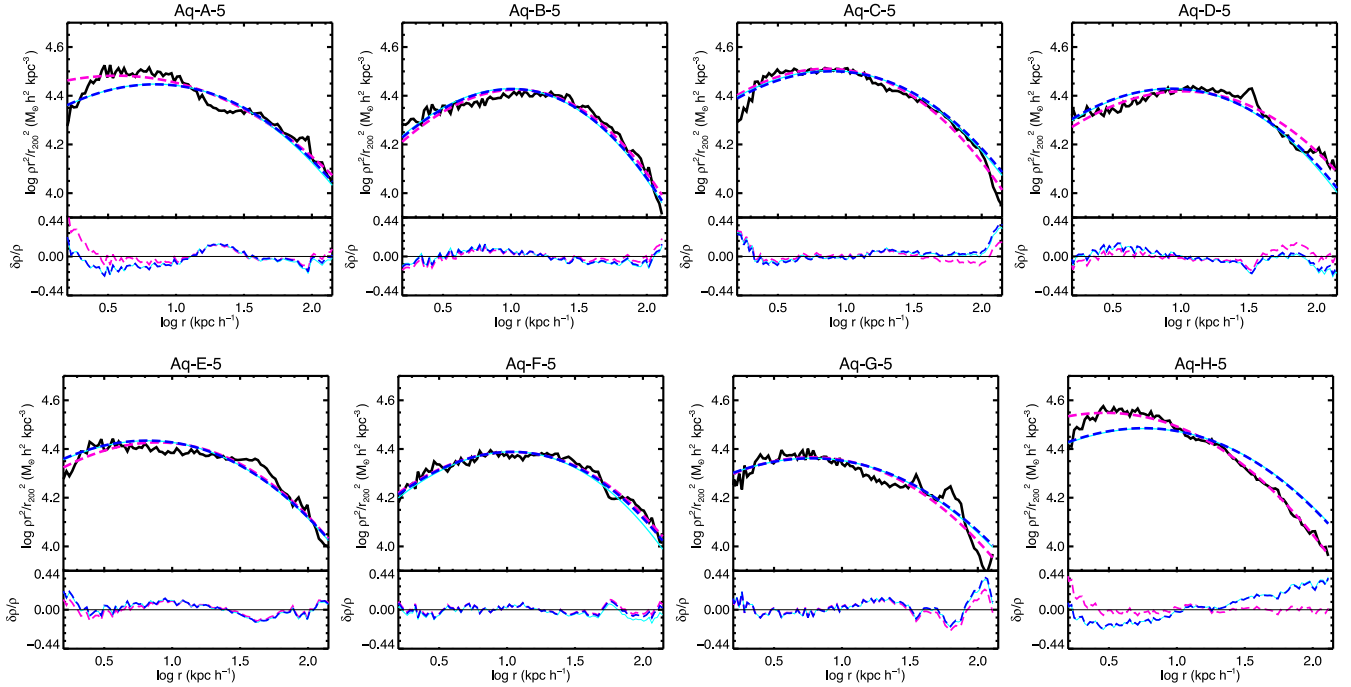


**Figure 2.** Einasto fitting parameters of the density profiles estimated for the total accreted stars in the haloes (green symbols), for accreted EMPs (blue symbols) and for the whole EMPs (cyan symbols) as a function the corresponding best-fitting parameters obtained for the DM profiles. The fitted linear regressions to the whole and accreted EMP parameters are also shown (thin and thick dashed lines). We include predicted  $\alpha$  and  $r_2$  predicted for the MW (magenta symbols), using available observed estimations under the hypothesis that MW halo formed mainly by accreted stars from low-mass satellites.

In order to explore at what extent the accreted stellar haloes can trace the DM ones, we correlate the Einasto fitting parameters for our metallicity-selected subsamples of accreted stars with those obtained for the corresponding DM profiles by Tissera et al. (2010). Statistically, significant correlations for the  $\alpha$ ,  $r_2$  and  $\rho_2$  parameters are measured, principally for the accreted and whole EMPs, which yield correlation factors within the range  $\sim 0.80$ – $0.85$ . These factors go down to  $\sim 0.50$ – $0.65$  for all accreted stars and VMPs. Hence, hereafter, we only analyse the EMPs. We fit linear regressions to the correlations defined by both EMP subsamples, by applying a Levenberg–Marquardt least-squares algorithm. In Fig. 2, we show the linear regressions determined by the accreted and whole EMPs. The rms values found for the  $\rho_2$  and  $n$  linear fittings are smaller

than  $\sim 0.02$ , but for the  $r_2$  characteristic scales, it is  $\sim 0.2$ . For the whole EMP subsample, we also find an increase of a factor of 4 in the rms of  $\rho_2$ .

We use these correlations to estimate the DM density profiles by using the fitting parameters of the EMPs and the linear relations shown in Fig. 2. Fig. 3 shows the predicted and simulated DM profiles, as well as the best-fitting Einasto DM profiles obtained directly from the simulated data. Although the shape of the DM profiles can be well estimated from those of the EMPs, the normalization is more difficult to predict without an independent estimation of, e.g. the total DM mass. For this reason, we renormalized the profiles so that the predicted and real DM profiles agree at  $r_2$ . We could instead use our predicted DM profiles to infer the DM mass of the systems;



**Figure 3.** Predicted DM profiles from the fitting relations shown in Fig. 2 for the accreted EMPs (blue dashed lines), the whole EMPs (cyan solid lines), the direct Einasto fits to the DM profiles (magenta dashed lines) and the real DM profiles from the simulations (black lines). In the lower panels, we show the residuals of the direct Einasto fit to the DM profiles and of those predicted by the accreted EMP and whole EMP profiles to the simulated DM profiles. Note that the predicted profiles have been renormalized to reproduce the DM profiles at  $\sim r_2$  (see the text for comments).

in this case, the DM mass is recovered with an  $\sim 0.7$  dex dispersion. This large dispersion reflects the problem with the normalization of the predicted profiles, indicating the need to have an independent measure of the DM mass. The small panels of Fig. 3 show the residuals of the DM profiles obtained by fitting the real data and those predicted by the EMPs to the simulated DM profiles. Except for Aq-A-5 and A-H-5, the galaxies with the largest *in situ* fractions, the residuals are always smaller than 10 per cent. These results show that the whole EMP subsample can also trace the shape of the DM profile. This is because it is dominated by accreted stars. We also tested the whole and accreted VMPs but in this case, there are larger dispersions and the predicted profiles are not always good.

The Einasto parameters obtained from the EMPs are in good agreement with those reported for the MW halo by Deason, Belokurov & Evans (2011b) and Sesar, Jurić & Ivezić (2011):  $\alpha \sim 0.60$ ,  $0.45$  and  $r_2 \sim 20$ ,  $22$  kpc, respectively. Instead, the total accreted stellar haloes yield profiles slightly more concentrated (i.e. smaller  $r_2$ ). This finding suggests that the MW halo might have been formed by a significant contribution of accreted stars from small satellites, with less than 20 per cent of *in situ* stellar mass. Under this assumption and using the reported observed values, our fitting relations predict the DM halo of the MW to have  $\alpha \sim 0.16$ – $0.13$  and  $r_2 \sim 13$ – $16$  kpc.

Even more, from recent measures of the DM density in the solar neighbourhood, we can normalize the simulated relation. The resulting MW predicted profile yields an enclosed DM mass at  $\sim 8$  kpc of  $\sim 3.9 \times 10^{10} M_\odot$  or  $\sim 6.7 \times 10^{10} M_\odot$ , if we assume  $\rho_{\text{dm}} = 0.0087 M_\odot \text{ pc}^{-3}$  (Read 2014) or  $\rho_{\text{dm}} = 0.015 M_\odot \text{ pc}^{-3}$  (Piffl et al. 2014), respectively. The agreement with the RAVE results by Piffl et al. (2014) is very encouraging. Detailed near-future observations of the MW stellar halo will help us to further constrain our models and predictions.

## 4 CONCLUSIONS

We analysed the density profiles of the stellar populations in eight MW simulated galaxies and compared them with those of their DM haloes. We focused on low-metallicity stars since they map the whole potential well with a frequency increasing outwards. These stars tend to come in low-mass satellites and hence did not experience dissipation within the host potential well.

The density profiles of accreted stars in the stellar haloes can be well fitted by an Einasto profile. Similarly, the density profiles of accreted stars with systematically lower  $[\text{Fe}/\text{H}]$  abundances can all be well fitted by Einasto profiles. This is not the case for the total stellar populations formed by *in situ* and accreted stars, since the former are much more concentrated, producing a total density profile which is best described by a broken power law.

The Einasto parameters obtained for the EMPs density profiles correlate well with the corresponding ones of the DM profiles (even if the *in situ* components are included, since they represent less than 20 per cent of the mass). The correlations get weaker as stars with higher metallicity are included. This can be understood since more metal-rich stars are contributed by more massive satellites and have more possibilities to have experienced dissipation and tend to be more centrally concentrated (Tissera et al. 2014). The obtained Einasto parameters for the EMPs are in good agreement with those reported for the MW stellar halo, suggesting that the latter might have been formed principally by the accretion of low-mass satellites. If we adopt the observed values under this hypothesis, the simulated correlations for the EMPs predict  $\alpha \sim 0.15$  and  $r_2 \sim 15$  kpc for the DM profile of the MW. Assuming observed measures of the local DM density to normalize predicted DM profile at  $\sim 8$  kpc yields an enclosed DM of  $\sim 3.9$ – $6.7 \times 10^{10} M_\odot$ , within this radius. These estimations are in agreement with results from the RAVE survey (Piffl et al. 2014), indicating the goodness of the predicted shape in

this region. Our predicted DM profile for the MW yields a total DM mass of  $\sim 7.3 \times 10^{11} h^{-1} M_{\odot}$ .

Our findings suggest that low-metallicity stellar haloes might store important information on the DM mass and the shape of DM profiles, helping to constrain the cosmological model and the details of galaxy formation models.

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