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Exchange of Microtubule Molecular Motors During Melanosome Transport in *Xenopus laevis*  
Melanophores is Triggered by Collisions with Intracellular Obstacles

**Authors:** Luciana Bruno · Mercedes Echarte · Valeria Levi

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Abstract The observation that several cargoes move bidirectionally along microtubules in vivo raised the question regarding how molecular motors with opposed polarity coordinate during transport. In this work, we analyzed the switch of microtubule motors during the transport of melanosomes in *Xenopus melanophores* by registering trajectories of these organelles moving along microtubules using a fast and precise tracking method. We analyzed in detail the intervals of trajectories showing reversions in the original direction of transport and processive motion in the opposite direction for at least 250 nm. In most of the cases, the speed of the melanosome before the reversion slowly decreases with time approaching zero then, the organelle returns over the same path moving initially at a very high speed and slowing down with time. These results

could be explained according to a model in which reversions are triggered by an elastic collision of the cargo with obstacles in the cytosol. This interaction generates a force opposed to the movement of the motor-driven organelle increasing the probability of detaching the active motors from the track. The model can explain reversions in melanosome trajectories as well as other characteristics of in vivo transport along microtubules observed by other authors. Our results suggest that the crowded cytoplasm plays a key role in regulating the coordination of microtubules-dependent motors.

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Keywords (separated by '-') Microtubule molecular motors - *Xenopus* melanophores - Coordination - Melanosome transport

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2 **Exchange of Microtubule Molecular Motors During Melanosome**  
3 **Transport in *Xenopus laevis* Melanophores is Triggered by**  
4 **Collisions with Intracellular Obstacles**

5 **Luciana Bruno · Mercedes Echarte ·**  
6 **Valeria Levi**

7  
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9 **Abstract** The observation that several cargoes move  
10 bidirectionally along microtubules in vivo raised the ques-  
11 tion regarding how molecular motors with opposed polarity  
12 coordinate during transport. In this work, we analyzed the  
13 switch of microtubule motors during the transport of mel-  
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18 original direction of transport and processive motion in the  
19 opposite direction for at least 250 nm. In most of the cases,  
20 the speed of the melanosome before the reversion slowly  
21 decreases with time approaching zero then, the organelle  
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23 speed and slowing down with time. These results could be  
24 explained according to a model in which reversions are  
25 triggered by an elastic collision of the cargo with obstacles in  
26 the cytosol. This interaction generates a force opposed to the  
27 movement of the motor-driven organelle increasing the  
28 probability of detaching the active motors from the track.  
29 The model can explain reversions in melanosome trajec-  
30 tories as well as other characteristics of in vivo transport along  
31 microtubules observed by other authors. Our results suggest

that the crowded cytoplasm plays a key role in regulating the  
coordination of microtubules-dependent motors.

**Keywords** Microtubule molecular motors ·  
*Xenopus* melanophores · Coordination ·  
Melanosome transport

**Introduction**

Molecular motors are responsible for the transport of a  
wide variety of components, which are positioned in the  
cytoplasm in a precise spatio-temporal manner. There are  
three families of motors involved in the transport of car-  
goes in the cell cytoplasm: myosin motors which move  
along actin filaments; as well as kinesin and dynein motors  
that move along microtubules toward their plus- and  
minus-end, respectively (reviewed in [1]).

The involvement of two motors of opposite polarity in  
organelle transport along microtubules raised the question  
about how transport is regulated along these cytoskeleton  
tracks. In several cellular systems it was observed that  
cargoes initially transported toward the microtubule minus-  
end frequently changed their direction and moved toward  
the plus-end and vice versa (see for example, [2–8]).

Three different models have been proposed to explain  
these experimental observations (reviewed in [1, 9, 10]).

According to the first model, opposed-polarity motors  
are involved in a tug-of-war with the stronger motor  
determining the direction of motion at any particular  
moment. This model predicts, for example, that the  
impairment of the plus-end directed transport should  
improve the instantaneous transport in the opposite direc-  
tion. However, it has been shown that mutations which  
damage the functionality of kinesin-II during pigment

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65 organelle transport in melanophores cells or of dynein or  
66 the dynactin complex in the case of lipid droplets transport  
67 in *Drosophila* embryos do not improve the transport in the  
68 opposite direction [10, 11].

69 The exclusionary presence model considers that while  
70 both types of motors can bind to a cargo, they cannot both  
71 do it at the same time. However, the total amount of either  
72 plus- or minus-end directed motors that is attached to a  
73 pigment organelle in frog melanophores does not change  
74 when stimulating the transport toward the plus- or the  
75 minus-end of the microtubules [11] in opposition to what it  
76 would be expected according to this model.

77 These and other experimental finding from several cel-  
78 lular systems support the coordination model to explain  
79 regulation of transport along microtubules (reviewed in  
80 [9]). According to this model, kinesin and dynein are  
81 simultaneously present on cellular cargo and their activities  
82 are coordinated so that when plus-end motors are active,  
83 minus-end motors are not, and vice versa.

84 The precise mechanism of coordination is still unknown;  
85 however, some protein complexes have been proposed to  
86 mediate the coordination between motors [12–14]. Dynactin  
87 is a potential candidate for coordination of microtubule-  
88 dependent motors in frog melanocytes [15]. Minus- and  
89 plus-end motors compete for binding to the same region of  
90 this protein complex and impairment of dynactin abolishes  
91 both plus- and minus-end motion of several bidirectional  
92 cargoes [8, 15]. Moreover, dynactin increases the proces-  
93 sivity of kinesin-2 [16] and of cytoplasmic dynein, which is  
94 not processive in the absence of this protein complex [17].

95 Regardless of the exact nature of how coordination is  
96 achieved, the process that triggers the switch of motors  
97 remains unexplained. In this article, we address this issue  
98 by studying reversions in trajectories of organelles trans-  
99 ported by microtubule-dependent motors in *Xenopus laevis*  
100 melanophores.

101 Melanophores cells are one of the cellular systems  
102 commonly used to study the function of molecular motors  
103 in vivo (reviewed in [18]). The major physiological task of  
104 these cells is to move pigment organelles called melano-  
105 somes in the cytoplasm, allowing animals to display color  
106 change. *Xenopus* melanophores have melanosomes filled  
107 with the black pigment melanin, and therefore these  
108 organelles can be easily imaged and discriminated from  
109 other cellular components using brightfield transmission  
110 microscopy without the need of any contrast generation  
111 technique or the use of fluorescent probes.

112 The transport of melanosomes is regulated by signaling  
113 cascades initiated by the binding of specific hormones to  
114 cell-surface receptors, which results in the modulation of  
115 cAMP concentrations in the cytoplasm [19, 20]. Therefore,  
116 one can stimulate melanosome movement toward or  
117 away from the cell center by using appropriate hormones

118 to decrease or increase the concentration of cAMP,  
119 respectively. While melatonin increases the amount of  
120 cAMP triggering aggregation of melanosomes towards the  
121 perinuclear region, MSH (melanocyte stimulating hormone)  
122 reduces cAMP concentration and thus promotes melano-  
123 somes dispersion in the cytoplasm [18]. Importantly, the  
124 observation of bidirectional motion of melanosomes during  
125 aggregation and dispersion indicates that kinesin and dynein  
126 motors alternates in the transport of these cargoes. To  
127 achieve the characteristic distributions of melanosomes  
128 observed during aggregation and dispersion, minus-end  
129 directed motion dominates during aggregation while plus-  
130 end directed motion dominates during dispersion [9].

131 Melanosome transport in *Xenopus* melanophores is well  
132 characterized in terms of molecular motors participating in  
133 the movement. Transport along microtubules is driven by  
134 kinesin-2 [21] and cytoplasmic dynein [22], and myosin-V  
135 transports melanosomes along actin filaments [23]. It is  
136 possible to eliminate the contribution of myosin-V to  
137 melanosome movement by depolymerizing actin filaments  
138 with latrunculin B. The remaining movement is entirely  
139 microtubule dependent and therefore, latrunculin-treated  
140 cells can be used to study the contribution of microtubule  
141 motors cytoplasmic dynein and kinesin-2 to organelle  
142 transport without the contribution of the second transport  
143 system [11].

144 In this work, we analyzed the switch of microtubule  
145 motors during transport of melanosomes in *Xenopus* mel-  
146 anophores by registering trajectories of these organelles  
147 moving along microtubules using a fast and precise  
148 tracking method. We analyzed in detail the intervals of  
149 trajectories showing reversions in the original direction of  
150 transport that can be associated with a switch between the  
151 microtubule-dependent motors responsible for the trans-  
152 port. The experimental data could be explained considering  
153 that reversions are triggered by an elastic collision of the  
154 cargo with obstacles in the cytosol. Our results suggest that  
155 the crowded cytoplasm plays a key role in regulating the  
156 switching of microtubules-dependent motors.

## 157 Materials and Methods

### 158 Cell Culture and Samples Preparation for Imaging

159 Immortalized *Xenopus laevis* melanophores were cultured  
160 as described [24]. In order to track the movement of  
161 individual organelles, the number of melanosomes in the  
162 cell was reduced by treatment with phenylthiourea [11].

163 For microscopy measurements, cells were grown for  
164 2 days on 25-mm round polylysine-coated coverslips  
165 placed into 35-mm plates in 2.5 ml of the medium. Before  
166 observation, the coverslips were washed in serum-free 70%

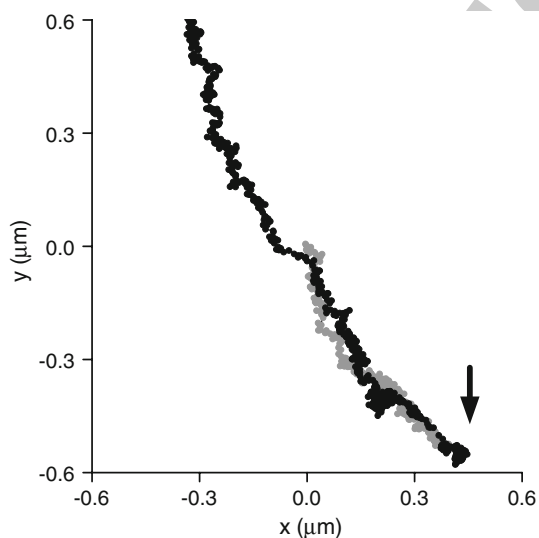


167	L-15 medium and mounted in a custom-made chamber	Reversion-Finder Method	216
168	specially designed for the microscope. The cells were		
169	treated with 10 $\mu$ M latrunculin B (Biomol International,	We define long-term reversions as those segments of the	217
170	Plymouth Meeting, PA) for at least 30 min to depolymerize	trajectories in which the melanosome travels back over the	218
171	actin filaments. Melanophores were stimulated for aggrega-	same path for at least 250 nm (i.e., the average radius of a	219
172	tion or dispersion with 10 nM melatonin or 100 nM	melanosome [24]) during more than 1 s. Importantly, this	220
173	MSH, respectively. Samples were observed between 5 and	distance is much longer than the level of noise commonly	221
174	15 min after stimulation. All measurements were per-	observed in the trajectories ( $\sim 10$ nm). This criterion allows	222
175	formed at 21°C.	excluding spontaneous, short-term reversions due to, e.g.,	223
		thermal jittering or external noise from the analysis.	224
176	Microscope Setup	To automatically detect reversions in trajectories we	225
		scored each pair of a trajectory data points with a param-	226
177	Tracking experiments were carried out in an Olympus IX70	eter $D_{ij}$ defined as $D_{ij} = l_{ij}/d_{ij}$ , where $d_{ij}$ is the point-to-	227
178	microscope using a 60 $\times$ water-immersion objective	point Euclidean distance and $l_{ij}$ is the distance along the	228
179	(numerical aperture = 1.2) under illumination with a tung-	trajectory between the same pair of data points.	229
180	sten-halogen lamp. A CMOS camera (Pixelink, Ottawa,	Points belonging to a reversion will have high $D_{ij}$ val-	230
181	Ontario, Canada) was attached to the video port of the	ues. We constructed an algorithm that select regions with	231
182	microscope for imaging the cells. Movies were registered at	$D_{ij}$ values higher than a threshold set to 10. These regions	232
183	a speed of 100 frame/s. The pixel size was 120 nm and was	were considered part of reversions if they also fulfill the	233
184	set following the criteria discussed in Thompson et al. [25].	requirements presented above. To speed up the computa-	234
185	Experiments are controlled by the acquisition program	tional time without losing critical information, only 1 of	235
186	Globals for Images, program developed at the Laboratory	10 contiguous data points were considered in this initial	236
187	for Fluorescence Dynamics (UCI, Irvine, CA).	analysis.	237
		The procedure described above was implemented using	238
188	Pattern-Recognition Algorithm for Tracking	Matlab (The MathWorks, Natick, MA) routines.	239
189	Melanosomes		
		<b>Results and Discussion</b>	240
190	The pattern-recognition tracking routine is described in	Reversions in Trajectories of Melanosomes During	241
191	[26]. Briefly, the program starts the tracking routine dis-	Aggregation and Dispersion	242
192	playing the first frame of the image stack under analysis.		
193	The operator chooses the target melanosome by simply	Melanophores were incubated with latrunculin B as indi-	243
194	clicking on top of its image. Doing so, the program sets the	cated in Materials and Methods in order to depolymerize	244
195	initial coordinates of the melanosome and generates an	the actin filaments. Aggregation or dispersion of melano-	245
196	intensity pattern that consists of the average intensity	somes was induced by addition of melatonin and MSH,	246
197	obtained from the first 10 frames of a region containing the	respectively. We recorded 10–20 movies per cell after	247
198	melanosome image. This pattern is stored in the computer	incubating them for at least 5 min and no more than 15 min	248
199	memory to be used during the calculation of the melano-	in the presence of either melatonin or MSH.	249
200	some position through the image stack. In the following	To have a precise description of melanosome motion we	250
201	frames, the pattern is shifted around the position deter-	used a fast tracking routine that allow us to recover the	251
202	mined for the particle in the previous frame and a	melanosome position with 2 nm precision and 10 ms	252
203	parameter $\sigma$ that scores the absolute intensity differences	temporal resolution. One of the advantages of the method	253
204	between image and pattern is calculated. The $\sigma$ value will	is that it does not assume any intensity distribution for the	254
205	be minimal when the image in the frame matches the	particle image and thus it can be used to track, for example,	255
206	pattern features, thus it can identify the particle of interest	micrometer-sized particles such as melanosomes. We have	256
207	from other structures present in the frame [27]. Another	previously used this method to study melanosome transport	257
208	advantage is that the method does not need a theoretical	along microtubules [26] and actin filaments [29].	258
209	expression for the intensity distribution of the particle as	Three hundred total trajectories of melanosomes moving	259
210	Gaussian fitting methods [28]. The particle position is	after stimulating cells for aggregation or dispersion were	260
211	calculated with sub-pixel resolution by determining the	determined from the movies using the pattern-recognition	261
212	position corresponding to the minimum value of $\sigma$ with a	algorithm described above and were classified according to	262
213	parabolic interpolation. By using this method, we could	the direction in which the melanosome was initially	263
214	recover the position of 500 nm particles with 2 nm preci-		
215	sion and 10 ms temporal resolution [26].		

264 moving. Microtubules minus-ends are attached to the  
 265 centrosome, which is located near the nucleus. Therefore,  
 266 those trajectories showing a preferential direction toward  
 267 the perinuclear region correspond mainly to organelles  
 268 transported by cytoplasmic dynein while those moving  
 269 toward the cell periphery are generally transported by  
 270 kinesin 2. In this article, we assigned inward and outward  
 271 motion of melanosomes to transport driven by dynein and  
 272 kinesin, respectively. It is important to mention that while  
 273 most of inward and outward trajectories will in fact cor-  
 274 respond to the assigned motors, we cannot be sure about  
 275 the identity of the motor in every analyzed trajectory.

276 Figure 1 shows an example of a trajectory obtained for a  
 277 melanosome moving after stimulating the cells with mel-  
 278 atonin. In this particular trajectory, the organelle moves  
 279 toward the cell nucleus probably by the action of cyto-  
 280 plasmic dynein and it suddenly reverts its direction being  
 281 transported backwards by kinesin-2 toward the cell  
 282 periphery. A similar behavior can also be observed in the  
 283 supplemental movie 1, which shows a representative  
 284 tracking experiment.

285 To quantify the probability of reversion during in vivo  
 286 transport, we constructed a reversion-finder algorithm that  
 287 allow us to detect those regions of the trajectories in which  
 288 organelles originally moving in a given direction show  
 289 afterwards a continuous sustainable motion in the opposite  
 290 direction for more than 250 nm during at least 1 s (see  
 291 Materials and Methods). These arbitrary thresholds were  
 292 set to exclude from the analysis dubious reversions that  
 293 could be consequence of other processes.



**Fig. 1** Reversions in melanosome trajectories. Example of a trajectory obtained for a melanosome that was initially transported by cytoplasmic dynein (gray) and suddenly reverts its direction moving finally toward the cell periphery by the action of kinesin-2 (black). The switch point is marked with an arrow. The trajectory lasted 9.8 s

294 In total, 30 and 40% of the trajectories determined for  
 295 melanosomes during dispersion and aggregation, respec-  
 296 tively, presented long-range reversions indicating that they  
 297 are frequent processes during organelle transport (Table 1).

298 We determined the relative reversion probability of  
 299 organelles driven by cytoplasmic dynein by calculating the  
 300 ratio of minus-to-plus end reversions relative to the number  
 301 of minus-end trajectories. A similar procedure was followed  
 302 to calculate the reversion probability of melanosomes  
 303 transported by kinesin-2. Table 1 shows that the reversion  
 304 probability of each of the motors is similar during aggrega-  
 305 tion and dispersion suggesting that the reversion  
 306 mechanism is independent from the stimulation condition of  
 307 the cells.

308 On the other hand, the probability of switching motors  
 309 for melanosomes initially driven by kinesin-2 is approxi-  
 310 mately twice the reversion probability of those driven by  
 311 cytoplasmic dynein. In contrast to the increasing amount of  
 312 information regarding the structure, regulation and bio-  
 313 physical properties of kinesin-1 and cytoplasmic dynein,  
 314 little is known about the properties of kinesin-2 obscuring  
 315 the interpretation of this result. Thus, the causes of this  
 316 higher reversion probability of kinesin driven melanosomes  
 317 remain to be determined.

#### A Model for Wave-Shaped Reversions

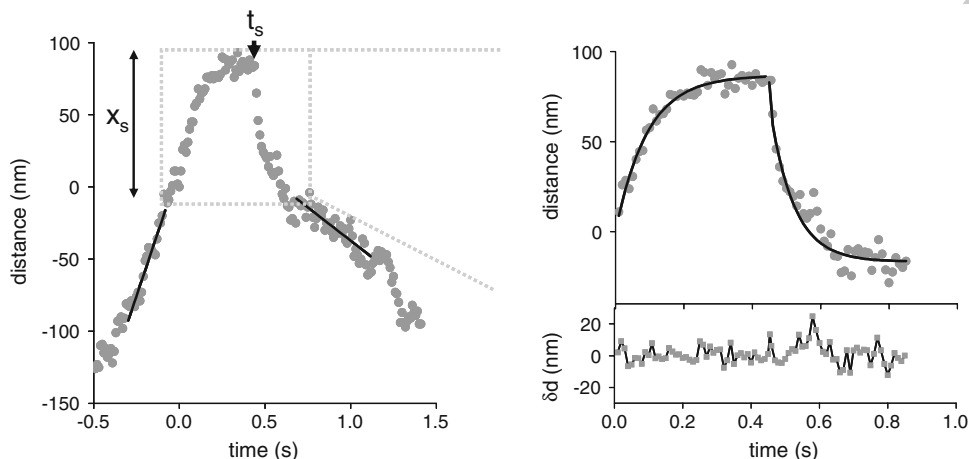
318 Figure 2 shows the distance traveled by a melanosome as a  
 319 function of time during a representative reversion. The right  
 320 panel of the figure shows that the speed of the melanosome  
 321 before the reversion slowly decreases with time approaching  
 322 zero then, the melanosome returns over the same path  
 323 moving initially at a very high speed and slowing down with  
 324 time. We found that 70% of the reversions present these  
 325 wave-shaped distance versus time plots. The instantaneous  
 326 speed during the first 30 ms after the reversal point ( $t_s$  in  
 327 Fig. 2) was in average  $1,900 \pm 150$  nm/s. This high speed  
 328 cannot be explained by only considering the transport by  
 329 microtubule-dependent molecular motors, which move at  
 330 speeds in the range of 250–1,000 nm/s in melanophores of  
 331 *Xenopus laevis* [11, 26].

332 Wave-shaped reversions present characteristics similar  
 333 to those observed for trajectories of beads attached to  
 334 motor molecules moving along microtubules in optical  
 335 trapping experiments. The motor moves the beads far from  
 336 the equilibrium position at the center of the trap while the  
 337 trap applies an opposing force making the motor to slow  
 338 down and stop when reaching stall force. It then detaches  
 339 from the track and the bead moves fast to the equilibrium  
 340 position at the trap center (see for example, [30, 31]).

341 Similarly, we propose that intracellular components in  
 342 the overcrowded cytoplasm such as cytoskeletal filaments  
 343 can act as obstacles which elastically interact with motor-  
 344

**Table 1** Statistics of melanosome transport

	Plus-end directed melanosomes (%)	Trajectories presenting reversions (%)	Wave shape-reversions (%)	Reversion probability for kinesin	Reversion probability for dynein
Dispersion	74	40	73	0.41	0.25
Aggregation	40	30	70	0.46	0.17



**Fig. 2** Reversions characterization. Representative region of a trajectory showing a wave-shaped reversion. The distance traveled by the organelle was calculated as described in Materials and Methods and is represented as a function of time. The stall distance  $x_s$  and the reversion time  $t_s$  are marked with arrows. Continuous lines

show the processive motion of the organelle before and after the interaction with the obstacle. Continuous lines in the right panel show the fitting of Eqs. 3 and 4 to the experimental data. The bottom panel shows the residuals obtained from the fitting

345 driven organelles. Melanosomes are in average 500 nm-  
 346 diameter suggesting a high probability of encountering  
 347 obstacles in the overcrowded cytoplasm.

348 The stochastic motion of the complex organelle-motor is  
 349 governed by the Langevin equation [32, 33], which takes  
 350 into account the Brownian force produced by collisions  
 351 with other molecules as well as the viscous drag and other  
 352 external forces. If the time scale of the Brownian force is  
 353 much smaller than the time scale of other forces that might  
 354 be involved in the motion of the motor-organelle complex,  
 355 the temporal evolution of the position of the organelle can  
 356 be described at long enough times in terms of a non-sto-  
 357 chastic equation [34]. In this regime, the forces acting on  
 358 the organelle are given by,

$$m \frac{d^2x}{dt^2} = -\gamma \frac{dx}{dt} + F_{\text{motor}} + F_{\text{ext}} \quad (1)$$

360 where  $x$  is the average distance traveled by the organelle,  $m$   
 361 is its mass,  $\gamma$  is a drag coefficient,  $F_{\text{motor}}$  is the force exerted  
 362 by the motors driving the organelle and  $F_{\text{ext}}$  is an external  
 363 force.

364 When the number of active motors and the viscous drag  
 365 do not change during the transport, the speed of the orga-  
 366 nelle reaches a constant value  $v'_d$  in the absence of external  
 367 forces. The system described by Eq. 1 will reach this

stationary or Stokes regime [35] after a time  $t_r$  equal to  $m/\gamma$  [36]. This behavior is observed in experimental trajectories as periods of rectilinear motion of constant velocity, called “runs” [37].

Considering that the organelle moving at  $v'_d$  encounters an obstacle and that their interaction is elastic:

$$m \frac{d^2x}{dt^2} + \gamma \frac{dx}{dt} + \kappa x = F_{\text{motor}} \quad (2)$$

where  $\kappa$  is the elastic constant characterizing the interaction.

Taking into account that the average melanosome diameter is 500 nm [24], its density is  $\sim 1.2$  g/ml [38] and  $\gamma \sim 10^{-8}$  Ns/m for the drag coefficient in water [36],  $m/\gamma$  would be  $\sim 10^{-2}$   $\mu$ s, which is significantly lower than the temporal resolution of the tracking method we used. Since the viscosity of the cytosol is higher than that of aqueous solutions (see references in [39]),  $m/\gamma$  will be even lower than the value calculated above indicating that  $t_r$  is achieved instantaneously considering the time resolution of our experiments. In this condition, the dynamics of the system is reduced to an overdamped behavior.

The distance traveled by the organelle before the switch of direction is given by the solution of Eq. 2, which in the overdamped limit reduces to:

$$x(t) = \frac{\gamma}{\kappa} v_d \left[ 1 - e^{-\frac{\kappa t}{\gamma}} \right] \quad 0 < t < t_s \quad (3)$$

392 Where  $x(t = 0) = 0$ . This equation shows that the moving  
393 melanosome slows down as a consequence of the opposing  
394 force introduced by the obstacle. This increasing, elastic  
395 force will also increase the probability of detachment of the  
396 active motors from the track [40, 41]. We postulate that at  
397 the reversal position  $x_s$ , the active motors detach from the  
398 microtubule and the melanosome starts moving in the  
399 opposite direction driven by opposed-polarity motors.

400 This proposition is based also on the facts that: (1) there  
401 are no evidences in the literature that microtubule-depend-  
402 ent motors responsible for melanosomes transport can  
403 move backward along microtubules and, (2) dynein and  
404 kinesin are not active simultaneously as was also men-  
405 tioned in the “Introduction” section. The elastic energy  
406 released after the detachment of the first motors will also  
407 contribute to the backward motion of the melanosome.

408 By following an analysis similar to that described  
409 before, it can be demonstrated that the distance traveled by  
410 the organelle after the switch of direction is given by:

$$x(t) = \left[ x_s + \frac{\gamma}{\kappa} v'_d \right] e^{-\frac{\kappa}{\gamma}(t-t_s)} - v'_d \frac{\gamma}{\kappa} \quad t_s < t \quad (4)$$

412 where  $t_s$  is the switch time,  $x(t = t_s) = x_s$  and  $v'_d$  is the  
413 speed of the motors that drive the organelle after the  
414 reversion.

415 Equations 3 and 4 consider that if several motors are  
416 driving the organelle, they would act in concert transport-  
417 ing the cargo together until detaching from the track or  
418 after attaching back to it. Welte et al. [14] measured in  
419 vivo the stall force of lipid droplets transported in *Dro-*  
420 *sophila* embryos during different stages of development.  
421 They determined that the stall force changes among these  
422 stages in a quantized fashion, consistent with a variation in  
423 the number of active motors driving the droplets. They did  
424 not observe at any of these development stages a distri-  
425 bution of stall forces compatible with motors detaching  
426 one-by-one. This result is different from what it would be  
427 expected according to a model describing in vitro transport  
428 of cargos driven by multiple copies of motors under load  
429 [42]. The reason that might explain this divergence is that,  
430 as it is indicated by Welte et al. [14], specific proteins  
431 would enforce the coordination of same polarity motors in  
432 vivo and control the number of actively engaged motors.

### 433 Trajectories Analysis

#### 434 *Wave Shape-Reversions Are Characteristic of Fast-Moving* 435 *Organelles*

436 We fit Eqs. 3 and 4 to the segments obtained before and  
437 after the reversal position, respectively, considering the

438 same set of  $\gamma$  and  $\kappa$  for each analyzed reversion. These  
439 parameters characterize the interaction of the organelle  
440 with the obstacle and the medium and, in a first approxi-  
441 mation, can be considered to be the same before and after  
442 the encounter.

443 The fitting of the equations was done by minimizing the  
444 function:  $S = \sum_i (X(t_i) - x_i)^2$ , where  $X(t_i)$  is the prediction  
445 of the model for the position of the melanosome at time  $t_i$ ,  
446 and  $x_i$  is the corresponding experimental data. Fitting  
447 routines were written within the MatLab programing  
448 environment (The MathWorks, Natick, MA), using the  
449 genetic algorithm toolbox.

450 Figure 3a shows the distribution of  $v'_d$  obtained by fitting  
451 Eqs. 3 and 4 to wave-shaped reversions during aggregation  
452 and dispersion. We then classified  $v'_d$  values according to  
453 the motor that was driving the organelle before the rever-  
454 sion and found that the distribution of this parameter was  
455 not significantly different for kinesin- and dynein-driven  
456 melanosomes (>0.77 confidence level, Kolmogorov–Smir-  
457 nov test [43]).

458 The histogram of  $v'_d$ —which is obtained by only  
459 analyzing regions of the trajectories presenting rever-  
460 sions—can be compared to the speed distribution we have  
461 previously obtained by analyzing the complete trajectories  
462 of melanosomes driven by dynein and kinesin [26]. Briefly,  
463 in that previous work we found that melanosomes speed  
464 followed a multi-peak distribution that could be explained  
465 considering that this parameter depends linearly on the  
466 number of active motors. We also found that melanosomes  
467 transported by a single dynein or kinesin molecule were the  
468 predominant population during aggregation and dispersion  
469 and that their speed was  $\sim 250$  nm/s.

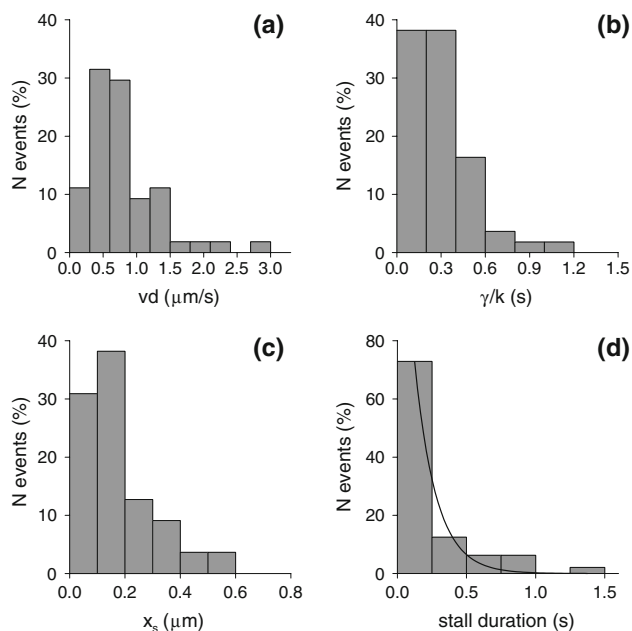
470 We cannot discard the presence of multiple peaks in the  
471  $v'_d$  distribution; however, we do not have enough infor-  
472 mation to statistically analyze this possibility since  
473 reversals represent only a small region of the trajectories.

474 On the other hand, Fig. 3a shows that the  $v'_d$  distribution  
475 presents a maximum at  $\sim 0.5$   $\mu\text{m/s}$  indicating that wave-  
476 shaped reversions are detected only in those regions of the  
477 trajectories in which melanosomes are moving fast and are  
478 probably transported by more than 1 copy of the active  
479 motor.

#### 480 *The Switch Occurs Once the Motor Reaches Stall* 481 *Condition*

482 Equations 3 and 4 show a strong dependence with  $\gamma/\kappa$  ratio  
483 but not with the independent parameters and thus, they  
484 cannot be recovered individually from the fitting.  $\gamma/\kappa$  rep-  
485 represents the time constant of the elastic response, i.e., low  
486 values of  $\gamma/\kappa$  imply that the obstacle stiffness is high and  
487 thus the organelle stops shortly after their interaction. This





**Fig. 3** Distribution of parameters characterizing reversions in melanosomes transport direction. **a** Distribution of melanosomes speed before interacting with the obstacle. **b** Distribution of characteristic times of the interaction with the obstacle. The mean  $\gamma/k$  value is  $290 \pm 30$  ms.  $v_d$  and  $\gamma/k$  values were obtained by fitting Eqs. 3 and 4 to wave-shaped reversions in melanosomes trajectories. **c** Distribution of stall distances calculated by using Eq. 5. The average stall distance was  $200 \pm 30$  nm. **d** Distribution of stall duration. Stall duration is defined as the interval of time the organelle spends in stall conditions before it reverts its direction of motion. The continuous line corresponds to the fitting of an exponential-decay function, with a characteristic time constant of  $160 \pm 2$  ms. Characteristic parameters of the distributions are expressed as the mean  $\pm$  standard error. Histograms were constructed setting the bin size to the value determined by the commonly used criterion [53]:  $\text{bin size} = 3.49 \frac{s}{n^{1/3}}$  where  $n$  is the number of data and  $s$  is an estimate of the standard deviation

488 parameter shows a wide distribution with a mean value of  
 489  $290 \pm 30$  ms (Fig. 3b). The average  $\gamma/k$  value obtained in  
 490 this work is in the order of the values expected according to  
 491 the typical elasticities of intracellular elements [36, 44, 45]  
 492 and of  $\gamma$  values determined in the cytoplasm [39]. This  
 493 result confirms that the response time of the elastic inter-  
 494 action is higher than the time resolution of the tracking  
 495 method allowing the direct observation of the time evolu-  
 496 tion of the system.

497 This result is not compatible with a model in which  
 498 reversions are initiated by a sudden stop of the active  
 499 motor. If this was the case, wave-shaped distance versus  
 500 time plots would be the consequence of the motion of the  
 501 melanosome caused by inertia once the motor stops mov-  
 502 ing. The time constant of the inertial effects is lower than  
 503  $10^{-2}$   $\mu$ s as we discussed previously. If we also consider the  
 504 stiffness of the motor to stretching  $\kappa_e$  [45, 46] the time  
 505 constant would be  $\gamma/\kappa_e \sim 0.03$  s. These values are

506 significantly lower than that measured in this work  
 507 (Fig. 3b) and thus this hypothesis can be ruled out.

508 According to Eq. 2, we can define the stalling distance  
 509  $x_s$  as the theoretical distance traveled by the melanosome  
 510 from the moment it contacts the spring-like obstacle until  
 511 its maximal compression:

$$x_s = \frac{\gamma}{\kappa} v_d \quad (5)$$

513 Figure 3c shows the distribution of stalling distance  
 514 determined by using this equation. The average value was  
 515  $200 \pm 30$  nm.

516 We can also estimate the distance traveled by the mel-  
 517 anosome since it starts interacting with the elastic obstacle  
 518 until it switches the direction, i.e., the switch distance. The  
 519 algorithm described in “Materials and Methods” section  
 520 can easily recognize the reversion position and the initial  
 521 point of the interaction can be detected by the sudden  
 522 decrease of the organelle speed.

523 We calculated the ratio between the stall and switch  
 524 distances and observed that this ratio is not significantly  
 525 different from 1 (Student’s t test, 95% confidence). This  
 526 result suggests that reversions can only occur once the  
 527 motors responsible for the motion reach stall conditions.

### The Detachment of the Organelle Is Stochastic

528 To further characterize the mechanism that triggers the  
 529 exchange of motors we measured the stall duration, i.e., the  
 530 period that the organelle spends in stall conditions before it  
 531 reverts its direction. Since the start of a plateau was not  
 532 always sharply defined because some curves approached it  
 533 asymptotically, we arbitrarily determined the stall duration  
 534 from the moment the organelle travel at speeds lower than  
 535 50 nm/s to the switch time  $t_s$ .  
 536

537 Figure 3d shows that the stall duration distribution fol-  
 538 lows an exponential-decay function with a characteristic  
 539 time of  $0.16 \pm 0.02$  s. This behavior indicates that the  
 540 detachment probability during stall is constant as would be  
 541 expected according to the stochastic nature of the switch.

542 This result agrees with those obtained in optical trapping  
 543 experiments of kinesin-attached beads moving along  
 544 microtubules: Coppin et al. [41] also observed an exponen-  
 545 tial-decay distribution for the stall duration measured when  
 546 the beads are far from the trap center. Interestingly, these  
 547 authors also verified that the dissociation rate of the motors  
 548 from the track is independent of load once the motor stalls.

549 Our model is consistent with the observation of the  
 550 independence of the reversion probability with the stimu-  
 551 lation state of melanophores (Table 1) since reversions  
 552 only depend on the probability of encountering an obstacle.  
 553 This probability would not depend on the stimulation state  
 554 of melanophores since the viscoelastic properties of the

555 cytoplasm of *Xenopus laevis* melanophores are not signif- 582  
 556 icantly different during aggregation and dispersion as 583  
 557 suggested by preliminary results from our group using 584  
 558 single particle tracking microrheology (for a description of 585  
 559 this technique see [47, 48]). 586

#### 560 *The Second Motor Attaches to the Track After Complete* 561 *Relaxation of the Elastic Interaction*

562 In half of the analyzed wave-shaped reversions, the second 590  
 563 motor attaches to the track once the melanosome stops 591  
 564 moving (i.e.,  $v'_d = 0$   $\mu\text{m/s}$ ). Visual examination of these 592  
 565 reversions shows that after elastic relaxation, melanosomes 593  
 566 generally stay still for a short period of time ( $<0.2$  s) and, 594  
 567 immediately after that, almost all melanosomes start 595  
 568 moving in the opposite direction in a stationary fashion 596  
 569 (95%). The mean speed of this processive motion was 597  
 570  $300 \pm 30$  nm/s, which is compatible with a melanosome 598  
 571 moving with a single motor copy. 599

572 Since the fast, backward motion usually takes  $\sim 300$  ms 600  
 573 to complete and occurs near the microtubule along which the 601  
 574 melanosome was originally moving, there is a high proba- 602  
 575 bility that this microtubule is very close to the organelle and 603  
 576 thus active transport will resume on the same microtubule. 604  
 577 However, in some cases melanosomes could reattach to a 605  
 578 different microtubule. If the second microtubule is located 606  
 579 close to the initial microtubule, the reversion will be detected 607  
 580 by the reversal finding algorithm. Otherwise this motion will 608  
 581 be observed not as a reversion but as a change of direction.

Therefore, we are probably underestimating the frequency of 582  
 events involving switches of microtubule motors. 583

On the other hand, in the last year it has been demon- 584  
 strated that myosin-V diffuses along microtubules [49] and 585  
 enhances significantly the processive run length of kinesin- 586  
 1 when both motors are present on the same cargo. These 587  
 authors proposed that myosin-V can act as a tether pre- 588  
 venting kinesin from diffusing away from the track. In the 589  
 particular case of melanophores, it might be possible that 590  
 myosin-V could also tether melanosomes to microtubule; 591  
 however, there are not experimental evidences to support 592  
 this hypothesis. 593

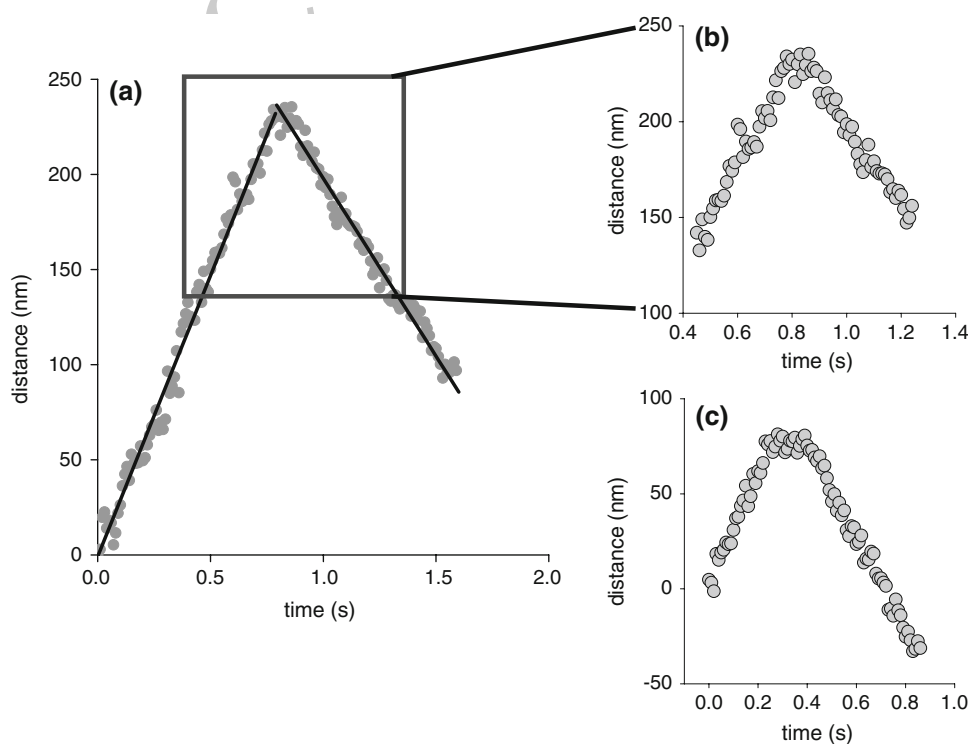
30% of the data shows  $v'_d$  values with a similar distri- 594  
 bution of that observed for  $v'_d$  (not shown). The rest of the 595  
 data could not be classified because they did not show a 596  
 satisfactory fitting after the reversion point. This could be 597  
 probably consequence of other processes affecting the 598  
 trajectories of the organelles (e.g., a collision with a second 599  
 obstacle). 600

#### 601 Peak-Shaped Reversions

In this work, we have analyzed wave-shaped reversions 602  
 which, as we mentioned before, are characteristic of fast- 603  
 moving organelles. We also detected some reversions 604  
 (24%) which show distinct properties with respect to wave- 605  
 shaped reversions (Fig. 4). In these peak reversions, the 606  
 speed was approximately constant before and after the 607  
 switch time. Interestingly, these reversions were observed 608

**Fig. 4** Peak-shaped reversions.

**a** Distance traveled by a melanosome as a function of time for a representative peak-shaped reversion. The speeds of the organelle before and after the switch were 300 nm/s and 200 nm/s, respectively. **b** Zoom of the reversion showed in (a) near the switch point. The stall duration is  $\sim 0.1$  s. **c** Simulated peak-shape reversion. The trajectory was obtained using Eqs. 3 and 4 with parameters values  $v'_d = 320$  nm/s,  $\gamma/\kappa = 0.03$  s,  $v'_d = -250$  nm/s and stall duration equal to 160 ms, which corresponds to the characteristic stall duration determined in Fig. 3d. Random noise with the same distribution and amplitude of the experimental one was added to the simulated data to make the comparison easier



609 for organelles moving at average speeds of  $320 \pm 40$  nm/s,  
610 compatible with transport with a single motor copy [26].

611 These distinct characteristics may suggest that a different  
612 switching mechanism is responsible for peak reversions.  
613 However, Eq. 5 shows that the stall distance decreases with  
614 both  $v_d$  and  $\gamma/\kappa$ . We observed that by using a set of  $v_d$  and  $\gamma/\kappa$   
615 values slightly lower than the mean parameters obtained  
616 from wave-shaped reversions we do recover plots with  
617 similar properties of those corresponding to experimental  
618 peak-shaped reversions. For example, the simulated rever-  
619 sion showed in Fig. 4c resulted from using a  $\gamma/\kappa$  value that is  
620 10% lower than the mean value determined from Fig. 3b.  
621 Importantly, the  $v_d$  and  $\gamma/\kappa$  set of values that resulted in the  
622 peak-shaped reversion shown in this example are within the  
623 range of the experimental values determined for waved-  
624 shaped reversions (Fig. 3a and b).

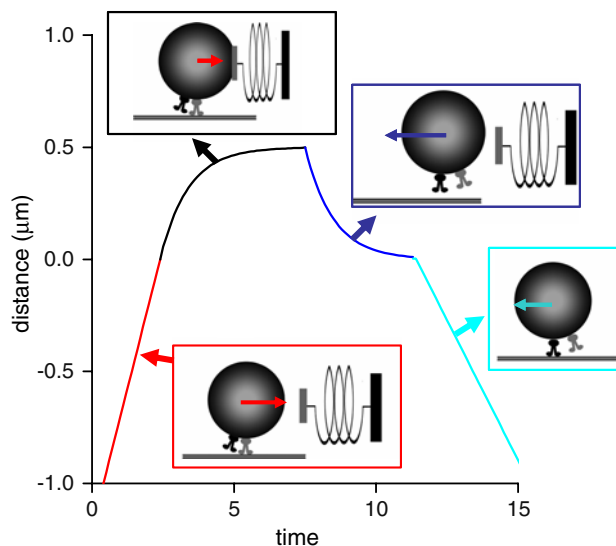
625 We would have expected a higher proportion of peak-  
626 shape reversions since most of the time melanosomes are  
627 transported at  $\sim 250$  nm/s [26]. However, these organelles  
628 are expected to show shorter stall distances as we dem-  
629 onstrated above and shorter run lengths as it is observed  
630 during the transport of lipid droplets [37]. Thus, these  
631 reversions are probably not detected by the reversion-finder  
632 algorithm.

### 633 Concluding Remarks

634 In this work, we have studied the characteristics of trajec-  
635 tories of melanosomes moving along microtubules and  
636 focused on those regions showing reversions in the direction  
637 of transport to get insight into the mechanism that triggers  
638 the switch between microtubule motors of opposed polarity.

639 Most of the studied reversions could be explained  
640 according to the model schematically represented in Fig. 5.  
641 According to our results, an obstacle in the cytosol intro-  
642 duces an extra load to the active motors that make them slow  
643 down and stop when the force generated by the motor  
644 molecules equals the opposite force, i.e., the drag force plus  
645 the force exerted by the spring-like obstacle. Once this stall  
646 condition is reached, the active motor stochastically detaches  
647 from the microtubule track. In this condition, the free  
648 organelle moves backwards at high initial speeds and slows  
649 down as a consequence of the friction with the medium.  
650 Once the organelle stops moving and, since it is near the  
651 track, opposite polarity motors can attach to the microtubule  
652 and transport the organelle in the opposite direction.

653 Our model considers that the melanosome behaves as a  
654 stiff bead. In a recent article, Guo et al. [50] measured the  
655 Young modulus of melanosomes and found that they are  
656 considerably higher than the modulus of organelles with  
657 cytoplasm ( $<1$  MPa) and approaching values of the mod-  
658 ulus of protein crystals ( $\sim 100$  MPa). Despite these  
659 measurements were done on melanosomes of human



**Fig. 5** Scheme of the switch mechanism triggered by an elastic interaction between a motor-driven melanosome and an obstacle. The organelle travels along the microtubule at a constant speed (red) until it approaches to the obstacle and start interacting with it (black). As a consequence, the melanosome slows down. Once the opposing force equals the stall force of the motor, it detaches from the track and start moving against the viscous drag (blue). After relaxation of the system, the opposite polarity motor attaches to the track and continues moving in opposite direction (cyan). The scheme considers the case in which  $v'_d = 0$ . To simplify the scheme, a single copy of each family of motors is represented

660 retinal pigment epithelium, their properties are not expect-  
661 ed to be very different from those of *Xenopus laevis*  
662 melanophores and thus the latter can be considered as stiff  
663 organelles.

664 The postulated model gives an explanation to the trigger  
665 of the switch of motors, but does not explain how the  
666 coordination of the motors is achieved. We can hypothesize  
667 that the collision determines the end of a run of a given  
668 polarity motor and the regulatory proteins (e.g., dynactin)  
669 involved in the coordination mechanism would define  
670 which motor will be turned on after this collision and keep  
671 motors of opposite polarity in the inactive state.

672 The model postulated in this work agrees with previous  
673 results from Gross et al. [37] who postulated that there is a  
674 mechanism that ends runs of dynein-driven organelles  
675 before the motors fall off the microtubules which acts with  
676 a constant probability per unit distance, and is typically  
677 coupled to a switch in travel direction. They also proposed  
678 that a similar mechanism governs plus-end motion, and its  
679 regulation controls the net direction of transport.

680 Also, the model gives a possible explanation to previous  
681 results from Rogers et al. [24] who observed that mela-  
682 nosomes transport in vitro process over longer distances  
683 that in vivo. The main difference between both assays is  
684 the environment in which the transport occurs: while in

685 vitro assays are done in aqueous solutions, the cytoplasm is  
686 a crowded, viscoelastic medium and the probability of  
687 finding obstacles that decreases the run length and increa-  
688 ses the probability of switching motors is higher.

689 As it is explained by Gross [1], the frequent changes of  
690 directions observed for several cargoes in the cell cyto-  
691 plasm seems to be an inefficient mechanism to move the  
692 cargoes to a given, particular location. In the same way,  
693 transport toward a particular region of the cell will be  
694 delayed if the encounters with obstacles trigger the switch  
695 of microtubule-dependent motors. However, in the absence  
696 of a reversion mechanism, motor-driven organelles would  
697 remain in a fixed position after encountering an obstacle  
698 until the obstacle or the complex moves away. Since the  
699 cytoplasm is an overcrowded medium with a high concen-  
700 tration of obstacles, the efficiency and regulation of the  
701 transport will be highly impaired in this situation.

702 On the other hand, if the encounter with the obstacle  
703 triggers the exchange of motors, the cargoes would back up  
704 and move in the opposite direction. Eventually, it may  
705 switch to a different microtubule that crossed the path and  
706 find an alternative route to evade the obstacle. This mech-  
707 anism would contribute to avoid “traffic jams” such as those  
708 observed in axons, which have been proposed as the crucial  
709 injury in a range of neurodegenerative diseases [51, 52].

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714 P. Balenzuela.

## 715 Supporting File

716 Movie 1. Typical tracking experiment showing a reversion  
717 in a melanosome trajectory. The movie shows a  
718  $24 \times 24 \mu\text{m}^2$  region of the cell (in pseudo color); mela-  
719 nosomes are observed in yellow. The black line shows the  
720 trajectory recover by the pattern-recognition algorithm.  
721 The total duration of the tracking experiment is 20 s.

722

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