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Optical Nanoparticle Sorting Elucidates Synthesis of Plasmonic Nanotriangles

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ABSTRACT

We investigate the optical and morphological properties of gold nanoparticles grown by reducing a gold salt with Na₂S. Lasers are tuned to the observed plasmon resonances, and the optical forces exerted on the nanoparticles are used to selectively print individual nanoparticles onto a substrate. This enables us to combine dark-field spectroscopy and scanning electron microscopy to compare the optical properties of single nanoparticles with their morphology. By arresting the synthesis at different times, we are able to investigate which type of nanoparticle is responsible for the respective resonances. We find that thin Au nanotriangles are the source of the observed NIR resonance. The initial lateral growth of these triangles causes the plasmon resonance to

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3 redshift into the NIR, whereas a subsequent thickening of the triangles and a concomitant
4 truncation lead to a blueshift of the resonance. Furthermore we find that the nanotriangles
5 produced have extremely narrow linewidths (187 ± 23 meV), show nearly isotropic scattering and
6 are stable for long periods of time. This shows their vast potential for applications such as *in-vivo*
7 imaging and bio-(chemical) sensing. The method used here is generally applicable to other
8 syntheses, and shows how complex nanostructures can be built up on substrates by selectively
9 printing NPs of varying plasmonic resonances.
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21 KEYWORDS

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24 Optical printing, gold nanoparticles, cancer therapy, nanoshells, plasmon, nanoparticle sorting
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30 TEXT

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32 Noble metal NPs possess remarkable properties that have led to applications in fields ranging
33 from solar energy harvesting¹⁻² and photodetection³⁻⁶ to bio-(chemical) sensing,⁷ nanomedicine⁸
34 and cancer therapy.⁹⁻¹⁰ One of the main advantages these astounding particles have is the
35 collective oscillation of their conduction band electrons, generally known as a localized surface
36 plasmon (LSP). This leads to an enhanced optical cross section whose resonance position can be
37 tuned by modifying not only the material, but also the size and shape of the NP. Especially
38 interesting for biomedical applications are NPs whose resonance is shifted to the near infrared
39 (NIR) window, a region of minimum absorption in biological tissues between 750 and 900 nm.¹¹
40 Nanoshells, comprising a thin metallic film surrounding a dielectric core have proven especially
41 good for this, as the resonance position can be tuned by controlling the ratio of the core to shell
42 radius.¹² One of the proposed synthesis methods of nanoshells consists in the reduction of an Au
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3 salt, H₂AuCl₄, by Na₂S (Figure 1a).¹³ It was said that this leads to spherical nanoshells, less than
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5 40 nm in size, with resonances in the NIR window, a strong sensitivity to the surrounding
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7 medium and an extremely narrow line width.¹⁴⁻¹⁵ The actual existence of these NPs, however, is
8
9 a matter of dispute, as some researchers argue that the observed features are either caused by
10
11 non-spherical Au NPs¹⁶⁻¹⁷ or NP-aggregates.¹⁸⁻¹⁹ Settling this dispute is no easy feat and requires
12
13 a combination of strategies. Firstly, the synthesis needs to be carried out in a manner that one can
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15 investigate the NPs present at any given time during the procedure, as their optical properties
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17 change during their growth process. Next, the NPs at each time step need to be separated and
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19 their optical properties correlated to their morphology. In case there are many types of particles
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21 in the dispersions then one needs a large amount of statistics to be able to draw significant
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23 conclusions. Optical forces present an extremely beneficial method for the separation of the
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25 individual NPs. Using the strong scattering force acting on plasmonic NPs at resonance it has
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27 been possible to print such particles at specific locations onto hard substrates,²⁰⁻²² onto
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29 phospholipid membranes²³⁻²⁵ or even to inject them into cells.^{24, 26} This effect, termed optical
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31 printing, uses a laser tuned to the resonance maximum of the plasmonic NP and consequently the
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33 strong axial optical force to propel the NPs along the Poynting vector. Only nanoparticles with
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35 particle plasmons in resonance with the laser wavelength experience a force strong enough to
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37 propel them towards the surface where they are then affixed. Once a nanoparticle is printed onto
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39 the substrate, the printed nanoparticle then interacts with the laser beam, preventing subsequent
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41 printing of additional nanoparticles.^{22, 27} Additionally, it has been shown that in a dispersion
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43 comprising two kinds of metallic NPs, each kind could be propelled in a different direction based
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45 on their localized surface plasmon resonance (LSPR), a first step towards NP sorting.²⁸
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3 In this manuscript we develop a strategy to settle the aforementioned dispute. The synthesis of
4 the NPs is carried out and arrested at specified times by addition of Na₂S. The NPs are then
5 separated and fixed at specific locations on a substrate through an all-optical approach. By
6 employing multiple lasers we can selectively print NPs whose LSPRs overlap with the laser
7 wavelength. We can then combine dark field Rayleigh scattering and scanning electron
8 microscopy (SEM) to compare the optical properties of the printed NPs with their morphology.
9 Through this method we found that the NPs responsible for the strong NIR peak are not
10 nanoshells, but actually thin Au nanotriangles.¹⁶⁻¹⁷ Additionally, we found that unlike in large
11 nanoprisms, their optical scattering is nearly isotropic, likely due to their small size.²⁹ Moreover,
12 the nanotriangles possess a very narrow line width and also show a strong sensitivity to the
13 refractive index of the surrounding medium. All of these features are prerequisites for sensing
14 applications, for which these nanotriangles are highly adept. Finally, we propose a mechanism to
15 explain the observed dynamics of the UV-Vis spectra of the NP dispersions during synthesis.
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36 RESULTS AND DISCUSSION

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38 Fabrication of the Au NPs was carried out in a two-step reaction according to previously
39 published methods.¹³⁻¹⁴ Briefly, 10 mL of a 2 mM HAuCl₄ solution were mixed with 12 mL of
40 fresh 1mM Na₂S solution. After two minutes 2 mL of 1 mM Na₂S solution were added, giving a
41 total molar S/Au ratio of 0.65. The reaction was then allowed to evolve for 3 hours and
42 monitored every two minutes by means of UV-Vis spectroscopy (Figure 1b). The initial UV-Vis
43 spectrum showed a strong peak at 527 nm with a shoulder on the short wavelength side and a
44 very weak signal in the NIR region at 820 nm. As the synthesis progressed, the NIR signal
45 rapidly redshifted within approximately 20 minutes while simultaneously gaining in intensity
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3 (Figure 1c). This peak subsequently blueshifted, initially increasing in intensity and then
4 decreasing again until the peak merged with that of the strong peak in the visible region (see
5 Supporting Information Figure S1). The signal in the green spectral region increased in intensity
6 while only slightly redshifting to 533 nm during the synthesis. This signal is characteristic of
7 small, spherical solid Au NPs with diameters in the range of 20-40 nm,³⁰ while the NIR peak is
8 the source of the dispute mentioned previously and whose origin we wish to clarify with this
9 work.
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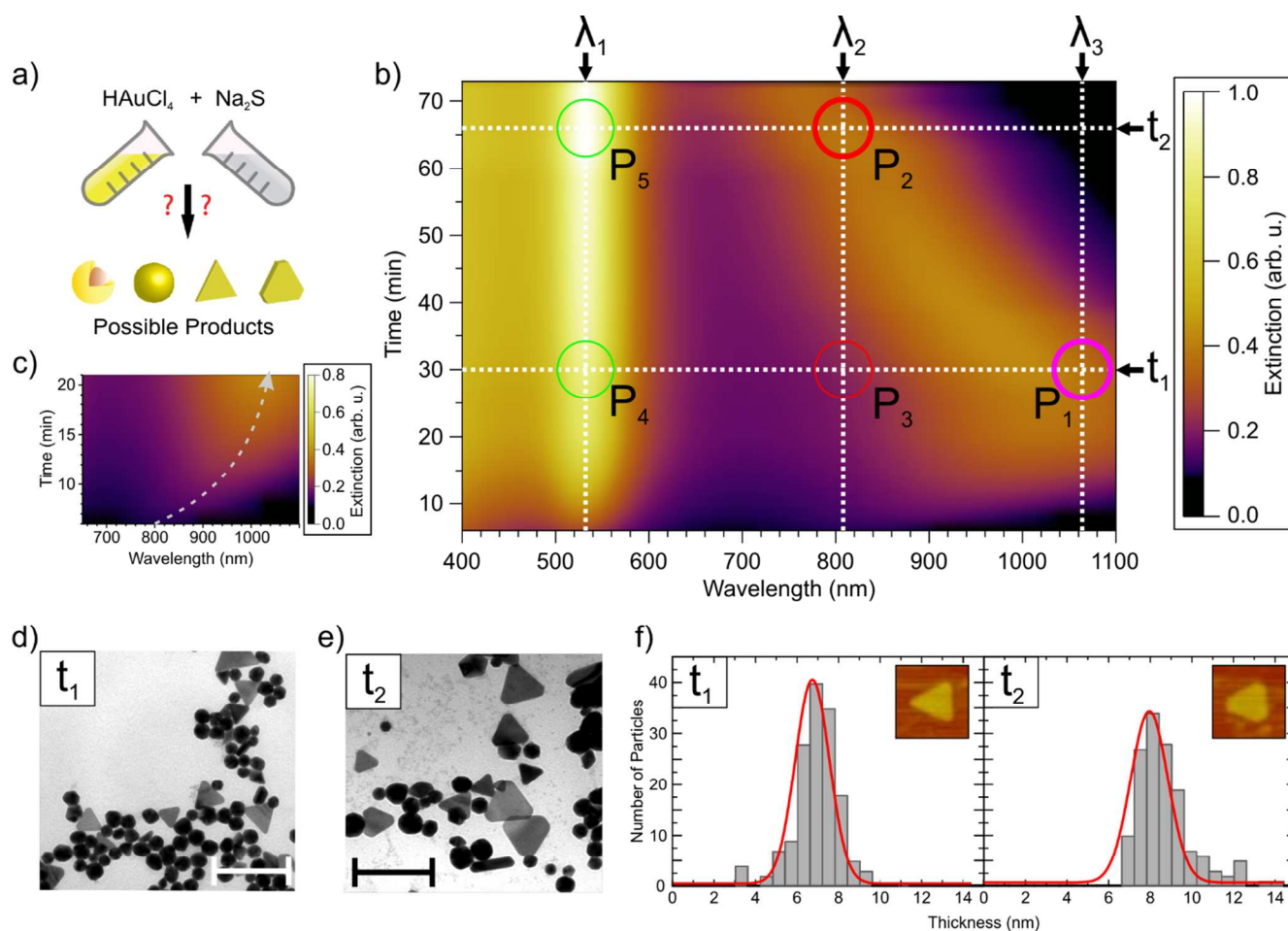


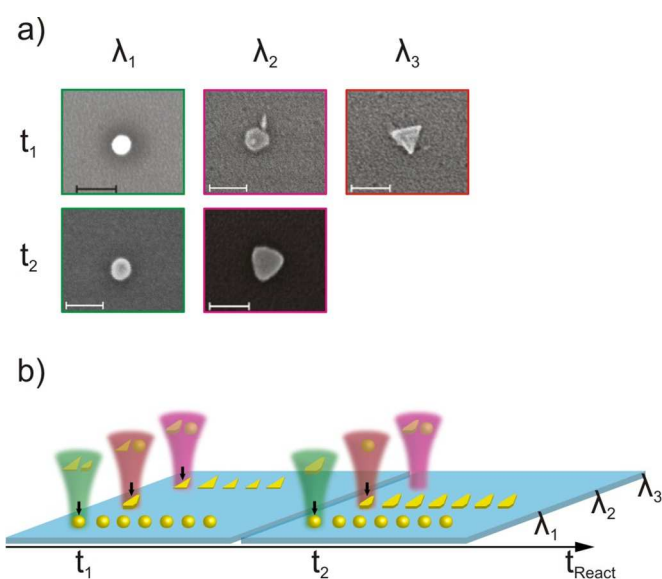
Figure 1. (a) Scheme of the nanoparticle synthesis by reduction of HAuCl_4 by Na_2S . (b) Temporal evolution of the extinction spectrum during the synthesis. (c) Magnification showing the initial redshift during the first 20 minutes of the synthesis. TEM images of nanoparticles when arrested at (d) t_1 and (e) t_2 . The scale bar is 100 nm. (f) AFM height distribution of nanotriangles from different syntheses, when arrested either at t_1 or t_2 . Box size is 150 nm.

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3 In some of the previous studies the initial redshift was purported to stem from the growth of
4 the core of the nanoshells. After 20 minutes the growth of the core terminated and the shell
5 began to grow in thickness. This was said to be the reason for the subsequent observed blueshift
6 of the NIR peak. To investigate these claims we wished to look into the NPs responsible for the
7 NIR resonance at the point of the maximum redshift (time t_1) and then at a later time during the
8 blueshift (time t_2), at an optimal position for biological applications, and at which these NPs have
9 been traditionally analyzed.¹³⁻¹⁹ To this end we repeated the synthesis several times, letting it
10 progress until the NIR peak position reached the desired wavelength. The synthesis was then
11 arrested by adding an additional volume of 1 mM Na_2S solution, also leading to high long-term
12 stability of the suspensions (cf. Figure S2).³¹ In each synthesis there were slight variations in the
13 UV-Vis spectra, such as the maximum redshift or the speed of evolution; however the general
14 trends were always the same. As a first step we looked at the morphology of the ensemble by
15 means of transmission electron microscopy (TEM) on the two aforementioned dispersions. TEM
16 images of both samples showed large areas covered with tiny (~ 5 nm) NPs (see Figure S3),
17 which we thereupon removed *via* selective centrifugation. In the purified samples we found that
18 the overwhelming majority of remaining NPs (75%) were spherical and between 15 and 40 nm
19 in size (Figure 1d, e). The remaining 25 % of NPs turned out to be equilateral triangles. In the
20 sample stopped at t_1 the majority of these NPs showed were sharp-tipped with side lengths
21 ranging from 60 to 100 nm (Figure 1d). The low contrast of the triangles indicates that they were
22 much thinner than the spheres. This was confirmed by atomic force microscopy (AFM)
23 measurements, which found these triangles to be 6.8 ± 0.7 nm thick (cf. Figure 1f and S 4).
24 Additionally, the AFM measurements confirmed that the round NPs were actually spherical in
25 nature. In contrast to the t_1 sample the sample arrested at t_2 was found to contain predominantly
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3 triangles with strongly truncated or rounded corners (Figure 1e). While also having low contrast
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5 these triangles were found to be 8.8 ± 1.0 nm thick, which was slightly thicker than those
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8 observed in the t_1 sample (see Figure 1f and Figure S4). A small amount of nanorods and more
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10 complex looking NPs were also found in both samples. With so many NPs present in each of the
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12 samples it was impossible to distinguish which of these were responsible for which features
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14 observed in the extinction spectra without investigating them individually. Additionally, the
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16 existence of the nanoshells could neither be confirmed nor refuted by these methods.
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20 Therefore we wanted to determine in which way the individual NPs contributed to the UV-Vis
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22 spectra and additionally see if nanoshells were present. Our approach consisted of optically
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24 sorting the NPs depending on their LSPR, printing them individually on substrates and then
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26 comparing their morphology *via* SEM and their optical properties through dark field Rayleigh
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28 scattering. To this end we coupled three lasers into an optical dark field microscope whose
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30 wavelengths $\lambda_1 = 532$ nm, $\lambda_2 = 808$ nm, $\lambda_3 = 1064$ nm overlapped with the previously
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32 mentioned UV-Vis peaks (Figure 1b). A diluted drop of each of the dispersions was placed on a
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34 glass coverslip and imaged with a water immersion objective. The lasers were focused by the
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36 same objective onto the substrates, and printing monitored in the dark field microscope. Once a
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38 nanoparticle was printed, the substrate was shifted laterally by several microns. In this manner
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40 we were able to print rows of nanoparticles, which could then easily be found in the dark-field
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42 and electron microscopes (see Figures S5, 6). When using the 1064 nm laser, the printing could
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44 not be observed as easily in the dark field microscope. So, instead of waiting for a printing event,
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46 we continuously moved the substrate, in order to avoid long time exposure of printed
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48 nanoparticles. The end result however was very similar as for the other two cases.
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3 We wanted to concentrate on two specific points in the UV-Vis spectra, namely the NIR peaks
4 in the t_1 and t_2 dispersions, marked as P_1 and P_2 in Figure 1b. At P_1 NPs were printed rapidly
5 within a matter of seconds using the 1064 nm laser, and all turned out to be sharp-tipped
6 triangles with a large variation in side lengths centered on 70 nm, as confirmed through SEM
7 (Figure 2a).
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39 Figure 2. (a) Nanoparticles printed from
40 samples arrested at t_1 and t_2 using laser with the
41 wavelengths λ_1 , λ_2 and λ_3 respectively. The
42 scalebars are 100 nm. (b) Scheme of the
43 selective printing process.
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3 Likewise, at P₂ single NPs were printed easily employing the 808 nm laser. These turned out to
4 also be triangles albeit with truncated or rounded edges (Figure 2a). Since there was also a weak
5 signal present at P₃ we also investigated the NPs responsible for this resonance. Printing events
6 only occurred after several minutes, and the printed NPs turned out to be truncated or rounded
7 triangles, like at P₂. This indicates that there are hardly any NPs present in the solution with a
8 plasmon resonance resonant with the laser. No round particles were printed at any of these
9 points. However, when investigating P₄ and P₅ with the 532 nm laser, we were able to rapidly
10 print single spherical NPs, approximately 40 nm in diameter. This method proved to be highly
11 specific as each laser printed NPs of only one type of morphology (Figure 2b, Figures S5, S6). In
12 none of the cases were the printed structures shown to be aggregates of NPs.
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27 In order to check whether the laser affects the nanoparticles during the printing process, we
28 calculated the maximum temperatures of the gold particles by means of the finite element
29 method (FEM, Comsol Multiphysics). With the laser powers used, the maximum temperature
30 increase in the truncated triangles is only around 65 °C. For the sharp-tipped nanotriangles the
31 temperature increase amounts to 245 °C (Figure S7). However this is significantly below the
32 melting temperature of bulk gold and also below that of gold nanoparticles.³²⁻³³ Additionally, in
33 heating experiments of gold nanorods, it was seen that while for thermal heating nanorods
34 transformed completely to spheres at 250 °C within an hour, for laser-induced heating no
35 changes were observed for heating temperatures of up to 700 °C. Since our experiment more
36 resembles the laser-induced heating, as the nanoparticles only experience these temperatures for
37 a very short time due to the constant motion of the substrate, reshaping of the NPs should be
38 negligible.³⁴⁻³⁵
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For a full investigation of the individual NPs, their respective morphology needed to be compared to their optical properties. To this end we acquired Rayleigh scattering spectra of optically printed single NPs. All of the spherical NPs showed a strong LSPR around 530 nm, with a width of approximately 400 meV comparable to that of the observed peak in the ensemble

UV-Vis measurements (Figure 3a).

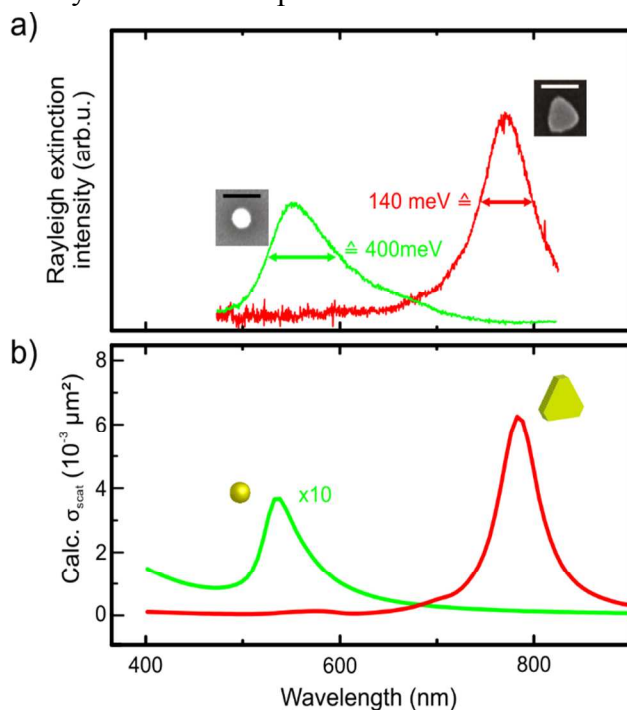


Figure 3. (a) Single particle scattering spectra of printed nanoparticles using the laser wavelengths λ_2 and λ_3 . Note the considerably decreased linewidth of the triangular nanoparticle indicating that the huge linewidth of the ensemble is due to inhomogeneous broadening. Scalebars are 100 nm. (b) Scattering cross sections calculated via FDTD of the printed nanoparticles matching the expected peak positions. Simulation parameters were radius = 20 nm for the sphere and side length = 80 nm, clip = 10 nm and thickness = 9 nm for the clipped triangle.

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18 In contrast to this, the truncated nanotriangles all had resonances located between 750 and 800
19 nm, which were significantly narrower, on the order of only 187 ± 23 meV. This corresponds
20 well to the values reported previously on the purported Au₂S/Au nanoshells¹⁵ and suggests that
21 the broad NIR peak in the UV-Vis spectra was caused by a large inhomogeneous broadening of
22 the nanotriangles in the ensemble. Polarization-dependent measurements showed that the
23 scattering from the nanotriangles was nearly isotropic, a feature generally ascribed to spherical
24 NPs (cf. Figure S8). The narrow linewidth measured should allow for extremely sensitive bio-
25 (chemical) sensing applications, as slight shifts of the refractive index of the medium
26 surrounding the NPs will lead to a large shift of the plasmon resonance (see Figure S9). Based on
27 previous studies and using the values obtained from this analysis we carried out finite difference
28 time domain (FDTD) simulations on the observed NP shapes.³⁶⁻⁴² Using the images obtained
29 from SEM, we are able to reproduce the spectra of the triangles and spheres (Figure 3b).
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46 To corroborate that the nanotriangles alone can reproduce the recorded UV-Vis spectra and
47 explain the red- and subsequent blueshift of the NIR peak during synthesis we carried out a
48 series of simulations and calculations. Statistics of triangle sizes acquired from TEM analysis
49 showed that the side length of the nanotriangles grew during the redshift, reaching a maximum at
50 the time t_1 (Figure 4a). Additionally, the spread in side lengths became considerably broader with
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3 time. Simulated LSPR of sharp-tipped nanotriangles smaller than 50 nm is very weak and
4 located between 650 and 750 nm (Figure 4b). This LSPR redshifts and greatly increases in
5 intensity as the nanotriangles grow, reaching 1000 nm for triangles with 100 nm side length (cf.
6 Figures S10, S11). For truncated triangles, the LSPR blueshifts as the degree of truncation
7 increases, as depicted for triangles of nominal side length 80 nm (Figure 4c). An increasing
8 thickness of both types of nanotriangles, which was measured with an AFM, also leads to a
9 blueshift of the LSPR peak (cf. Figures S12, S13). FDTD simulations of the spheres showed an
10 LSPR around 530 nm, confirming that they are solid Au spheres.^{14-15, 43} Combining these
11 findings together we can fully explain the observed UV-VIS spectra (Figure 4d). To achieve this
12 we estimated the distribution of triangle sizes obtained from the TEM, AFM and SEM
13 measurements. We used this to weight the single particle cross sections, taking into account that
14 the ratio of spheres to triangles was approximately 3:1. We compared the resulting individual
15 spectra with the ensemble extinction spectrum obtained from the UV-Vis spectra at t_1 (Figure
16 4d). The heights of the individual cross sections are normalized with respect to the cross sections
17 of the spheres. The intensity of the individual contributions from the nanotriangles nicely follows
18 the NIR peak, and the intensity ratio of the signal from the spheres to that of the triangles also
19 agrees with the experimental findings, validating the approach we have taken. This shows that
20 the existence of solid Au nanotriangles can fully explain all of the observations made *via* UV-Vis
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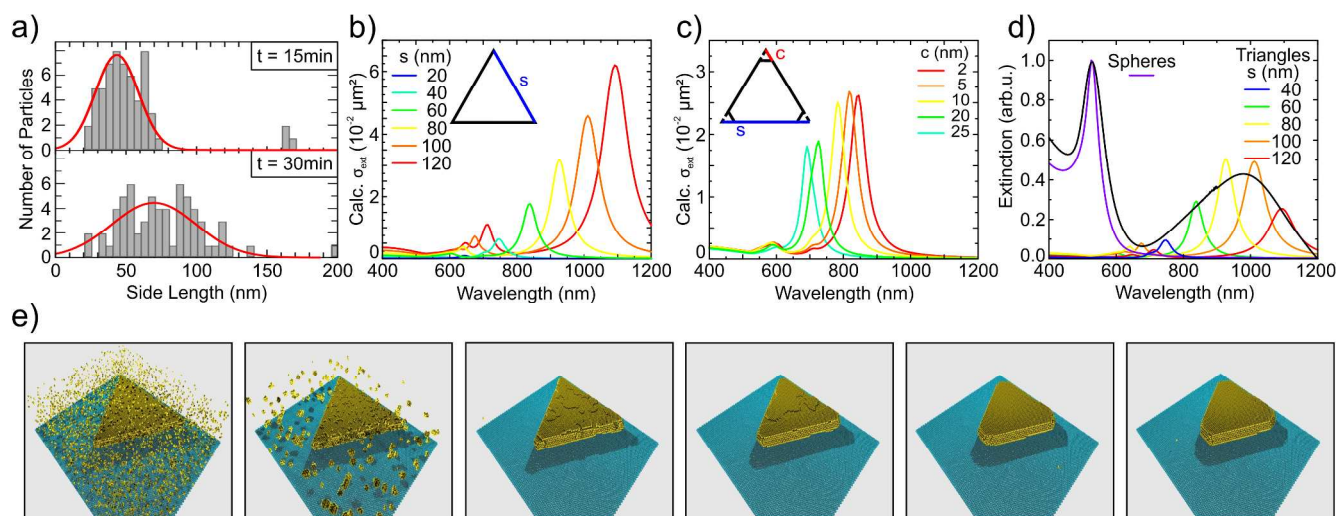


Figure 4. (a) Side lengths of nanotriangles found before (above) and at (below) the maximum red shift.

(b) FDTD calculations of extinction cross sections for sharp triangles of thickness, $d = 7\text{ nm}$, with

varying side length, s . (c) Calculated extinction cross sections of clipped triangles with $d = 9\text{ nm}$, $s =$

80 nm for different clip length, c . (d) Weighted spectrum of simulated nanoparticles (colored) in

comparison to the measured ensemble spectrum (black). The size distribution of sharp triangles was

taken from Fig. 4a, $t = 30\text{min}$ and the intensities from Fig. 4b were weighted correspondingly. The

spectrum of spherical 40nm gold nanoparticles was multiplied by three due to the ratio of spheres to

triangles of roughly 3:1. The measured ensemble spectrum was normalized to the peak of the

simulated spherical nanoparticles. (e) Monte Carlo simulation of the evolution of a sharp thin

nanotriangle to a truncated thicker one for zero, 100, 1000, 2500, 5000 and 10000 Monte Carlo steps.

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15 While the synthesis of different nanotriangles has been long investigated,⁴⁴⁻⁴⁶ studies on the
16 structural changes of the NPs as their synthesis progresses are lacking. It is known that particle
17 shape results from two competitive phenomena: growth far from equilibrium (fast Au atom
18 aggregation) that results in tip formation, and thermodynamics that tends to smooth out these
19 instabilities by slow surface diffusion processes.⁴⁷ We believe that the truncation of originally
20 sharp-tipped nanotriangles is the result of both intraparticle Au atom diffusion and the Au
21 incorporation from non-reduced Au species (from excess reactant) and the small NPs present in
22 the solution through an Ostwald ripening process. Thus, depending on the reaction time, one of
23 these two phenomena prevails, explaining the initial redshift (kinetics) and the final blueshift
24 (thermodynamics) experimentally observed.
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38 We performed Monte Carlo (MC) simulations to mimic the structural transformations that
39 sharp nanotriangles may experience in the presence of an Au atom source (for details see:
40 Supporting Information). For the initial condition we started with a triangular Au NP with sharp
41 edges surrounded by a random dispersion of Au atoms occupying 1% of the available sites
42 (Figure 4e). Once the simulation commenced these Au atoms rapidly agglomerated, forming
43 small clusters. These then condensed on the nanotriangles through an Ostwald ripening process.
44 Mass transport by surface diffusion subsequently led to Au atoms wandering from the tips and
45 sides to the large top surface, increasing the thickness and rounding the sharp tips. Considering
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3 all of the obtained results we have devised a model to explain the observed dynamics of the UV-
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5 Vis spectra during the synthesis. Initially the synthesis produces two main types of NPs, solid Au
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7 spheres and thin triangles. As the synthesis progresses the spheres become slightly larger in size,
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9 while the triangles became significantly larger. Once the latter reach side lengths of ≈ 50 nm,
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11 their extinction cross sections become large enough to be discernible from the spheres in the UV-
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13 Vis spectra. As the synthesis progresses triangles become larger, but at different rates, so the
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15 spread in sizes becomes greater, leading in turn to a larger inhomogeneous broadening. The
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17 triangles then become truncated and thicker, producing the observed blue-shift and explaining
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19 the slight decrease in inhomogeneous broadening, as the LSPR positions of thicker nanotriangles
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21 are not quite as dependent on side length as the thinner ones are. This model also agrees with
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23 experiments showing the formation and growth of silver nanotriangles.⁴⁸
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31 CONCLUSIONS

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34 In summary we have used optical printing and single particle spectroscopy to investigate the
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36 synthesis of Au NPs by the reduction of HAuCl_4 with Na_2S . UV-Vis spectra showed a strong
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38 peak in the NIR that first redshifted in time and subsequently blueshifted. By stopping the
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40 synthesis at different times and employing lasers tuned to the UV-Vis peaks of these dispersions,
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42 we were able to determine which NPs were responsible for the observed resonances. Single NPs
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44 whose resonance overlapped with the laser wavelength were printed onto substrates. Rayleigh
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46 scattering spectra of these single NPs were acquired and compared with SEM images. We found
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48 that, contrary to previous reports, the NPs responsible for the NIR peak are neither nanoshells
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50 comprising an Au shell around an Au_2S core nor Au NP aggregates, but actually thin, solid Au
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52 nanotriangles. By arresting the reaction at different times we studied the structural evolution of
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3 the nanotriangles and the corresponding effect on their optical properties. The initial redshift is
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5 caused by the lateral growth of these triangles. The subsequent blueshift is a result of the
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7 simultaneous increasing thickness and degree of rounding or truncation of the nanotriangle tips
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9 as was confirmed by MC simulations. FDTD calculations confirmed measurements showing that
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11 the linewidth of the nanotriangles is extremely narrow, less than 190 meV, and that small
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13 variations in size, thickness and truncation lead to the large inhomogeneous linewidth of the
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15 ensembles. These optical properties along with the long-term stability of the nanotriangles show
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17 the great potential of the Au nanotriangles in various applications, such as *in-vivo* imaging and
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19 bio-(chemical) sensing. Optical forces have proven an important tool not only for surface
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21 patterning of substrates but also as a means to investigate chemical reactions and optical
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23 properties of NPs. Additionally, the use of multiple laser wavelengths tuned to individual
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25 plasmon resonances will permit surface structuring simultaneously with a multitude of different
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27 pre-formed plasmonic NPs, leading to further possible applications in chemical and bio-sensing,
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29 fluorescence enhancement and photodetection.
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39 METHODS

40 **Gold Nanoparticle Synthesis**

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42 Gold nanoparticles were synthesized at room temperature in a two-step reaction by reduction
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44 of H_{Au}Cl₄ (G4022 Sigma Aldrich) with Na₂S (407410 Sigma Aldrich), as reported elsewhere.¹⁴
45
46 In a typical procedure, 10 mL of a 2 mM H_{Au}Cl₄ solution were mixed with 12 mL of fresh 1mM
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48 Na₂S solution (pH = 6). After two minutes 2 mL 1 mM Na₂S solution were added, giving a total
49
50 molar S/Au ratio of 0.65. The reaction was then allowed to evolve for different times, between 5
51
52 minutes and 3 hours, and then 14 mL of 1 mM Na₂S solution was added to arrest its evolution.
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Nanoparticle printing

For the particle printing experiments, cw-lasers of wavelengths 532nm (Coherent Verdi V10, Coherent, USA), 1064nm (Cobolt Rumba 05-10, Cobolt, Sweden) and a Ti:Sapphire Oscillator (Spectra-Physics Tsunami, Spectra-Physics, USA) pumped by a 532nm Spectra-Physics Millennia V (Spectra-Physics, USA) operated in cw mode at 808nm were sequentially coupled into an upright Zeiss AxioTech 100 (Zeiss, Germany) microscope. The laser beams were focused by a Zeiss Achroplan 100x/1.0W water dipping objective (Zeiss, Germany) which was dipped directly into the diluted nanoparticle dispersion of interest. In order to correct for the chromatic aberration of the objective appearing in the NIR, the 1064nm laser was prefocused by a lens with $f=1000\text{mm}$ placing the laser focus in the image plane. Nanoparticles were made visible by applying dark field conditions using a dark field condenser placed in the light path.⁴⁹

For printing at 808 and 1064nm a laser power around 15mW on the sample was used. The power was set so that the average time between two printing events was on the scale of 5s. At 532nm a lower laser power of around 5mW on the sample was enough due to the significantly higher concentration of resonant spherical particles. The particle dispersion was diluted 1:100 with MiliQ water before printing. More water was added during the experiment to compensate for evaporation keeping the particle concentration constant. The laser spot sizes were around $1\mu\text{m}$. After each printing event, the piezo stage was moved by several microns to avoid printing of multiple particles in close vicinity.

The area where the nanoparticles were printed was marked by a scratch to allow their later relocation for single particle spectroscopy and SEM. After the optical printing was performed samples were rinsed carefully with MiliQ water to remove the nanoparticle dispersion avoiding further random deposition.

Single particle spectroscopy

Single particle spectra were acquired using an upright Zeiss AxioLab.A1 microscope in dark field condition attached to a Princeton Instruments SP2300 monochromator (Princeton Instruments, USA) together with an LN-cooled CCD (Acton Spec-10, Princeton Instruments, USA).

Scanning electron microscopy (SEM)

The individual printed nanoparticles were imaged using a Zeiss Ultra 55 Plus SEM (Zeiss, Germany). In order to avoid electrical charging of the substrate a thin (~1nm) layer of a gold-palladium alloy was sputtered on top of the substrate with a Leica Sputter Coater EM SCD005 (Leica Microsystems, Germany) making the substrate's surface conductive.

Transmission Electron Microscopy (TEM)

TEM samples were prepared by depositing 7 μL of arrested nanoparticle dispersion onto a carbon coated copper TEM grid (Ted Pella, USA) which was left to dry in the dark and subsequently measured with a JEOL JEM-1011 TEM (JEOL, Japan)

UV-Vis Spectroscopy

Extinction spectra of the nanoparticle dispersion at different stages of the synthesis were taken with an Agilent Cary 60 Spectrophotometer.

Atomic Force Microscopy (AFM)

AFM experiments were performed in air in tapping mode with a multimode microscope operated by a Nanoscope V control unit, both from Veeco Instruments (Bruker Corporation, USA). Silicon tips from NanoWorld (Switzerland) were used in all measurements at $f_0 = 285$ kHz; $k = 42$ N/m. Samples were prepared by drop-casting the nanoparticle dispersions on freshly cleaved HOPG substrates (Ted Pella Inc., USA).

Simulations of effective extinction and scattering cross sections:

Effective cross sections were calculated using the software Lumerical FDTD Solutions (Lumerical, Vancouver, Canada). The simulation environment measured 2000 x 2000 x 2000 nm³ applying 24 perfectly matched layers with a reflection coefficient of 0.0001. The gold nanoparticles were modeled by interpolating the dielectric functions of gold and water from the measured data of Johnson and Christy⁵⁰ and Palik⁵¹ respectively and placed in the middle of simulation volume filled up with water. The particles were surrounded by a mesh with a cell size of 0.05nm. The linearly polarized electromagnetic plane-wave was injected in positive z-direction, resembling the dark field illumination used in the experiments. The simulation time was chosen to be 1000 fs after which convergence was attained.

Simulation of particle temperatures

The temperatures the nanoparticles have during the printing process were calculated by using finite elements method (Comsol Multiphysics 5.2, COMSOL, Sweden). Therefore we determined the laser power density by first imaging the laser spot and then fitting a 2D Gaussian function. We then integrated this Gaussian over the previously calculated absorption cross-section of the particular particles at the printing wavelength. Knowing the total laser power this gives us the absorbed laser power density which is converted to heat.

In COMSOL we modeled the dimensions of the particular sharp and clipped triangles in a water surrounding. Giving the structure the calculated absorbed laser power density leads to the calculated equilibrium temperature.

Monte Carlo Simulations

We have implemented a kinetic Monte Carlo model in which particles are able to perform diffusion events between nearest neighbors in an FCC lattice. Transition probabilities

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3 $P(c_{ini} \rightarrow c_{fin})$ from a given initial configuration c_{ini} into c_{fin} were taken according to the transition
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5 state theory, *i.e.*

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$$P(c_{ini} \rightarrow c_{fin}) = \nu \exp[-E_{act}(c_{ini} \rightarrow c_{fin})/k_B T]$$

9
10 where ν is the attempt frequency set equal to 5×10^{12} Hz, $E_{act}(c_{ini} \rightarrow c_{fin})$ is the activation
11 energy for the transition, k_B is the Boltzmann constant, and T is the absolute temperature.

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13 We have chosen a simple model⁵²⁻⁵³ for the activation energies, in which a linear dependence
14 of the energy barrier with the first-neighbor coordination number (n) of the diffusing particle at
15 the initial state is assumed: $E_{act} = E_b \cdot n$, where E_b can be interpreted as the bond energy between
16 two nearest neighbor particles (in this work, we have set $E_b = 0.1$ eV and $T = 300$ K).

24 ASSOCIATED CONTENT

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28 **Supporting Information.** The following material is included: Extinction spectra showing:
29 synthesis after 24 hours, long-term stability of samples after synthesis was arrested. Additional
30 TEM, AFM, SEM images depicting nanoparticles found in the samples. Comparison of
31 experimental and calculated polarization-dependent scattering spectra on nanotriangles.
32 Calculated extinction spectra showing the dependence of the spectra on surrounding medium,
33 nanotriangle thickness, degree of truncation, and size. This material is available free of charge
34 via the Internet at <http://pubs.acs.org>.
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Author Contributions

The manuscript was written through contributions of all authors. All authors have given approval to the final version of the manuscript. ‡These authors contributed equally.

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ABBREVIATIONS

LSP, localized surface plasmon; LSPR, localized surface plasmon resonance; NIR, near infrared; SEM, scanning electron microscopy; TEM, transmission electron microscopy; AFM, atomic force microscopy; FDTD, finite difference time domain; MC, Monte Carlo; NP, nanoparticle; Au, gold.

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