# Thickness determination of electrochemical titanium oxide ( $\mathrm{Ti} / \mathrm{TiO}_{2}$ ) formed in $\mathrm{HClO}_{4}$ solutions 

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#### Abstract

Anodic Titanium oxide films were potentiodinamically grown on Ti foil and glass/Ti in $0.010 \mathrm{M} \mathrm{HClO}_{4}$ at $50 \mathrm{mVs}^{-1}$. The current density-potential curves $(j-E)$ showed that the oxide grows according to the physical model for high-field conduction. However, for final potential $\left(E_{f}\right)$ higher than $1.5 \mathrm{~V} \mathrm{vs} . \mathrm{Ag} / \mathrm{AgCl}$ (sat. KCl ) the oxygen evolution reaction becomes more significant, and the formation of bubbles prevented or made more difficult the application of in situ techniques for the simultaneous study of the thickness and optical properties of the anodic layer. We developed a method based on electron probe microanalysis (EPMA) to calculate oxide thickness using ex situ ellipsometry as a reference technique. The normalized intensity of the $\mathrm{O} \mathrm{K} \alpha$ peaks was measured for anodic oxides corresponding to $E_{f}$ values from spontaneous oxide up to 50 V , where a linear relationship was observed for a narrower range of final potentials. This behaviour was studied with Monte Carlo simulations. After calibration with ellipsometric results, to take into account sample damage during the electron irradiation, EPMA was applied as a method for thickness determination at $E_{f} \geq 1.0 \mathrm{~V}$. Once the method was established, ex situ thickness determinations became independent of the preparation method of the oxide layer, which represents a comparative advantage against ex situ ellipsometry.


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## 1. Introduction

In situ techniques provide a wide range of advantages regarding the study of the formation of thin oxide layers from a metal substrate [1] compared to ex situ procedures. However, when the reactions involved entail evolution of bubbles or dissolutionprecipitation mechanisms of layer formation, the experimental requirements for determining thickness in situ are more demanding. In this situation, ex situ techniques are a convenient tool, making a wider range of potentials accessible. In the case of EPMA, once the technique is calibrated, the intensity of the $\mathrm{O} \mathrm{K} \alpha$ characteristic line can be used to determine the thickness of electrochemical prepared oxides under different experimental conditions.

[^0]Electrochemical titanium oxide is formed mainly by galvanostatic, potentiostatic and potentiodynamic methods [2-14]. Poznyak et al. [15] investigated the morphology, structure and semiconducting properties of titanium oxide films formed by high-voltage pulsed discharge and compared them with the conventional galvanostatic anodization method, finding remarkably different physicochemical properties. Mazzarolo et al. [16] recently studied the anodic growth of titanium under the same conditions used for nanotube growth but in absence of fluoride ions. They determined the individual processes that occur during anodization up to 25 V in a two-electrode configuration using impedance measurements to calculate oxide thickness.

Anodic oxide thickness, mainly of valve metals, is usually calculated by charge integration from CVs obtained at different $E_{f}$ values [17-20]. Pouilleau et al. [12] studied the structure and composition of $\mathrm{TiO}_{2}$ films galvanostatically grown in $0.5 \mathrm{M} \mathrm{H}_{2} \mathrm{SO}_{4}$. They found thicknesses of between 23 and 62 nm for $E_{f}$ of $5 \mathrm{~V}, 10 \mathrm{~V}$ and 20 V ,
being highly sensitive to the environment. They attributed the variations of interfacial capacitance to variations in the thickness of the outer $\mathrm{TiO}_{2}$ layer, rather than to the modification of the composition of the surface compound. Their conclusions were based on surface analysis results, AFM and XPS. In previous works we studied anodic oxides formed in 0.1 M NaCl and compared them with each other and with respect to differences between thermal and sol-gel film growth [21-24]. It was found that the thickness of the oxide layer was the main factor determining the electrode response associated to the photocurrent quantum yields at $\lambda<390 \mathrm{~nm}$. In the present work, anodic titanium oxide films were potentiodynamically formed on two different titanium substrates in $0.010 \mathrm{M} \mathrm{HClO}_{4}$. The optical properties of the films prepared in a wide range of $E_{f}$ values (up to 60 V ) were obtained from ex situ ellipsometric measurements. The thickness of anodic oxide films for $E_{f}$ values up to 10 V , was calculated assuming that oxide film grows as a single anisotropic layer with constant optical properties. These results were compared to the thickness values obtained by charge integration from current density-potential curves $(j-E)$ (up to 1.5 V ) and EPMA (up to 10 V ). We have developed a comprehensive study of film thickness determination by EPMA, employing ex-situ ellipsometry as a reference technique for the calculation of anodic oxide thickness values. We have used Monte Carlo simulations to analyse the thickness range where proper results can be obtained considering the thin film hypothesis (TFH). This validated the use of the characteristic line of $\mathrm{O} \mathrm{K} \alpha$, detected under the irradiation of an electron probe, as a measure of oxide thickness.

## 2. Experimental

### 2.1. Materials

The working electrode was either a metallic titanium foil (purity $99.8 \%$ ) or a glass/Ti, where the titanium metal thickness was $2 \mu \mathrm{~m}$ (glass sputtering targets prepared by Mateck GmbH). The glass/Ti samples were manually cleaned with a sequence of acetone-ethanol-water washing. The foil samples were sonicated for 10 min in MilliQ water and a thermal oxide was formed at $450^{\circ} \mathrm{C}$ to isolate the back of the foil. The side of the electrode that has to be in contact with the electrolyte was polished with emery paper of different grades (1200 and 2500) and finished with $1 \mu \mathrm{~m}$ alumina. After mechanical polishing, two set of titanium foil samples were prepared: a) The samples labeled foil ${ }^{(1)}$ were maintained under nitrogen gas ( $99.999 \%$ purity) before and after the electrochemical oxide was formed. b) The samples labeled foil ${ }^{(2)}$ were stored one week before the electrochemical oxide was formed. The oxide film was grown potentiodynamically up to different $E_{f}(1,2,3 \mathrm{~V}$ for both substrates and $4,8,9,10,30,40,50$ and 60 V only for foils samples) at $50 \mathrm{mVs}^{-1}$. After oxide growth, the electrode was kept at the same $E_{f}$ value for 15 min in order to stabilize the oxide film. A $\mathrm{Ti} / \mathrm{TiO}_{2}$ additional sample prepared from a titanium disk and electrochemically grown up to 30 V in $1 \%$ boric acid and $5 \%$ phosphoric acid was only studied by EPMA. This sample was used to check the reliability of this technique to determine the oxide layer thickness. Potential cycling between -1000 and 1500 mV (and also from - 1000 to 0 V ) at $50 \mathrm{mVs}^{-1}$, starting from the open circuit potential (ocp) of the electrode towards more negative values was performed. To ensure minimal differences between the different samples, a stabilized voltammogram was recorded for each electrode, and the $j-E$ profile was compared with a reference voltammogram. A difference of less than $5 \%$ was accepted. Otherwise, pretreatment of the substrate and all the described procedures were repeated. The electrolyte used in all measurements was a $0.010 \mathrm{M} \mathrm{HClO}_{4}$ aqueous solution. The reference electrode was $\mathrm{Ag} / \mathrm{AgCl}$ (saturated KCl ). It was connected to the cell through a Luggin capillary filled with
the same solution as in the main cell compartment. This arrangement minimized possible contamination of the solution by chloride anions. Oxygen free solutions were obtained by continuously purging the electrochemical cell with nitrogen gas (99.999\% purity). All solutions, as well as glassware cleaning were prepared with water from a Milli-Q Millipore System. All chemicals were of analytical grade from Merck.

### 2.2. Cyclic Voltammetry (CV)

The electrochemical analyzer was an AUTOLAB-PGSTAT302N equipped with staircase scan module. In blank experiments, voltammetric curves were recorded at a scan rate $(v)$ of $50 \mathrm{mVs}^{-1}$ starting at the ocp and changing toward more negative values. The effect of double-layer charging was assessed by subtracting a given baseline from each steady state current voltammogram, $j_{s s}$. The results corresponding to the first voltammetric cycle are shown, otherwise it is indicated in the figure caption.

### 2.3. Ellipsometry

Optical measurements were performed ex-situ (in air) using a Rudolf Research rotating-analyzer type automatic ellipsometer (vertical type, 2000FT model), equipped with a 75 W tungsten lamp as the light source and a filter ( 546.1 nm ). The incidence angle used for all experiments was $70.00^{\circ}$. The light reflected by the surface was passed through the rotating analyzer and into a photomultiplier tube (PMT) for detection. The sinusoidal output of the PMT was digitized by a HP microcomputer and the values of the ellipsometric angles $\Psi$ and $\Delta$ were calculated by the computer from the three Fourier coefficients of the PMT signal. The PMT signal was averaged five times and the data were collected in operatortriggered mode. Ten measurements were taken for each sample. The optical data were fitted by using an anisotropic single-layer model, which assumes that oxide film growth takes place as a single anisotropic layer of constant optical properties growing on a flat bare titanium substrate. The fitting model uses the Simplex method [25] to minimize the error function
$\left.G=\frac{1}{m} \sum_{i}\left[\left(\Delta_{i}^{\exp }-\Delta_{i}^{\text {calc }}\right)^{2}+\left(\Psi_{i}^{\exp }-\Psi_{i}^{\text {calc }}\right)^{2}\right)\right]$
where the superscript "calc" denotes the values calculated from the model. No assumptions were made about the values of the optical properties of the oxide when fitting the experimental response, allowing the Simplex algorithm to vary the refractive index ( n ) and the extinction coefficient $(\mathrm{k})$ in the normal (or ordinary; $\boldsymbol{n}_{\text {ord }}=n_{\text {ord }}-k_{\text {ord }} i$ ) and parallel (or extraordinary; $\boldsymbol{n}_{\text {extr }}=n_{\text {extr }}-k_{\text {extr }} i$ ) directions with respect to the surface of the anisotropic layer. Film thickness is an additional parameter that is also optimized in the fitting procedure; once the optical response is fitted in terms of the given model, both the optical properties and the thicknesses of the oxide films are obtained. The $\mathrm{TiO}_{2}$ films were prepared by anodization of Ti foils, as described in section 2.1. Growth at high voltage took place using a two-electrode setup with a 173 EG\&G PAR potentiostat/galvanostat (Princeton Applied Research, USA) coupled to a 175 EG\&G PAR signal generator and a 7090A Hewlett Packard plotting recorder.

### 2.4. Scanning Electron Microscopy (SEM)

The surfaces of the titanium foils samples were studied with a field emission scanning electron microscope Carl Zeiss Ligma before and after anodic oxide formation, at the Laboratorio de Microscopía y Análisis por Rayos X (LAMARX) of the National University of Córdoba. Two different secondary electron detectors
were used: the in-lens type for high magnification and the Everhard Thornley detector for low magnification.

### 2.5. Electron Probe MicroAnalysis (EPMA)

The JEOL JXA 8230 microprobe available at LAMARX was used for three different purposes detailed in the following subsections.

### 2.5.1. Test of purity

To check the purity of the Ti film in the glass/Ti sample (without electrochemical treatment) both 12 and 20 keV electron beams were used to excite the sample with an energy dispersive spectrometer for detection. According to the results obtained with the standardless quantification algorithm implemented in POEMA [26], the only impurity detected was potassium with a mass concentration of $0.19 \pm 0.03 \%$.

### 2.5.2. Determination of film thickness

When a flat sample is irradiated with an electron beam, the incident particles undergo multiple interactions with the atoms and are finally stopped within the material if it is thick enough (i.e., typically around one micrometer for electrons with incidence energy $E_{0}$ below 10 keV ). Among the different interactions undergone by the electrons, an inner shell ionization can be followed by a radiative decay. The number of characteristic X-rays produced after a particular decay is greater for the more abundant elements in the sample. In the configuration studied here, a nanometric film of $\mathrm{TiO}_{2}$ lies on a Ti substrate, which can be considered semi-infinite for the 4 keV electron beam used. With these ideas in mind, it can be stated that the thicker the oxide film, the more intense the $\mathrm{O} \mathrm{K} \alpha$ peak detected. Moreover, it is reasonable to assume that peak intensity $I$ will be proportional to the film mass thickness $\rho d$ for thin layers of mass density $\rho$ and linear thickness $d$, so that self-absorption and electron scattering can be disregarded [27]. The relationship between $I$ and film thickness was previously investigated by means of Monte Carlo simulations, as explained in section 3.4.1, in order to find a convenient value for $E_{0}$ and the thickness region of linearity between I and $\rho d$. The software package PENELOPE was used [28] for this end. The samples formed by electrochemical oxide growth on metallic titanium foils at all $E_{f}$ values, as described in subsection 2.1 , were irradiated with a 4 keV electron beam. In view of the strong overlapping between the $\mathrm{O} \mathrm{K} \alpha$ peak and the Ti-L lines, the measurement of spectra was carried out using the LDE1 crystal (with mylar window) of a wavelength dispersive spectrometer, whose high spectral resolution allows for a clear separation of the involved lines. A beam current $I_{0}$ of 80 nA was used, while the step size between channels was $10 \mu \mathrm{~m}$, with a dwell time of $1.3 \mathrm{sec} /$ channel, scanning a spectral range between 375 and 583 eV .

### 2.5.3. Study of the specimen current

The specimen current $I_{\text {spec }}$ is defined as the electric charge conducted from the sample to the ground per unit of time. Taking into account that electrons from the incident beam can be backscattered from the sample, and considering also the emission of secondary electrons from the sample surface, the following can be written:
$I_{0}=I_{\text {spec }}+I_{\eta}+I_{\delta}$
where $I_{\eta}$ and $I_{\delta}$ are the electric currents due to electron backscattering and secondary electron emission, respectively. Eq. (2) can be normalized by $I_{0}$, so that,
$1=\frac{I_{\text {spec }}}{I_{0}}+\eta+\delta$
where the so called backscattering and secondary electron coefficients $\eta$ and $\delta$, respectively, represent the number of electrons of each particular kind divided by the number of incident electrons.

These coefficients depend on the mean atomic number of the sample and $E_{0}$. Particularly, $\eta$ depends more markedly on the mean atomic number, while $\delta$ depends on $E_{0}$ and on surface characteristics. For these reasons, at constant values of $E_{0}$ and $I_{0}$, any change of the sample during the experiment would be evidenced by variations in $I_{\text {spec }}$. To study the stability of the sample during electron irradiation, the $I_{\text {spec }}$ values were recorded throghout the experiment for two materials: bulk rutile $\left(\mathrm{TiO}_{2}\right)$, and a $\mathrm{TiO}_{2}$ film anodically grown on a Ti substrate at $E_{f}=30 \mathrm{~V}$. It must be clarified that this film was grown by a different method than the one described in subsection 2.1, but this fact does not affect the conclusions obtained as shown below.

## 3. Results and Discussion

### 3.1. Cyclic Voltammetry measurements

Fig. 1 shows $j-E$ profiles obtained for different $\mathrm{Ti} / \mathrm{TiO}_{2}$ foil ${ }^{(2)}$ substrates, anodically oxidized in $0.010 \mathrm{M} \mathrm{HClO}_{4}$ under potentiodynamic conditions with different $E_{f}$ values. The physical model of the high-field conduction behavior of oxide growth can be used up to 1.5 V , as can be inferred from the steady state current observed after the onset in the $j-E$ curves, which is the characteristic kinetic behaviour of valve metals [29]. At more positive potentials than 1.5 V , the contribution from the oxygen evolution reaction becomes more significant, as it can be seen from the exponentiallike current increase observed when the potential turns to higher positive values. Table 1 shows the potential values of the different electrodes for the onset of the mentioned processes. A large difference is observed when comparing the $E_{\text {onset }}^{1}$ values of the samples laveled foil ${ }^{(1)}$ and foil ${ }^{(2)}$, indicating that the pre-treatment for foil ${ }^{(2)}$ induced spontaneous oxide growth to a final thickness similar to the glass/titanium sample. Even though initial thickness for the two foil samples are different, the $\Delta E$ is 0.62 V for both electrodes and the potential window is remarkably large when compared to the glas/Ti electrodes. To keep track of the general characteristics of the initial thickness, two samples with $E_{f}$ of 1.0 and 1.5 V were stored six months, and then the electrochemical oxide was grown up to $E_{f}=1.5 \mathrm{~V}$. The current started increasing from 1.5 V in $j-E$ profiles, indicating that a spontaneous oxide was formed with a thickness corresponding at least to 1.5 V for both samples. Marino et al. [30] investigated the potentiodynamic formation of $\mathrm{TiO}_{2}$ up to $E_{f}=5 \mathrm{~V}$ in $\mathrm{NaH}_{2} \mathrm{PO}_{4} / \mathrm{Na}_{2} \mathrm{HPO}_{4}$. They analysed the dissolutionreconstruction of different oxide films and calculated anodic oxide thickness by integration of the $j-E$ profiles. They found an anodization rate of $2.5 \mathrm{nmV}^{-1}$ and estimated the initial thickness of the film (from air oxidation of the polished electrode) to be in the range $\sim$ $1.9-3.0 \mathrm{~nm}$. They identified the oxides present in the films by X-ray diffraction and found that for the film grown up to $E_{f}=1.5 \mathrm{~V}$ the diffractogram peaks corresponded to orthorhombic $\mathrm{TiO}_{2}$, while for the films grown up to $E_{f}>3 \mathrm{~V}$ there were also peaks related to the

## Table 1

Open Circuit Potential ( $E_{\text {ocp }}$ ) from foil and glass $/ \mathrm{Ti} / \mathrm{TiO}_{2} / 0.010 \mathrm{M}_{\mathrm{HClO}}^{4}$ interfaces $E_{\text {onset }}^{1}$ and $E_{\text {onset }}^{2}$ refers to the start of oxide formation and oxygen evolution reactions, respectively. $\Delta E=\left|E_{\text {onset }}^{1}\right|-\left|E_{\text {onset }}^{2}\right|$. Electrode potential sweep rate: $50 \mathrm{mVs}{ }^{-1}$. Reference electrode $\mathrm{Ag} / \mathrm{AgCl}$ (sat. KCl ). The thermal oxide was formed at $400^{\circ} \mathrm{C}$ during 1 h .

| Electrode | $E_{\text {ocp }}(\mathrm{V})$ | $E_{\text {onset }}^{1}(\mathrm{~V})$ | $E_{\text {onset }}^{2}(\mathrm{~V})$ | $\Delta E$ |
| :--- | :--- | :--- | :--- | :--- |
| glass $/ T i / T i O_{2}$ | 0.29 | 0.89 | 1.0 | 0.11 |
| foil $^{(1)}$ | 0.020 | 0.30 | 0.95 | 0.62 |
| foil $^{(2)}$ | 0.30 | 0.85 | 1.5 | 0.62 |
| glass $/ \mathrm{Ti} / \mathrm{TiO}_{2}$ (thermally grown) | -0.080 | - | - | - |

foil ${ }^{(1)}$ (electrochemical oxide inmediately formed after pre-treatment or kept in oxygen-free atmosphere)
foil ${ }^{(2)}$ (stored one week before electrochemical oxide formation)


Fig. 1. Cyclic voltammogram curves of $\mathrm{Ti} / \mathrm{TiO}_{2}$ electrodes (each curve corresponds to a different foil ${ }^{(2)}$ ) in $0.010 \mathrm{M} \mathrm{HClO}_{4}$ at $50 \mathrm{mVs}^{-1}$. The anodic oxide was potentiodinamically formed by increasing the positive potential limit $\left(E_{f} / \mathrm{V}\right.$ vs $\mathrm{Ag} / \mathrm{AgCl}($ sat KCl$)$ ). Inset: $j-E$ profiles of foil ${ }^{(1)}$ (red line) and foil ${ }^{(2)}$ (black line) samples in the -1.0 V to 0.6 V range.
monoclinic mixed oxide $\mathrm{Ti}_{3} \mathrm{O}_{5}$. From the integration of current vs. time profiles (from the $E_{\text {onset }}^{1}$ to $E_{\text {onset }}^{2}$ potential region) the charge due to anodic oxide growth was calculated. Assuming a fixed stoichiometry of $\mathrm{TiO}_{2}$, and using a value of $3.44 \mathrm{gcm}^{-3}$ for the titanium oxide density (see section 3.3 below), the thickness values where obtained using the following equation
$\Delta d_{o x}=\frac{M_{\text {oxide }} Q_{\text {oxide }}}{n F \delta}$
where $M_{\text {oxide }}$ is oxide molar mass, $Q_{\text {oxide }}$ oxide charge, $\delta$ is oxide density, $n$ the number of electrons and $F$ the Faraday constant. The procedure used to obtain the charge was as follows: a) the double-layer current density was substracted using the value at $E_{\text {onset }}^{1} \mathrm{~b}$ ) current density (current divided by the geometric area of the electrode) was integrated as a function of time for different $E_{f}$. c) thickness was calculated assuming perfect stoichiometry, four electrons, and oxide density of $3.44 \mathrm{gcm}^{-3}$ (see section 3.3 ).

Jaeggi et al. [20] studied the potentiodynamic growth at $5 \mathrm{~V} / \mathrm{s}$, of titanium oxidefilms in 1 M sulfuric and phosphoric acid. From the $j-E$ profiles they observed that $O_{2}$ evolution did not occur in the phosphate-containing electrolyte at least up to 60 V . In our case, the
$\mathrm{O}_{2}$ evolution reaction started at the lowest potentials ever recorded when compared to any conditions reported in the literature. Surface techniques used to characterize anodic oxide films have demonstrated that there is a direct correlation between the defects in the $\mathrm{TiO}_{2}$ crystalline lattice and the features observed in the CVs [12]. Galvanostatically formed anodic oxide in $0.5 \mathrm{M} \mathrm{H}_{2} \mathrm{SO}_{4}$ up to 5 V has a thickness of $23 \pm 2 \mathrm{~nm}, 35 \pm 4 \mathrm{~nm}$ at $E_{f}=10 \mathrm{~V}$ and $62 \pm 5 \mathrm{~nm}$ at $E_{f}=20 \mathrm{~V}$. Characteristic peaks of anatase and rutile were reported on Raman spectra of $\mathrm{TiO}_{2}$ films, formed at low voltages $(0-5 \mathrm{~V})$ [31,8]. They proposed a three layer model by ellipsometry in which an intermediate $\mathrm{TiO}_{x}$ layer is inserted between a TiO layer in contact with the Ti substrate and the $\mathrm{TiO}_{2}$ outer layer. It can be observed in the literature that the experimental conditions reported for anodic oxide growth contain high concentration of ionic species in solution (provided that anodic oxide formation was carried out in aqueous electrolytes). On the contrary, our experimental conditions correspond to a low concentration of a non-adsorbing electrolyte, where the characteristic features of the oxide surface/electrolyte interface observed in the -1.0 to 0 V range are more marked than in other electrolytes [21,22]. The inset in Fig. 1 shows the $j-E$ profiles in a narrowed range, where the current main features can be
associated only to the $\mathrm{Ti}(\mathrm{IV} / \mathrm{III})$ processes for these characteristic potentials. This may indicate the absence of hydrogen-containing species that could be incorporated in the oxide [7], and/or low content of non-stoichiometric species in the $\mathrm{TiO}_{2}$ lattice $[7,32,33]$. $\mathrm{Ti}_{2} \mathrm{O}_{3}$ may also be formed but in contact with water it is unstable and rapidly transforms to $\mathrm{TiO}_{2}[34,35]$, which is more likely to occur at a low concentration of the electrolyte. The inset in Fig. 1 also shows differences during the negative potential scan, where the reduction of $\mathrm{Ti}(\mathrm{IV})$ (hydr)oxide species at the surface in contact with the electrolyte to $\mathrm{Ti}(\mathrm{III})$ (hydr)oxide species occurs and is related to the re-oxidation process obtained during the positive potential scan, this current being more pronounced for foil ${ }^{(1)}$. The hydrogen adsorption-desorption process also takes place, demonstrating that there are several processes associated to this potential range. We consider that acording our experimental setup, the initial oxide thickness could be the main factor that determines the characteristic CV response of the first cycle, as is shown in the section 3.3. After anodic oxide growth, the films were subjected to several repetitive potentiodynamic cycles at $50 \mathrm{mVs}^{-1}$ between -1.0 V and 0 V , showing that the electrochemical behavior of the passive electrodes remained unchanged (data not shown). Interestingly, the magnitude of the cathodic/anodic currents was different for each sample, but the observed shoulders appear always at the same potential values. The second range studied for these $j-E$ profiles was from 5.0 V to 60 V . During the formation process, the anodic $\mathrm{TiO}_{2}$ films started dissolving from the foil edges towards the center of the electrode. This observation was visible from 10 V , where a dissolution-precipitation mechanism took place simultaneously with the formation of the $\mathrm{TiO}_{2}$ lattice. The former led to particle deposition, which was observed by SEM, as is shown in the next section.

### 3.2. Differences observed in anodic oxides formed at the two substrates

The potentiodynamic formation-dissolution process of the oxide was studied at both substrates, with the surfaces analyzed by SEM and EPMA after different $E_{f}$ values were applied. Fig. 2 shows SEM images of titanium foils before (Fig. 2a) and after (Fig. 2b, 2c and 2d) anodic oxide formation. Titanium foil displays a homogeneous surface at the initial state (Fig. 2a), where numerous stripes can be observed, most likely induced by mechanical polishing. Fig. 2b shows an SEM image of the electrode surface corresponding to the interface region with a magnification of $2,300 \mathrm{x}$, where the right side corresponds to the non-immersed surface. Scanning towards the left, a different topography becomes evident and a well defined boundary edge is observed, where the brighter zone indicates the region where the anodic oxide was formed (corresponding to the electrode surface that was immersed in the electrolyte). Fig. 2c shows the anodic oxide zone within $10 \mu \mathrm{~m}$ from the interface. It can be observed that the stripes almost disappear, probably because the oxide growth rate is higher in those regions, and a stable long-range lattice of $\mathrm{TiO}_{2}$ is formed. When the anodic oxide zone analyzed was located far from the interface (Fig. 2d), the SEM image showed a great number of deposited particles, the characteristic size of which ranged between a few micrometers to about $100 \mu \mathrm{~m}$. These particles lay on a smoother surface compared to the surface corresponding to the other zones. The lamina-shaped particles could be interpreted as if they were formed from a side reaction that took place when the crystal was growing. Therefore, the film consists of a three-dimensional $\mathrm{TiO}_{2}$ compact layer, where $\mathrm{TiO}_{2}$ particles were deposited after the dissolution of a previously formed oxide and then a precipitation process occurred. Contrarily to the observations made by Jaeggi et al. [20], we found that the dissolution of anodic oxide is faster at the electrode edges for foil substrates. In addition, with increasing anodization potential i.e.,
increasing oxide thickness, pores cease to grow throughout the oxide, as was the case observed by Jaeggi et al. (data not shown). These facts could have a direct effect on anodic oxide thickness and its physicochemical properties. The glass/Ti sample has a regular texture composed by small grains ( $\sim 500 \mathrm{~nm}$ ) as shown in Fig. 3. The concentration of grain boundaries and the edges per unit area are visibly large and result in a high density when compared to the titanium foil surfaces. The local free energy of the anodic oxide formed on these grain boundaries, differs from that of bulk oxide, due to interfacial free energy contributions and the strength of the non-homogeneous electric field. As a result, these differences are more evident as the field becomes stronger, increasing the activity of the processes considered so far. This may be the main reason why oxide films could only be formed on titanium foil substrates with potentials higher than 5.0 V . The cross section of initial oxide and the anodic oxide corresponding to $E_{f}=3.0 \mathrm{~V}$ of these samples were also observed by SEM, but it was difficult to find a regular oxidized region after cutting the glass/ $\mathrm{Ti} / \mathrm{TiO}_{2}$ sample which suggests that the films were fragile, with low adhesion to the substrate, and could easily become detached (break off) from the metal/oxide interface. Therefore, the differences observed between the two substrates with respect to the $E_{o c p}$ values (see Table 1), could be a result of initial film morphology. The increase of thickness in this potential range, does not make a large difference in the morphological features of foil substrates, but in glass substrates the higher density of grains could produce an accelerated growth of the film through the joining of craters and pores. Nevertheless, the average thickness of the anodic films presented considerable local variations in thickness, especially in relation to the increase in voltage and variations in the proximity to the interface region. As a general trend, measurements showed that the films became thicker when $E_{f}$ increased, as will be shown in the next section. Film thicknesses obtained from voltammperometric measurements demonstrates that there is a linear relationship between oxide growth and $E_{f}$ when the double-layer and oxygen evolution currents are substracted from the $j_{T}$. With a naked eye, it can be observed that the anodic oxide formed from foils has less inhomogeneities at an anodization potential of 20 V compared to the corresponding oxide formed at higher potentials. Earlier works had always associated pores with nano-crystallites but recent studies demonstrated that the pores are formed by $\mathrm{O}_{2}$ bubbles in oxides formed in $\mathrm{H}_{2} \mathrm{SO}_{4}$ and phosphate-containing electrolytes [20].

### 3.3. Ellipsometry

Ex-situ ellipsometry measurements for different samples were carried out after the anodic oxide was potentiodynamically formed at different $E_{f}$ values. This procedure is less accurate than a CV measurement in which the same sample is used to sequentially form oxide of different thicknesses on its surface by applying an increasingly positive potential signal. The measurements were performed ex-situ since oxygen evolution produces light dispersion and interferes with the optical measurements at potentials higher than ca. $1.5-2 \mathrm{~V}$. The oxides were grown in a wide range of potentials (from 1.0 V to 60 V ) using a different sample (titanium foil ${ }^{(2)}$ ) for each $E_{f}$ value. This procedure provides a wide range of oxide thickness. This, in turn, allowed a very good fitting of $\Psi$ and $\Delta$ data values for obtaining reliable values of both optical properties and thickness values (see Fig. 4 below).

Fig. 4a shows the average $\Psi$ and $\Delta$ angles recorded for the $\mathrm{Ti} / \mathrm{TiO}_{2}$ /air interface after anodization to different $E_{f}$ values. Although the crystalline $\mathrm{TiO}_{2}$ oxides have highly anisotropic optical properties [36], all data sets could be reasonably fitted, with small error values, employing the anisotropic single layer model, which showed only a very slight anisotropy in the k values. The solid line in Fig. 4a represents the simulated curve for film growth


Fig. 2. SEM images of different magnifications of titanium foil electrodes. $E_{f}=40 \mathrm{~V} \mathrm{vs} \mathrm{Ag} / \mathrm{AgCl}(\mathrm{sat} \mathrm{KCl})$ in 0.010 M HClO 4 . a) initial oxide. b) interface region. c) anodic oxide zone within the $10 \mu \mathrm{~m}$ from the interface. d) anodic oxide zone more than $10 \mu \mathrm{~m}$ from the interface. The deposits observed in (d) correspond to $\mathrm{TiO}_{2}$-particles.
up to a thickness of around 150 nm using the optical parameters obtained from the best fit: $n_{\text {ord }}=2.38-0.004 i ; n_{\text {extr }}=2.38-0.006 i$. The value obtained for the refractive index of the anodic film ( $n_{f}$ $=2.38$ ) is lower than those reported for tetrahedral anatase ( $n_{\alpha}$ $=2.488 ; n_{\beta}=2.561$ ), tetrahedral rutile ( $n_{\alpha}=2.609 ; n_{\beta}=2.900$ ) or orthorhombic brookite ( $n_{\alpha}=2.583 ; n_{\beta}=2.584 ; n_{\gamma}=2.700$ ) [36]. The lower value obtained for the refractive index of the anodic oxide formed in the present conditions as compared with known crystalline oxides can be attributed to hydration, the incorporation
of anions or the lack of a crystalline structure [ 37,38 ]. On the other hand, the low values obtained for the extinction coefficients ( $\mathrm{k} \approx$ 0 ) indicate that the films are practically transparent, as found for other oxides formed on "valve" metals [37-39]. Thickness values, $d$, obtained also from the fitting procedure show a linear dependence with the formation potential (Fig. 4b). Under the anodizing and measurement conditions employed, an anodized rate $\alpha=2.5 \mathrm{~nm}$ $V^{-1}$ was obtained in the $1 \mathrm{~V} \leq E_{f} \leq 10 \mathrm{~V}$ range, which is of the same order of magnitude as those obtained by other authors ([19]


Fig. 3. SEM images of glass $/ \mathrm{Ti} / \mathrm{TiO}_{2}$ electrodes. a) initial $\mathrm{Ti} / \mathrm{TiO}_{2}$. b) $\mathrm{Ti} / \mathrm{TiO}_{2}\left(E_{f}=3 \mathrm{~V}\right)$.


Fig. 4. a) $\Psi-\Delta$ ellipsometric response during the growth of $\mathrm{TiO}_{2}$ films onto Ti foils up to different final potentials, $E_{f}$. The line represents the simulated curve connecting the anisotropic single-layer model with $n_{\text {ord }}=2.38-0.004 i ; n_{\text {extr }}=2.38-0.006 i$. b) Thickness $/ E_{f}$ dependence obtained from the fitting procedure. c) Mass-thickness relationship.
and references therein). A value of 1.7 nm was estimated for the thickness of the initial $\mathrm{TiO}_{2}$ film, from the extrapolation to the value of $E_{\text {onset }}^{1}$. This value for the spontaneously formed oxide is similar to that reported for other oxides [37-39]. The mass of $\mathrm{TiO}_{2}$ per $\mathrm{cm}^{2}$ generated during oxide formation was obtained from the anodic charge. The mass-thickness relationship (Fig. 4c) allowed then to estimate an oxide density value of $3.44 \mathrm{~g} \mathrm{~cm}^{-3}$. This value is lower than those reported for crystalline $\mathrm{TiO}_{2}$ (3.84; 4.26 and $4.17 \mathrm{~g} \mathrm{~cm}^{-3}$ for anatase, rutile and brookite, respectively [36] indicating that anodic films are amorphous and highly hydrated, as can be concluded from the optical properties obtained. This agree with results obtained from X-ray diffraction for $\mathrm{TiO}_{2}$ films formed in phosphate media, where crystalline orthorhombic $\mathrm{TiO}_{2}$ were detected for potentiostatically aged titanium oxide films while an amorphous structure was found for cycled oxide films [40].


Fig. 5. X-ray emission of $\mathrm{Ti} / \mathrm{TiO}_{2}$ samples corresponding to different final potentials.

### 3.4. EPMA

Fig. 5 shows X-ray emission spectra as a function of photon wavelength $\lambda$ of $\mathrm{Ti} / \mathrm{TiO}_{2}$ electrodes. The peak intensity of $\mathrm{O} \mathrm{K} \alpha$ increases as $E_{f}$ values increase. The intensity corresponding to the $\mathrm{Ti} \mathrm{L} \alpha, \beta$ and $\mathrm{L}_{l, \eta}$ doublets show an opposite trend. Both results were expected, since the thicker the oxide layer, the more oxygen atoms are expected to be available for the incoming electrons and less titanium atoms from the substrate can be ionized by the electron beam.

### 3.4.1. Determination of the growth rate

The net maximum values of normalized intensity (I) corresponding to the $\mathrm{O} K \alpha$ peaks were obtained after background subtraction by the beam intensity current (always around 80 nA ). The acquisition time was $1.3 \mathrm{sec} /$ channel for all the spectra. The results are plotted against $E_{f}$ in Fig. 6. A linear behavior can be observed for the thinnest oxide layers, for which the effect of selfabsorption and electron scattering can be neglected under a thin film hypothesis (TFH). This hypothesis of linearity between oxygen peak intensity and mass thickness was studied by Monte Carlo simulations of 4 keV electrons impinging on a $\mathrm{TiO}_{2}$ film lying above a semi-infinite Ti substrate, using the software package PENELOPE2011. The corresponding results can be seen in Fig. 7. A linear fit, with zero intercept, can be fitted with a regression coefficient


Fig. 6. Net maximum at the $0-K \alpha$ peak normalized by the beam current as a function of $E_{f}$. Dots: experimental data; solid line: linear function fitted to the first three points; dashed line: fit to the whole range in the abscissa. Squares: $\mathrm{Ti} / \mathrm{TiO}_{2}$ prepared as explained in sec. 2.1; triangle: $\mathrm{Ti} / \mathrm{TiO}_{2}$ prepared using a different procedure.


Fig. 7. Monte Carlo $\mathrm{O}-\mathrm{K} \alpha$ peak intensities vs. $\mathrm{TiO}_{2}$ layer thickness to check the thin film hypothesis.
of 0.99987 , including data up to 30 nm . The observed linearity suggests that the TFH assumed should be met in this range. The linear function obtained in Fig. 6 can be expressed as:
$I=8.40+2.35 E_{f} / V$
The linear range of thicknesses $d$ shown in Fig. 7, along with the results obtained by ellipsometry to relate $d$ with $E_{f}$, allowed us to assume a linear behavior of $I$ as a function of $d$ in the range of validity of Eq. (5), i.e., up to $E_{f}=4 \mathrm{~V}$. Assuming 1.7 nm for the spontaneous oxide (i.e. the thickness corresponding to the onset of 0.89 V ), as suggested by the ellipsometry results, Eq. (5) evaluated in $E_{f}=0.89 \mathrm{~V}$ leds to $10.49=c 1.7$, where the proportionality factor $c$ amounted to an increment $\Delta I$ of $6.2 \mathrm{~nm}^{-1}$. On the other hand, Eq. (5) leads to $\Delta I=2.35 E_{f} / V$. From these results, a growth rate $\Delta \mathrm{d}=0.38$ $\Delta V \mathrm{~nm} / \mathrm{V}$ was found. This figure was much lower than expected, which suggests that the sample was altered by the electron beam. In fact, a possible migration of oxygen atoms from the beam impact point would affect the growth rate determined by EPMA.

### 3.4.2. Oxide reduction by electron irradiation

Oxide reduction by electron irradiation was studied by monitoring the specimen current $I_{\text {spec }}$, as mentioned in section 2.5.3; the results of a $\mathrm{Ti} / \mathrm{TiO}_{2}$ electrode, where the oxide was grown anodically at $E_{f}=30 \mathrm{~V}$, are shown in Fig. 8. Even when the electrolyte used to prepare this sample was different from the one employed for the other samples studied here, the decrease observed in $I_{\text {spec }}$ depended only on the nature of the sample and not on its pretreatment. The fact that the nature of this sample was not strongly dependent on


Fig. 8. Evolution of the specimen current during spectrum acquisition for $\mathrm{Ti} / \mathrm{TiO}_{2}$.
the preparation procedure is corroborated by the data plotted in Fig. 6, where its $\mathrm{O} \mathrm{K} \alpha$ response (triangle) is indistinguishable from the corresponding data for the sample prepared as indicated in section 2.1 (squares). As can be seen in Fig. 8, after initial instability probably due to electric charging, a monotonous decrease of around $5 \%$ occurred, revealing a change in the sample, particularly, at its surface. It must be emphasized that the beam current $I_{0}$ was measured before and after the irradiation, and that it was markedly stable, even over periods of several hours. For this reason, the only magnitudes that can be thought to vary in the Eq. (3) of section 2.5.3 and to cause the variation in $I_{\text {spec }}$ are the $\eta$ and $\delta$ coefficients. Moreover, the region where backscattered electrons were produced was much larger than the source of secondary electrons. The latter corresponds to the surface, i.e., to the oxide layer, while the former mainly comprised the Ti substrate, which did not change during the experiment. Thus, the variation of $I_{\text {spec }}$ can be attributed to a variation of $\delta$, which, in turn, was due to the modification of the oxide layer at the beam impact point.

According to calculations made with aluminium, the production of secondary electrons is mainly caused by plasmon decay and by the excitation of conduction electrons [41]. Therefore, it is reasonable to consider that when the number of nearly free electrons increases in the irradiated material, i.e., the sample becomes a better conductor, the coefficient of secondary electrons will be greater. $\mathrm{TiO}_{2}$ is an n-type semiconductor; when heated by the incident electron beam [42], atoms of the electronegative element ( O ) are loosened leaving an excess of atoms of the electropositive element (Ti). These extra atoms introduce additional energy levels from which it is easy for the electrons to jump to the conduction band. In particular, the conductivity $\sigma$ varies with the temperature $T$ according to $\operatorname{Aexp}(-\epsilon / k T)$, where $A$ is a constant, $k$ is the Boltzmann constant and $\epsilon$ is the energy necessary to create an electron-hole pair. This behavior was experimentally verified by Earle [42] and by Padma et al. [43] for the particular case of $\mathrm{TiO}_{2}$ films grown on pure Ti. In light of the previous discussion, it can be concluded that the observed decrease of $I_{\text {spec }}$ can be attributed to a reduction of the oxide layer at the beam impact region. This result is in agreement with the work of Kern et al. [44]; they found that when a $\mathrm{TiO}_{2}$ foil supported by a steel substrate is irradiated by electrons, a reduction to $\mathrm{Ti}_{2} \mathrm{O}_{3}$ is observed for a current density of $0.8 \mathrm{~A} / \mathrm{cm}^{2}$. It must be noted that a lower current density of $0.32 \mathrm{~A} / \mathrm{cm}^{2}$ was used here. The variation of $I_{\text {spec }}$ shown in Fig. 6 allows to observe the reduction process during electron irradiation; as can be seen, reduction took place during the whole acquisition period, which comprised around two hours.

### 3.4.3. Correction of the EPMA results

Ex-situ ellipsometry measurements require a family of samples prepared by the same procedure involving several different thicknesses to be able of determine a relationship between the optical properties measured and the film thickness. On the other hand, to determine film thickness by EPMA, only the unknown sample is required, provided a calibration was previously performed with similar samples, independently of their method of preparation. Unfortunately, as discussed in subsection 3.5.2, the electron probe used to investigate the sample also produced its modification. For these reasons, it is of interest to find a relationship between the thicknesses obtained by EPMA, which involve a systematic error, and the more reliable data determined by ellipsometry. Thus, performing the adequate correction, a further determination by EPMA in a $\mathrm{TiO}_{2}$ layer on a pure Ti substrate would produce reliable values of thickness, independently of the technique used to prepare the sample. The dashed curve shown in Fig. 6 represents a fit to all the data up to saturation, the expression obtained being
$I=\alpha+\beta \tanh \left(\gamma E_{f}\right)$

Table 2
Thicknesses obtained for different final potential of oxide formation $\left(E_{f}\right)$ calculated from ellipsometry and EPMA techniques.

| $E_{f}(\mathrm{~V})$ | $d$ EPMA $(\mathrm{nm})$ | $d$ ellipsometry $(\mathrm{nm})$ |
| :--- | :--- | :--- |
| 1.0 | 1.1 | 2.0 |
| 2.0 | 5.1 | 4.5 |
| 4.0 | 10.2 | 9.5 |
| 8.0 | 19.0 | 19.5 |

with $\alpha=8.976, \beta=36.86$ and $\gamma=0.05625 V^{-1}$. On the other hand, as shown in subsection 3.3, the results obtained by ellipsometry exhibit a linear behavior between thickness and formation voltage in the $0-10 \mathrm{~V}$ range:
$d=a+b\left(E_{f}-0.89 \mathrm{~V}\right)$
with $a=1.7 \mathrm{~nm}$ and $b=2.5 \mathrm{~nm} / \mathrm{V}$. By eliminating the variable $E_{f}$, thickness can be related to the measured $\mathrm{O} \mathrm{K} \alpha$ maximum as follows:
$d=b\left[\frac{1}{2 \gamma} \ln \left(\frac{1+\frac{I-\alpha}{\beta}}{1-\frac{I-\alpha}{\beta}}\right)-0.89 \mathrm{~V}\right]+a$
The thicknesses obtained by EPMA and corrected (Eq. (8)) are presented in Table 2, along with the corresponding results obtained by ellipsometry (Eq. 7). The only values shown belong to the 0 10 V range, for which Eq. (7) is valid. The correction given in Eq. (8) produced results close to the ellipsometry values, except for the lowest voltage considered, for which the relative uncertainties are greater.

## 4. Conclusions

From the results reported here, a set of conclusions can be drawn concerning anodic oxide formation on titanium electrodes:

- The CV technique proved to be more useful than both, EPMA and ellipsometry, for determining differences in the thickness of the native oxide layer. The delay observed in the current onset for foil ${ }^{(2)}$ and glass/Ti/TiO $O_{2}$ electrodes, compared to foil ${ }^{(1)}$ samples (which were kept under controlled atmosphere), indicates a thicker initial oxide. For the last samples an initial thickness of 0.2 nm was calculated considering their $E_{\text {onset }}^{1}$ value, whereas a value of 1.7 nm was obtained for foil ${ }^{(2)}$ samples by ellipsometry.
- The presence of a $\mathrm{Ti}(\mathrm{IV} / \mathrm{III})$ process and the stability of the oxide layer under our experimental conditions were also well characterized by CV. The repetitive cycling in the -1.0 to 0 V potential range did not produce modifications in the characteristic response of (hydr)oxide surface groups, where both the anodic and cathodic charge transfer processes remained at the same peak potential values.
- High accuracy results were obtained by ellipsometry for the oxides electrochemically grown. The information about film thickness of foil and glass/Ti/TiO 2 substrates obtained by this technique corresponded to the average of a large area, averaging local fluctuations in film thickness. Moreover, the sample was not altered during the analysis. Finally, the ellipsometry technique requires a family of samples prepared by the same procedure involving several different thicknesses to be able of determe a relationship between the optical properties measured and the film thickness.
- The $\mathrm{O} \mathrm{K} \alpha$ peak intensity increased as $E_{f}$ increased, which was used as a characterization procedure. The analysis of the $I_{\text {spec }}$ as a function of time allowed the detection of a modification in the original oxide caused by electron beam irradiation during sample analysis at the beam impact region, which resulted in a reduction of $\mathrm{TiO}_{2}$ to $\mathrm{Ti}_{2} \mathrm{O}_{3}$. To overcome the obstacle of this modification of
the sample during the study, a recalibration of the data was performed using the more accurate data from ellipsometry. After this correction, reliable results could be obtained by EPMA for single samples of $\mathrm{TiO}_{2}$ films on titanium, that were prepared using different electrochemical conditions. The procedure of preparing a set of samples with the same optical properties was then avoided.
- Reliable results of titanium oxide electrochemically grown up to different thicknesses in $0.010 \mathrm{M} \mathrm{HClO}_{4}$ were obtained by EPMA using ex-situ ellipsometry as a reference technique.
- There are many applications using EDMA that determine surface morphology, microstructure and chemical composition of anodic titanium oxide [45,46]. We developed a non conventional use for EPMA, where the thickness of thin films could be determined. However, limitations to respect of materials composed of several elements, produce x-rays with overlapping peak positions (by both energy and wavelength) that must be separated.
- In principle, the developed method can be applied to other stratified materials, provided the sample damage caused by the electron impact is not too high. The problem of a possible peak overlapping can always be overcame using a wavelength dispersive spectrometer and an adequate peak selection.


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