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Strategies to tailor the architecture of dual Ag/Fe-oxide nano-heterocrystals - Interfacial and morphology effects on the magnetic behavior

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3 **Strategies to tailor the architecture of dual Ag/Fe-oxide nano-heterocrystals - Interfacial**
4 **and morphology effects on the magnetic behavior**
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Abstract

Bifunctional nanostructured architectures have shown appealing properties, since a single entity can combine the diverse properties of its individual constituents. Particularly, by growing Fe-oxide domains over Ag nanoparticles, the plasmonic and superparamagnetic properties can be combined in a single particle. Beyond the multifunctionality of this system, there are several properties that emerge from intrinsic factors, such as: interface and/or morphology. In this study, we present the synthesis protocols to obtain two sets of heterocrystals, each one with different morphology: dimer and flower-like. In addition, the magnetization behavior of these hybrid nano-heterocrystals is investigated and discussed. These nanomaterials were built by a seed assisted heterogeneous nucleation process, carried out in organic solvents of high boiling point, using the same batch of silver nanoparticles with a mean size of 6 nm as seeds, and tuning the electron-donor capacity of the reaction environment at the thermal decomposition of the iron precursor. Ag/Fe₃O₄ heterocrystals with dimer and flower-like morphologies were obtained. The synthesis protocols for generating these types of nanomaterials are discussed step-by-step. Structural and morphological properties were determined by transmission electron microscopy (TEM), X-ray diffraction (XRD) and X-ray absorption fine structure (XAFS). DC magnetization results suggest that the silver/magnetite coupling generates an increase of the blocking temperature in comparison to those obtained from pure magnetite. This behavior could be linked to a possible increase in the magnetic anisotropy produced by an additional disorder at the Ag-Fe₃O₄ interface. The higher interface area of the Ag/Fe₃O₄ heterocrystals with flower-like architecture leads to a higher blocking temperature and a stronger magnetic anisotropy. These results are supported by AC susceptibility data.

Keywords: hybrid nanostructures, interface effects, magnetic-plasmonic properties, superparamagnetism.

Introduction

In the last years, a prominent part of the research in nanochemistry has focused on the preparation of multifunctional and heterostructured nanocrystals, capable of combining two or more inorganic materials into a single system through a direct atomic bond at the interface and without the presence of a molecular bridge [1-3]. In these new heterocrystals, the properties of the different components can be combined or can give rise to new features associated to morphology or interface effects, factors that may partially govern the physicochemical behavior of these new kind of systems.

Nanoparticles that combine a noble metal (Au, Ag) with ferrite type oxides (Fe_3O_4 , CoFe_2O_4 , Co_3O_4 , MnFe_2O_4) are among the systems that have gained major interest. In these peculiar nanostructures, the optical properties of metallic nanoparticles and the magnetic properties of Fe-oxide counterparts can be coupled. This dual behavior opens the door to new applications in several fields, including catalysis [4, 5], antimicrobial agents [6, 7], magnetic and photo/magnetic hyperthermia therapies [8, 9] and surface-enhanced Raman scattering substrates (SERS) [10], among others.

The most common process to produce heterostructured nanocrystals is the so-called seed-assisted method, where a second component is heterogeneously nucleated over previously formed nanoparticles (“seeds”) of a first component. In this process, the heterogeneous nucleation of the second material is favored over homogeneous nucleation, especially when there is some compatibility between the cell parameters of the crystals [10]. Depending on the nucleation events and the growth evolution, the formed heterostructures can show distinct architectures, such as core-shell, dimer or flower-like. Many efforts have been dedicated to the understanding of the synthesis conditions and mechanisms that lead to these different morphologies [2, 11-13], mainly because the macroscopic properties of the system can be controlled or tuned by the interface type and the interactions between the different constituents.

In this work, we study the synthesis and the magnetic response of two Ag/ Fe_3O_4 heterostructures with different morphologies: dimer and flower-like type, prepared by a seed-assisted route in

organic solvents of high boiling point. In particular, we investigate how the solvent nature, associated to its electron-donor capacity, determines the final morphology. Correspondingly, we performed a complete magnetic characterization by means of DC and AC magnetization techniques to understand how the heterostructures architecture and the distribution of the Fe_3O_4 domains define the macroscopic behavior of the samples. A discussion on the structural and optical properties is also presented to reinforce the results and the conclusions.

Experimental Details

Materials. Silver nitrate (AgNO_3 , $\geq 99\%$, Sigma-Aldrich), sodium acetate (NaCH_3COO , $\geq 99\%$, Anedra), oleylamine ($\text{C}_{17}\text{H}_{31}\text{-CH}_2\text{NH}_2$, $\geq 99\%$, Sigma-Aldrich), oleic acid ($\text{C}_{17}\text{H}_{31}\text{-COOH}$, 90%, Sigma-Aldrich), 1-octadecene ($\text{C}_{18}\text{H}_{36}$, 90%, Sigma-Aldrich), diphenyl ether ($(\text{C}_6\text{H}_5)_2\text{O}$, 98%, Sigma-Aldrich), iron(III) acetylacetonate ($\text{Fe}(\text{acac})_3$, 99%, Sigma-Aldrich) and 1,2-hexadecanediol ($\text{C}_{16}\text{H}_{34}\text{O}_2$, 90%, Sigma-Aldrich). All chemicals were used as received without further purification.

Methods. $\text{Ag}/\text{Fe}_3\text{O}_4$ nanoparticles were synthesized by a seed assisted heterogeneous nucleation process in organic solvents of high boiling point, as described below.

Silver Seeds Nanoparticles (AgNPs) synthesis: AgNPs were prepared according to previous reports [8, 12]. Briefly, 2 mmol silver acetate (AgCH_3COO) were obtained by stoichiometric precipitation of AgNO_3 and NaCH_3COO . The so-obtained white powder was dissolved in 2 mmol oleic acid, 2 mmol oleylamine and 5 mL 1-octadecene. The mixture was heated up to 175 °C under continuous stirring for 30 min to allow AgNPs formation. Finally, the system was cooled down to room temperature and the AgNPs were precipitated by the addition of ethanol and isolated by centrifugation at 3000 RPM for 15 minutes. The isolated solid was washed 3 times before redispersed in hexane.

$\text{Ag}/\text{Fe}_3\text{O}_4$ heteroparticles formation: AgNPs were used as fixed seeds for Fe_3O_4 heterogeneous nucleation. 1 mmol Ag (as AgNPs), 0.3 mmol $\text{Fe}(\text{acac})_3$, 1.5 mmol 1,2-hexadecanediol, 0.75 mmol oleic acid, 1.5 mmol oleylamine and 5 mL solvent (1-octadecene, ODE or phenyl ether,

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3 PHE) were added. The solution was maintained at 125 °C for 1 h under continuous stirring, and
4 then heated (heating rate = 5 °C/min) up to 260 °C for 60 min under Ar flow to allow Fe₃O₄
5 formation. Finally, the system was cooled down to room temperature and the nanoparticles were
6 precipitated by the addition of ethanol and isolated by centrifugation at 3500 RPM for 30
7 minutes. The isolated solid is finally washed 3 times to then redisperses it in toluene.
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13 Two samples were prepared, Ag/Fe₃O₄ ODE (reaction carried out in 1-octadecene) and
14 Ag/Fe₃O₄ PHE (reaction carried out in phenyl ether). Furthermore, a control sample of
15 homogeneously nucleated Fe₃O₄ without the addition of the Ag seeds was prepared (see
16 supporting information).
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20 21 22 **Characterization**

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25 Phase composition and crystallographic characterization were carried out via X-ray diffraction
26 (XRD) on a standard Rigaku diffractometer with Cu K- α radiation. X-ray absorption fine
27 structure (XAFS) spectroscopy patterns for the Fe K-edge (7.112 keV) were recorded at room
28 temperature in fluorescence mode with the colloidal samples positioned at 45° with respect to the
29 X-ray beam. These measurements were performed at the Brazilian Synchrotron Light Laboratory
30 (CNPEM/LNLS, Campinas, Brazil). Absorption spectra, carried out on colloidal samples
31 (nanoparticles dispersed in toluene), were recorded at room temperature by means of a UV-Vis-
32 NIR Shimadzu UV-1603 using quartz cuvettes with 10 mm optical path length. Morphology and
33 size distribution of the samples were determined by transmission electron microscopy (TEM) at
34 the CNPEM/LN Nano (Campinas, Brazil) in a JEOL JEM 2100F with a field emission gun (FEG)
35 operating at 200 kV and at the IFW (Dresden, Germany) with a FEI-Tecnaï microscope
36 operating at 200 kV. The samples for TEM experiments were prepared by dripping 10 μ L of
37 nanoparticle dispersion onto a carbon-coated copper grid (Ted Pella) and drying the solvent at
38 room temperature. DC magnetization and AC magnetic susceptibility measurements were
39 recorded with a MPMS-XL7 Quantum Design SQUID magnetometer and a PPMS Quantum
40 Design magnetometer, respectively. The samples for magnetic measurements were prepared by
41 mixing 50 μ L of the colloidal dispersion containing the Ag/Fe₃O₄ heterocrystals and one drop of
42 1-octadecene. This mixture was stirred until the toluene was evaporated. After that, a piece of
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3 cotton (appr 15 mg) was impregnated with the Ag/Fe₃O₄ sample and finally placed in the
4 SQUID/PPMS holder sample.
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8 **Results and Discussion**

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11 The two-step procedure employed for the Ag/Fe₃O₄ heteroparticles synthesis can be
12 conveniently used to study the reaction mechanism and the conditions that lead to the formation
13 of different heterostructure morphologies.
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19 The AgNPs (produced separately in a first step) were incorporated to the synthesis medium prior
20 to the thermal decomposition reaction and the formation of the iron oxide phase. One-step
21 reactions have been successfully developed to produce heterocrystals by exploiting the
22 difference between the formation temperatures of the metallic nanoparticle and the iron oxide
23 [14-16]. In spite of this, we have decided to employ a two-step procedure so we can directly
24 relate the heterogeneous nucleation events and final heteroparticle structure with the reaction
25 conditions, i.e. solvent electron-donor capacity, and detach them from other variables, such as
26 AgNPs size and concentration, for example. Reports on two-steps heterostructure synthesis have
27 investigated different options for the iron source, including iron acetylacetonate, iron oleate or
28 iron pentacarbonyl. However, the precursor nature does not seem to affect the final particle
29 conformation because a common transformation to iron oleate occurs at low temperatures (prior
30 to iron oxide nucleation) regardless the original reagent [17, 18].
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41 Fig. 1 shows a representative TEM image of the prepared AgNPs. The synthesized AgNPs show
42 a spheroidal-like shape with a narrow size distribution. Histograms were obtained by counting
43 more than 200 silver nanoparticles from several TEM images, which were well fitted using a
44 lognormal function (Fig. 1(b)). From this, a mean diameter $\langle D \rangle = 6$ nm with a rather narrow size
45 distribution ($\sigma = 0.2$) was determined. These results confirm the well-dispersed nature of the
46 initial AgNPs colloidal suspension obtained in the first step of the synthesis. The *fcc* structure of
47 the silver nanocrystals was examined by X-ray diffraction, as it is shown in Fig. S1
48 (Supplementary Information).
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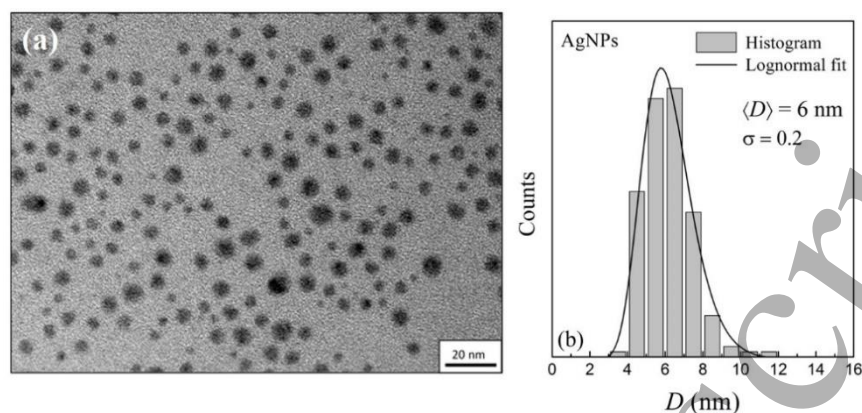


Figure 1. (a) TEM micrograph of AgNPs. (b) Size distribution histogram and the lognormal fit.

Afterwards, these AgNPs were used as fixed seeds for the formation of the Ag/Fe₃O₄ heteroparticles in both reaction solvents. Fig. 2 and 3 show the TEM and HR-TEM of the two samples. The difference in the atomic number between the two components makes it possible to distinguish iron oxide (lighter particles) from Ag (darker particles) in the images. For the heterostructures prepared in 1-octadecene (Ag/Fe₃O₄ ODE) dimer-like heteroparticles are predominant (Fig. 2), with just one iron oxide crystal attached per Ag seed. On the other hand, for heterostructures prepared in phenyl ether (Ag/Fe₃O₄ PHE) flower-like heteroparticles predominate (Fig. 3), with various iron oxide crystals surrounding a central Ag seed. In both TEM images (Fig. 2(a) and Fig. 3(a)), some non-attached iron oxide particles can be seen. However, the presence of these particles decreased dramatically from previous samples by carefully tuning the relation between Ag seeds and Fe precursor and the heating rate. Fig. 2(d) and 3(d) present the HR-TEM images for samples Ag/Fe₃O₄ ODE and Ag/Fe₃O₄ PHE, respectively. Additional images are displayed in the Supporting Information (Fig. S2 and S3). Notice that by applying the Fourier transformation technique on a simple dimer-like particle (inset on Fig. 2(d)), characteristic crystallographic planes for magnetite and silver phases can be detected, i.e., (2 2 2) plane for the Fe₃O₄ and (2 0 0) plane for the Ag. For the sample Ag/Fe₃O₄ PHE, the Fourier transformation (inset on Fig. 3(d)) reveals two characteristic crystallographic planes for magnetite phase: (2 2 0) and (3 1 1). However, none of the crystallographic planes of the metallic silver are distinguished, which is probably related to the absence of the oriented planes of this material in the obtained HR-TEM image.

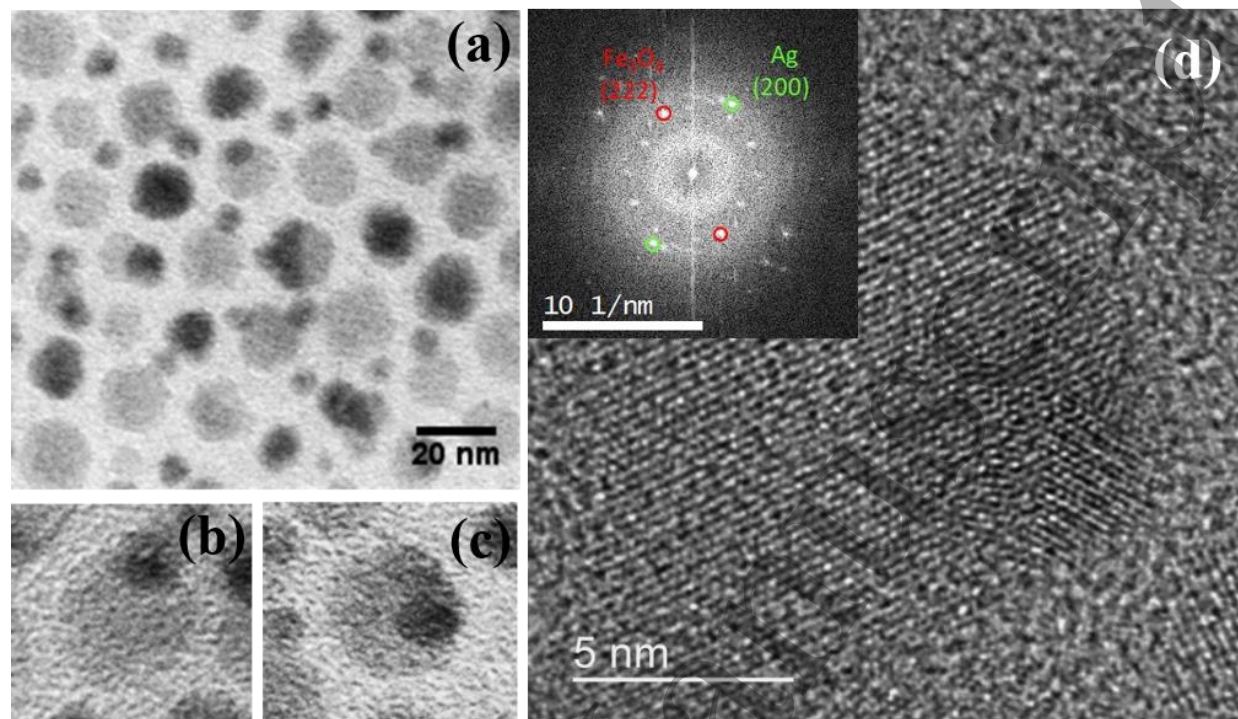


Figure 2. (a-c) Representative TEM images of Ag/Fe₃O₄ heterostructures with dimer-like configuration showing the iron oxide nanoparticles attached to a silver seed. (d) HR-TEM image of a single dimer-like nanoparticle and corresponding Fourier transform (*inset*).

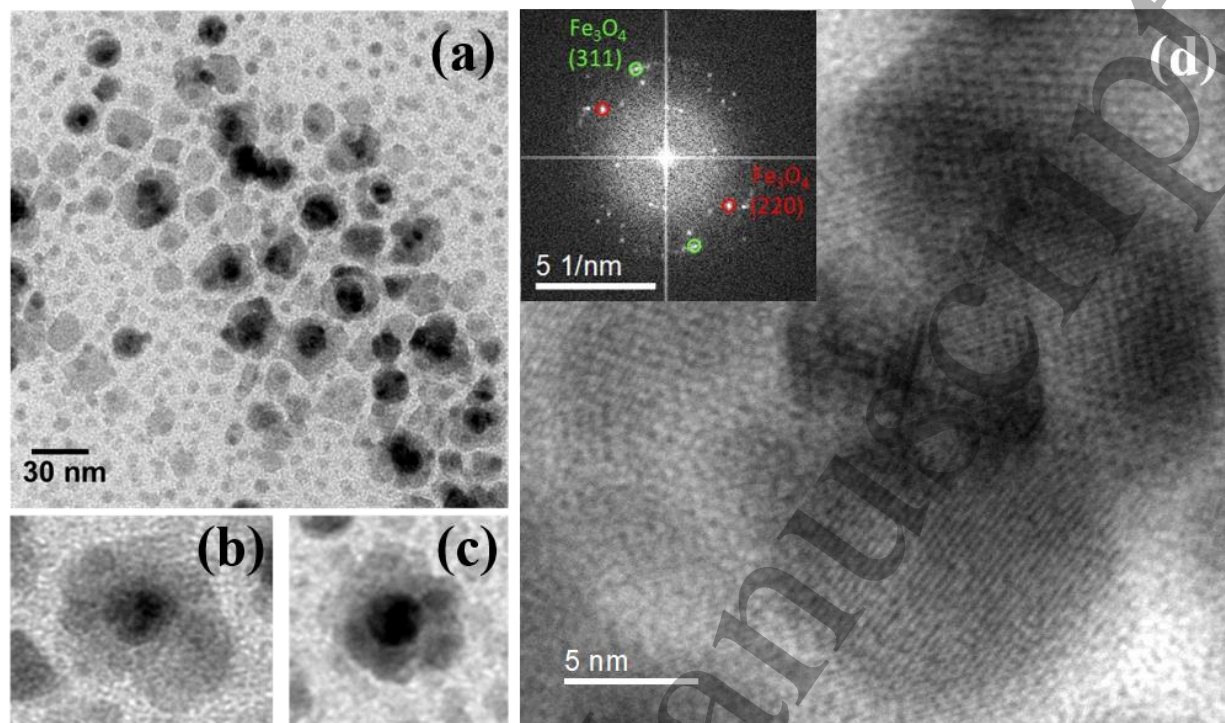


Figure 3. (a-c) Representative TEM images of Ag/Fe₃O₄ heterostructures with flower-like configuration showing the iron oxide nanoparticles surrounding a silver seed. (d) HR-TEM image of a single flower-like nanoparticle. (d) HR-TEM image of a single dimer-like nanoparticle and corresponding Fourier transform (*inset*).

Fig. 4(a) shows the XRD patterns of the dried heterostructures. As expected, two contributions from two sets of diffraction peaks can be observed in the diffractograms. One set can be assigned to the *fcc* spinel-inverse iron oxide phase (Fe₃O₄ and/or γ -Fe₂O₃) and the other one corresponds to *fcc* metallic Ag. Thus, one can confirm the presence of both crystalline structures in the samples.

Optical properties were studied in terms of UV-Visible absorption and the resulting spectra are shown in Fig. 4(b). The surface plasmon resonance peak associated to the Ag seeds in the heterocrystals is weakened compared to that obtained from the measurement of pure AgNPs. This behavior is due to the shielding effect caused by Fe₃O₄ [19, 20]. Comparing both samples, the effect is more pronounced for Ag/Fe₃O₄ PHE than for Ag/Fe₃O₄ ODE. This difference can be related to the morphology dissimilarities, as flower-like structures with a larger Ag-Fe₃O₄ interface can produce a larger shielding effect compared to dimer-like particles. A similar decay

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3 of the absorption peak was also observed in Au/Fe₃O₄ particles [19, 21]. A control sample of 10
4 nm magnetite particles was also investigated for comparison (synthesis and characterization
5 details are presented in the Supplementary Information), the respective UV-Vis spectra (red line
6 in Fig. 4(b)) is consistent with subwavelength sized dielectric spheres [22].
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11 In most reports concerning the studies of surface plasmon resonance on metallic/magnetic hybrid
12 nanoparticles, including Au/Fe₃O₄ and Ag/Fe₃O₄, a red-shift of the metal particle surface
13 plasmon resonance peak (SRP) is evidenced. According to some authors, this behavior results
14 due to the strong coupling of iron oxide, which exhibits higher values of the real part of the
15 refractive index with respect to the organic solvents [23, 24]. Strikingly, the absorption peak of
16 our prepared Ag/Fe₃O₄ heteroparticles stays at almost the same position as that of AgNPs. This
17 behavior was also observed by Huang et. al. [20] and can indicate a weak interaction between the
18 building blocks. Some authors also suggest that Fe₃O₄ growth over the Ag seed does not always
19 occur epitaxially, leading to nearly independent surface plasmonic properties of the metallic
20 nanoparticles [19]. Also, some kind of centrifugation sorting during Ag/Fe₃O₄ washing steps
21 may have privileged AgNPs with a UV-shifted peak position. Moreover, since the modification
22 of the size of the scattering objects are known to affect the intensity and width of the SPR peak,
23 the observed differences could also be related to the minor changes in the Ag seeds size during
24 the second synthesis step.
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37 Measurements of X-ray absorption fine structure (XAFS) were carried out in order to examine
38 the Fe-oxidation states of the iron oxide domains in the heterostructures. A control sample of
39 Fe₃O₄ nanoparticles was also analyzed for comparison. Fig. 4(c) and (d) show the profiles of the
40 X-ray absorption near edge spectra (XANES) and extended X-ray absorption fine structure
41 (EXAFS), respectively, determined for the samples Ag/Fe₃O₄ ODE and Ag/Fe₃O₄ PHE, as well
42 as for the iron oxide control sample. According to these results, the pre-edge energy position, the
43 pre-peak features and the EXAFS profiles of both heteroparticles systems are compatible with a
44 structure that combines ferric (Fe²⁺) and ferrous (Fe³⁺) ions [14]. Furthermore, by comparing the
45 obtained results from the Ag/Fe₃O₄ with those one from the Fe₃O₄ control sample, one can affirm
46 that the environment of Fe-atoms located in the heterostructures are disposed to form an iron
47 oxide with a stoichiometric Fe₃O₄ phase.
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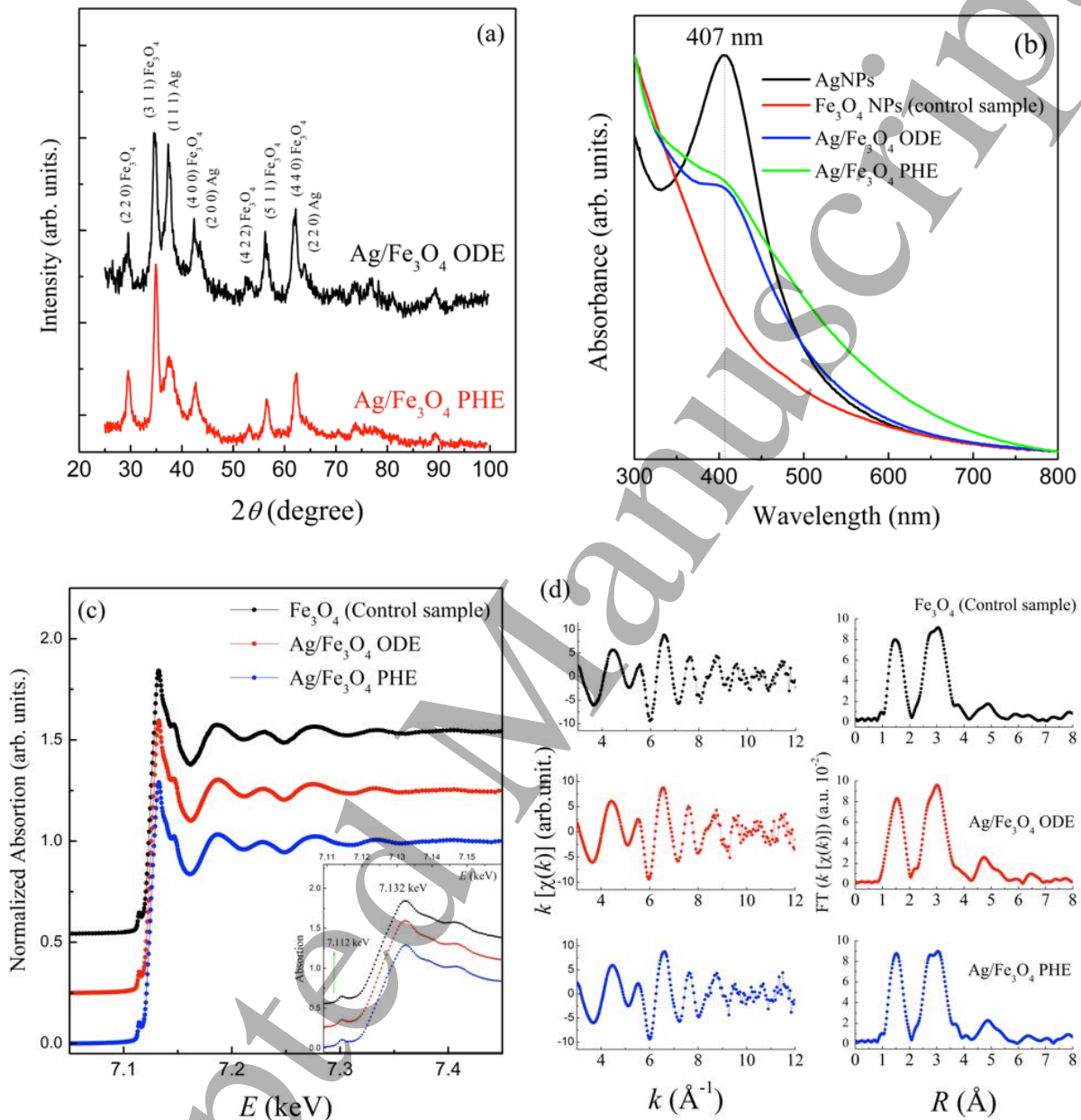


Figure 4. (a) XRD patterns of sample Ag/Fe₃O₄ ODE (dimer-like morphology) and Ag/Fe₃O₄ PHE (flower-like morphology). The phase and Miller index of principal peaks are indicated. (b) UV-Vis absorption spectra of the samples. (c) Experimental EXAFS Fe K-edge spectra for samples Ag/Fe₃O₄ ODE, Ag/Fe₃O₄ PHE and Fe₃O₄. *Inset*: XANES region. (d) Fourier transformations obtained for the studied samples.

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5 The study and control of the reaction conditions leading to different heterostructures
6 morphologies has been a topic of intense research in the last years. In a pioneering work, H. Yu
7 and co-workers proposed that a charge compensation process within the metal nanoparticle (Au
8 in their case) governs and defines the heterostructure conformation [11]. The nucleation of the
9 iron oxide on the Au surface induces a positive charge at the growth interface, immediately
10 compensated by free electrons of the metal. This charge migration produces electron deficient
11 “restrained facets” and denies multi-nucleation process. However, a change in the medium’s
12 electron-donor capacity by using a slightly polar solvent could compensate the charge induced
13 by the polarized growth plane, and thus allow multi-nucleation. Therefore, the modification of
14 the reaction medium from octadecene to phenyl ether can modify the final heterostructure
15 conformation from dimer-like to flower-like nanoparticles. Fig. 5 shows a scheme of the possible
16 pathways for the seeded-growth synthesis of Ag/Fe₃O₄ heterocrystals starting from a mixture of
17 Ag seeds and Fe precursor. This hypothesis was supported by few authors for noble metal
18 (especially Au) heterostructures [11, 12, 20, 25], and also for other metal – metal oxide
19 combinations, such as FePt/MnO [26]. Flower-like nanoparticles were also produced in highly
20 polar solvents like poly-alcohols [10], supporting the hypothesis that the polarity of the synthesis
21 medium influences the morphology of the heterostructure. Nevertheless, additional reports
22 suggest that other reaction conditions like metal seeds – Fe precursor ratio [27] or metal seed
23 crystallinity [20] could also define the heterostructure configuration.
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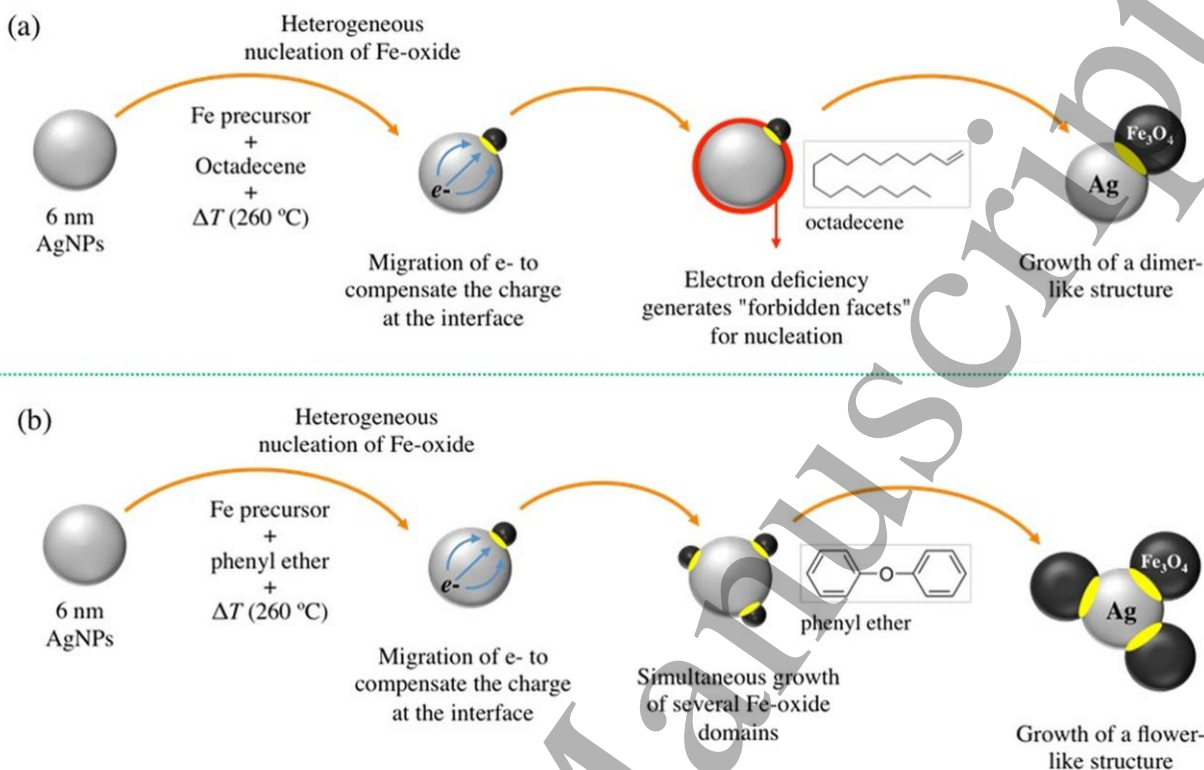


Figure 5. Schematic representation of the two possible pathways for the seeded-growth formation of Ag/Fe₃O₄ heterocrystals with different morphology. Heterogeneous nucleation of Fe₃O₄ to form (a) dimer or (b) flower-like structure.

Magnetization features of both Ag/Fe₃O₄ systems were investigated in order to understand the influence of the nanoparticle morphology on the macroscopic magnetic behavior. Because the magnetic measurements were performed in highly diluted samples, interparticle dipolar interactions are expected to be negligible. However, on the sample with flower-like morphology, the presence of several Fe₃O₄ domains attached to a single Ag seed can lead to intra particle dipolar and/or exchange interactions, because these magnetic domains can be in direct contact with each other and not spatially separated by the diamagnetic matrix.

Fig. 6 shows the magnetization vs. applied field (M vs. H) curves performed at 5 K and 100 K. Notice that the magnetization units are expressed in terms of emu per mass of heteroparticles, i. e., mass of Fe₃O₄ and Ag. While the saturation magnetization (M_s) values are lower than those reported for bulk magnetite, as expected for nanoscale magnetic materials, these are partially

underestimated because the additional mass of silver. The *S*-shaped curves can be related to the presence of monodomain magnetic iron oxide nanoparticles. At $T = 5$ K the magnetization loops display a coercive field of ~ 260 Oe for the Ag/Fe₃O₄ ODE sample and ~ 360 Oe for the Ag/Fe₃O₄ PHE sample, indicating that the nanoparticles are at the magnetically blocked regime at this temperature [28]. As can be noted, the coercive field at 5 K is larger in the Ag/Fe₃O₄ PHE sample than in the Ag/Fe₃O₄ ODE sample. This behavior can be understood from two main approaches: (i) Because the presence of several iron oxide domains attached to the same AgNP in the sample with a flower-like architecture promotes a strengthening of dipole-dipole interactions between the magnetic entities [12], which could produce a magnetic hardening in this sample. (ii) Because the effective coupling between the components can alter the arrangement of those magnetic moments located at the magnetic surface and thus modifies the magnetization dynamics. The coercive field in both samples drastically diminishes at 100 K, an indication that the magnetization behavior is close to the superparamagnetic regime [28].

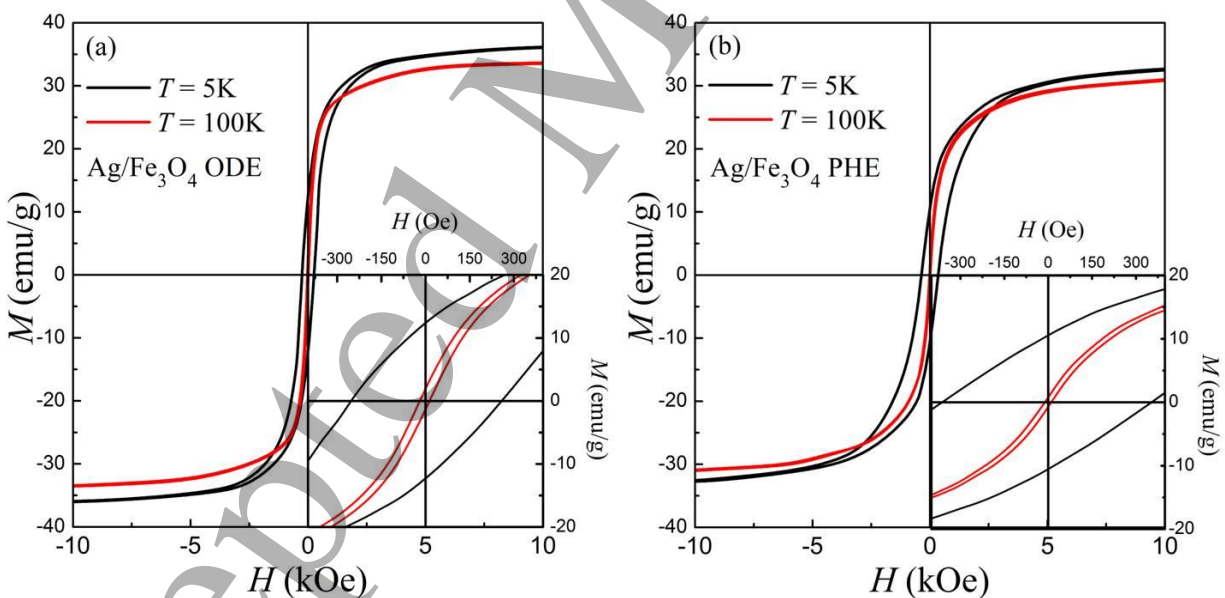


Figure. 6. Magnetization vs. applied magnetic field curves measured at 5 K and 100 K for samples (a) Ag/Fe₃O₄ ODE and (b) Ag/Fe₃O₄ PHE. *Insets*: zoom-in of the low field region showing the coercivity behavior.

Temperature dependence of the magnetization of Ag/Fe₃O₄ heterostructures, recorded under zero-field-cooling (M_{ZFC}) and field-cooling (M_{FC}) protocols, is presented in Fig. 7. In both curves, the ZFC magnetization display a clear and broad maximum value, at $T_m = 117$ K for Ag/Fe₃O₄ ODE and at $T_m = 197$ K for Ag/Fe₃O₄ PHE. Although in principle, T_m can be associated to the blocked-to-superparamagnetic transition of the iron oxide domains, the obtained values are approximately one magnitude of order larger than ones obtained for the magnetite control sample, which was found at $T_m = 27$ K (blue symbols in Fig. 7(a)) and whose shape can be easily linked to a weakly interacting monodomain system [28, 29]. The differences of T_m between the control sample and the more complex nanostructures can be explained considering three different phenomena: changes in the size of the magnetic nanoparticles [30], strengthening of dipolar interactions [31] and variations in the iron oxide phases [32]. However, for both heterostructures the sizes of the magnetic domains are quite similar; hence the magnetic phase remains unaltered and the dipole-dipole and superexchange interactions should not strong enough to generate a shift of T_m of 90 K and 170 K in the Ag/Fe₃O₄ ODE and Ag/Fe₃O₄ PHE samples, respectively (in comparison to the control Fe₃O₄ sample). Such interactions, especially in the sample with flower-like morphology, could affect the effective anisotropy of the structures, leading to a T_m displacement (as discussed previously) but in a weaker degree. Generally speaking, the T_m shift (as well as the ZFC/FC general curve shape) can be influenced by the morphology, metal seeds features and/or interface effects. For instance, some authors suggest that the coupling efficiency between metal/Fe-oxide surfaces can modify the superficial disorder [33, 34], the energy barrier (E_b) and the magnetocrystalline anisotropy (K) of the magnetic counterpart [35], altering the global magnetic behavior of the system. Therefore, two ZFC curves, with values for T_m matching those obtained from the experimental measurements of the Ag/Fe₃O₄ heteroparticles were simulated following the expression reported in [36], which is:

$$M_{ZFC}(T) \cong \frac{M_S^2 H}{3K} \left[\frac{E_{bm}}{k_B} \int_0^{T/T_{bm}} \frac{1}{T} y f(y) dy + \int_{T/T_{bm}}^{\infty} f(y) dy \right], \quad \text{Eq. 1}$$

where $f(y)$ is a distribution of energy barriers (adopted as lognormal-type), $y = E_b/E_{bm}$ being E_b and E_{bm} the energy barrier and the mean energy barrier, respectively. T_{bm} is the mean blocking temperature, M_S the saturation magnetization, H the external magnetic field (50 Oe), K the anisotropy and k_B the Boltzmann constant. These simulations allow us to approximately

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3 quantify the variations on E_{bm} and K_{ef} suggested by the mentioned factors. As expected, the E_{bm}
4 values extracted from the simulated curves (blue dotted lines in Fig. 7) are larger in comparison
5 to the one obtained from the ZFC fit of the Fe_3O_4 control sample: $\sim 3 \times 10^{-13}$ erg for Ag/ Fe_3O_4
6 ODE sample, $\sim 4.6 \times 10^{-13}$ erg for Ag/ Fe_3O_4 PHE sample, and $\sim 8 \times 10^{-14}$ erg for the Fe_3O_4 control
7 sample. It is worth noting that the increase in E_{bm} also raises the anisotropy constant, whose
8 mean value can be obtained from $K = E_{bm}/V$, being V the mean particle volume. Using this
9 expression and assuming a mean diameter of 8.6 nm for the Fe_3O_4 nanocrystals, the K values are
10 9.0×10^5 erg/cm³ and 1.4×10^6 erg/cm³ for dimer and flower-like samples, respectively. These
11 values are appr. 3.5 and appr. 6 times larger than the obtained for the Fe_3O_4 control sample (K
12 from fit procedure is 2.4×10^5 erg/cm³). All the obtained values are summarized in Table 1.
13 Larger values of E_{bm} and K can be partially attributed to an additional superficial disorder at the
14 Ag/Fe-oxide interface. In that interface, magnetic moments could be isotropically pinned and
15 still coupled with the magnetic moments of the magnetite core, generating a magnetic drag over
16 the core that leads to the increase of the anisotropy [37]. Consequently, the coupling effects will
17 be more significant for the sample with a greater interface between the two components,
18 Ag/ Fe_3O_4 PHE in this case.
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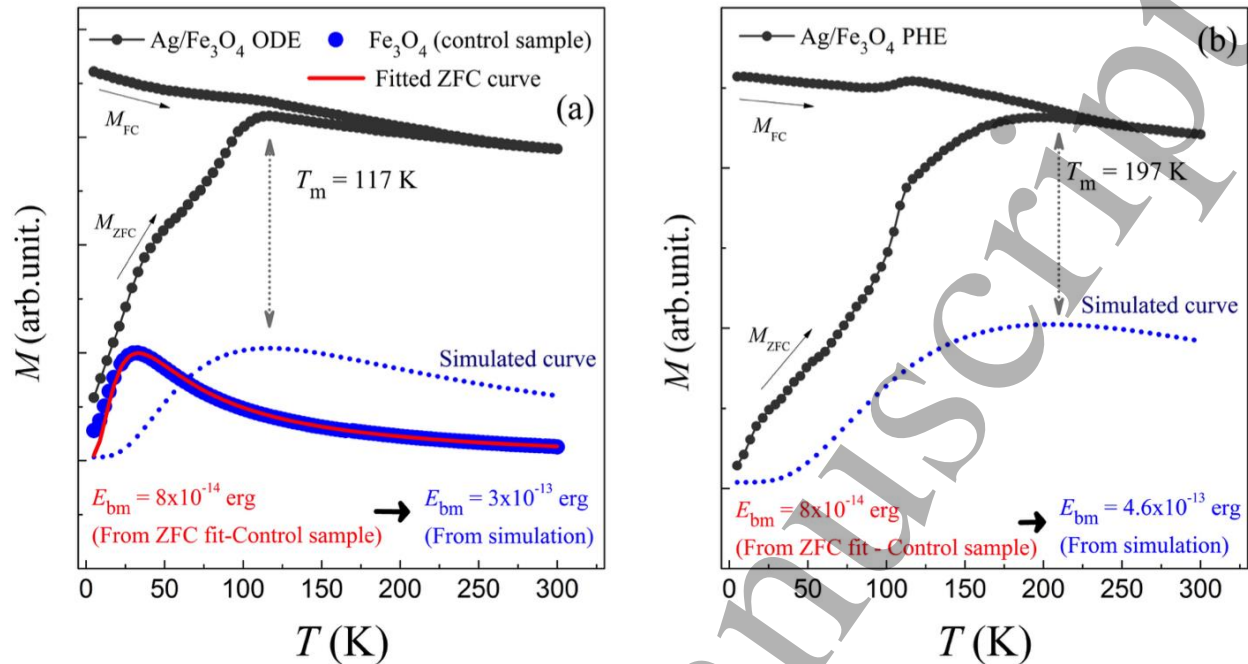


Figure 7. ZFC/FC curves taken in a magnetic field of 50 Oe for samples (a) Ag/Fe₃O₄ ODE and (b) Ag/Fe₃O₄ PHE. Blue filled symbols in (a) indicate the experimental ZFC curve of the magnetite control sample and the red solid line its respective fit according to Eq. 1. Blue dotted lines denote the simulated curves for the heterostructures according to Eq. 1.

Notice that it is not enough to just increase the value of the barrier energy (or K) to reproduce exactly the shape and the temperature evolution of the experimental ZFC curves obtained from the dual Ag/Fe₃O₄ heterostructures. In order to achieve this goal for these complex structures, several other factors need to be considered, such as the dipolar interactions, possible effects of spin or charge transfer [38, 39] and cation exchange effects [40], structural defects [41] among others.

Table 1. Parameters E_{bm} , T_{bm} , K and T_i obtained from the ZFC fit (Fe_3O_4) and from simulations ($\text{Ag}/\text{Fe}_3\text{O}_4$ heterostructures).

Sample	E_{bm} (erg)	T_{bm} (K)	K (erg/cm ³)	T_i (K)
Fe_3O_4 (Control sample)	8.0×10^{-14} ⁽ⁱ⁾	21 ⁽ⁱ⁾	2.4×10^5 ⁽ⁱ⁾	---
$\text{Ag}/\text{Fe}_3\text{O}_4$ ODE (Dimer-like morphology)	3.0×10^{-13} ⁽ⁱⁱ⁾	73 ⁽ⁱⁱ⁾	9.0×10^5 ⁽ⁱⁱ⁾	170
$\text{Ag}/\text{Fe}_3\text{O}_4$ PHE (Flower-like morphology)	4.6×10^{-13} ⁽ⁱⁱ⁾	120 ⁽ⁱⁱ⁾	1.4×10^6 ⁽ⁱⁱ⁾	200
⁽ⁱ⁾ Values extracted from the fit of the ZFC curve. ⁽ⁱⁱ⁾ Simulated values.				

Relevant information can also be qualitatively extracted from the FC magnetization curves. For example, the flattened behavior at low temperatures ($T \leq 50$ K) followed by a wider shoulder as the temperature increases could be a signature of limited nanoparticle aggregation and occurrence of a spin-glass and/or interacting superparamagnetic behavior [29, 42, 43]. Moreover, the manifestation of an irreversibility temperature (T_i), defined as the point in which the ZFC curve bifurcates from the FC one, indicates the presence of a distribution of blocking temperatures, as expected for size distributed magnetic monodomains. In summary, all examined properties suggest that the magnetic behavior of the studied samples are a consequence of a sum of several effects, leading to a complex interpretation of the magnetic results.

AC susceptibility measurements were performed to reinforce the results from DC magnetic characterizations. The real and imaginary components of the magnetic susceptibility vs. temperature (χ' vs. T and χ'' vs. T), recorded at four different frequencies, are shown in Fig. 8. Although the $\chi'(T)$ curves for $\text{Ag}/\text{Fe}_3\text{O}_4$ nanoparticles follow different trends, two main features can be highlighted in both systems: an initial shoulder at low temperatures followed by a broader peak near the T_m value obtained from the ZFC measurements. The nature of these events can be related with two independent blocked-to-unblocked system processes. The equivalent two events in the imaginary component of the magnetic susceptibility $\chi''(T)$ is reflected in the presence of

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3 two peaks, which are labeled as $T_{\text{MAX}(1)}^{\chi''}$ and $T_{\text{MAX}(2)}^{\chi''}$ in Fig. 8 (c and d). The initial shoulder at
4 $\chi'(T)$ is correlated to the first peak in $\chi''(T)$ as expected, and can be assigned to the blocked-to-
5 unblocked transition of those iron oxide nanoparticles non-linked or weakly-linked with the Ag
6 seeds. These temperature values are at similar positions to the ones observed in homogeneously
7 nucleated nanoparticles (control Fe_3O_4 sample), as shown in Fig. S5 (Supporting Information).
8 The second event, evidenced as a peak at $T \approx 110$ K for the Ag/ Fe_3O_4 ODE sample and $T \approx 170$
9 K for the Ag/ Fe_3O_4 PHE sample in the $\chi'(T)$ and $\chi''(T)$ curves, is related to the unblocking
10 processes of the Fe_3O_4 domains attached on the heterostructures or, in the case of flower-like
11 sample with the unblocked-blocked transition of those magnetic nanoparticles with higher inter-
12 particle interaction. Notice that these peaks are registered at higher temperatures for the sample
13 with flower-like morphology, an indication that in this sample the blocking/unblocking process
14 requires more thermal energy to occur. This behavior also supports the previous evidence of the
15 distinct interactions and possible alterations on the magnetocrystalline anisotropy between both
16 systems.
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30 The relative height of the peaks observed in the imaginary component of the magnetic
31 susceptibility has significant differences, i.e., for the Ag/ Fe_3O_4 ODE sample, $T_{\text{MAX}(1)}^{\chi''}$ is higher
32 than $T_{\text{MAX}(2)}^{\chi''}$, while for the Ag/ Fe_3O_4 PHE sample the opposite occurred. This difference can be
33 related to the ratio of homogeneous to heterogeneous Fe_3O_4 domains, indicating that the reaction
34 in phenyl ether avoids the homogeneous nucleation events and is more effective to produce
35 heterostructures.
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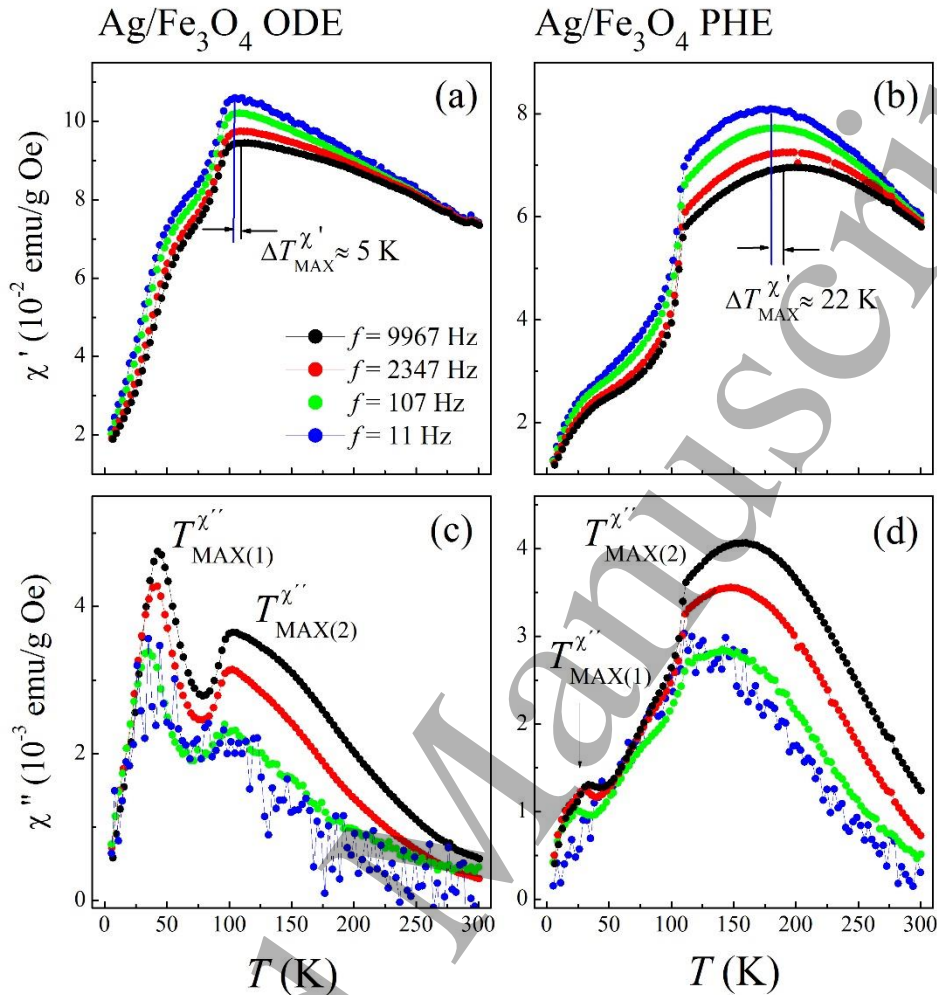


Figure 8. Temperature dependence of the real $\chi'(T)$ and imaginary $\chi''(T)$ components of the susceptibility of the Ag/Fe₃O₄ heterostructures.

From the AC magnetization results, it is also possible to observe T_{MAX} displacements as frequency function. Such displacements are rather small when compared to the ones observed in highly interacting nanoparticle systems, a qualitative indication that the magnetic entities can be in a weakly interacting state [29, 31, 44, 45]. To clarify this point, the frequency dependence of the peaks in $\chi'(T)$ was analyzed for both Ag/Fe₃O₄ systems by fitting the experimental data using the Vogel-Fulcher law, given by:

$$\tau = \tau_0 \exp(E_b/k_B(T - T_0)) \quad (1)$$

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3 where τ is the relaxation time that is associated to the frequency f through $\tau = 1/f$, τ_0 is the
4 characteristic time (ranging between 10^{-9} s and 10^{-12} s [46]), k_B is the Boltzmann constant, E_b is
5 the energy barrier and T_0 is a characteristic temperature that defines if the thermal energy ($T >$
6 T_0) or the interacting energy ($T < T_0$) dominates the magnetic moment dynamics and
7 magnetization process. By plotting the experimental data as $\ln(\tau)$ against $1/T_{MAX}^{\chi'}$ (being $T_{MAX}^{\chi'}$
8 the peak associated to the blocking-unblocking processes of heterogeneous nucleated Fe_3O_4
9 domains) and using Eq. 1 to fit them (Fig. 9), parameters with coherent physical meaning were
10 obtained. For E_b , the determined values are 8×10^{-14} erg for Ag/ Fe_3O_4 ODE and 2×10^{-13} erg for
11 Ag/ Fe_3O_4 PHE, which are closer to the ones extracted from the analysis of the ZFC curves.
12 Moreover, the obtained characteristic time τ_0 ($\sim 10^{-11}$ and $\sim 10^{-10}$ s for the Ag/ Fe_3O_4 ODE and
13 Ag/ Fe_3O_4 PHE samples, respectively) matches very well with those obtained from nanoparticle
14 systems with moderate or weak interacting magnetic nanoparticles [46]. The larger values of T_0
15 determined for the sample with flower-like morphology also indicate a strengthening of the
16 dipole-dipole and exchange forces due to intraparticle interactions. Using the Vogel-Fulcher law
17 to analyze τ_0 , E_b and T_0 of those magnetic particles attached to the Ag seeds supports the
18 previous conclusions obtained from DC magnetic studies, where larger energy barrier values
19 were detected for linked magnetic nanoparticles. A summary of the extracted magnetic
20 parameters is presented in Table 2.
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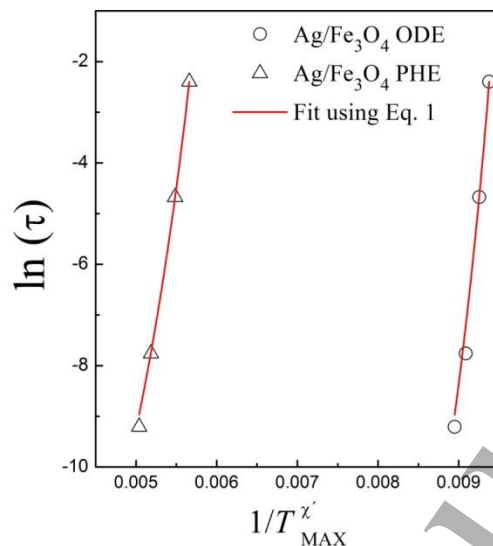


Figure 9. Plots of $\ln(\tau)$ vs. $1/T_{MAX}^{\chi}$ and the respective fit according to Eq. 1 (red solid line). Here, T_{MAX}^{χ} is the broader peak of the real component for every measured frequency.

Table 2. Extracted and calculated parameters from the AC magnetic susceptibility measurements.

Sample	T_{MAX}^{ZFC} (K)	T_{MAX}^{χ} (K)				$T_{MAX}^{\chi''}$ (K)								E_b (erg)	T_0 (K)
		f (Hz)				f (Hz) – peak 1				f (Hz) – peak 2					
		11	207	2347	9967	11	207	2347	9967	11	207	2347	9967		
<i>flower</i> <i>Ag/Fe₃O₄</i> <i>ODE</i>	108	104	105	106	107	35	35	41	42	98	100	102	104	8×10^{-14}	92
<i>dimer</i> <i>Ag/Fe₃O₄</i> <i>PHE</i>	205	180	183	188	190	23	24	31	34	144	145	149	157	2×10^{-13}	133

Conclusions

Controlled synthesis protocols were employed to obtain magnetic-plasmonic nanoheterostructures of Ag/Fe₃O₄ with two distinct morphologies. By controlling the polarity of the reaction medium, the charge compensation can be tuned, leading to particular architecture

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3 formations with Ag and iron oxide nanoparticles, resulting in dimer and flower-like
4 configurations. Despite the fact that the magnetite nanoparticles in the heterocrystals have
5 similar sizes and shapes, the magnetization response seems to reflect different phenomena. The
6 observed dissimilarities on the AC and DC magnetization responses can be related to several
7 factors. In the flower-like sample, the dipolar interactions between the Fe-oxide nanoparticles are
8 stronger than in dimer-like counterpart. Also, in flower-like configuration one could have
9 exchange interactions among Fe-atoms of neighboring particles, owing to the existence of
10 percolated regions. Comparing the magnetic parameters of both Ag/Fe₃O₄ samples with the
11 Fe₃O₄ control sample and performed simulations, we detected important modifications in the
12 magnetic response. The maximum of the zero-field-cooling curve was shifted around 90 K and
13 170 K for dimer and flower-like samples, respectively. Furthermore, the inferred magnetic
14 anisotropies also reveal important variations. While part of these changes can be understood in
15 terms of the strengthening of dipolar interactions or variations in the iron oxide phases, we also
16 addressed the observed magnetic complex behavior from the coupling efficiency between
17 metal/Fe-oxide and its role on the superficial disorder and hence on the magnetization response.
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References

- [1] Melinon P, Begin-Colin S, Duvail J L, Gauffre F, Boime N H, Ledoux G, Plain J, Reiss P, Silly F and Warot-Fonrose B, 2014 *Phys. Rep.* **543** 163-197.
- [2] Xia Y, Gilroy K D, Peng H C and Xia X 2017 *Angew. Chem.* **56** 60-95.
- [3]. Veverka M, Zaveta K, Kaman O, Veverka P, Knizek K, Pollert E, Burian M and Kaspá P 2014 *J. Phys. D: Appl. Phys.* **47** 065503
- [5] Costa L S and Zanchet D 2017 *Catal. Today.* **282** 151-158.
- [5] Wang C, Yin H, Dai S and Sun S 2010 *Chem. Mater.* **22** 3277-3282.
- [6] Amarjargal A, Tijing L D, Im I T and Kim C S 2013 *Chem. Eng. J.* **226** 243-254.
- [7] Gong P, Li H, He X, Wang K, Hu J, Tan W, Zhang S and Yang X 2007 *Nanotechnology.* **18** 285604.
- [8] Brollo M E F, Orozco-Henao J M, López-Ruiz R, Muraca D, Dias C S B, Pirota K R and Knobel M 2016 *J. Magn. Magn. Mater.* **397** 20-27.
- [9] Das R, Rinaldi-Montes N, Alonso J, Amghouz Z, Garaio E, García J A, Gorria P, Blanco J A, Phan M H and Srikanth H 2016 *ACS Appl. Mater. Interfaces.* **8** 25162–25169.
- [10] Sun L, He J, An S, Zhang J and Ren D 2013 *J. Mol. Struct.* **1046** 74-81.
- [11] Yu H, Chen M, Rice P M, Wang S X, White R L, Sun S 2005 *Nano. Lett.* **2** 379-382.
- [12] Moscoso-Londoño O, Muraca D, Tancredi P, Cosio-Castañeda C, Pirota K R, Socolovsky L M, 2014 *J. Phys. Chem. C.* **118** 13168-13176.
- [13] Y. Sun 2015 *Natl. Sci. Rev.* **2** 329–348.
- [14] Brollo M E, Lopez-Ruiz R, Muraca D, Figueroa S J, Pirota K R and Knobel M 2014 *Sci. Rep.* **4** 6839.
- [15] Chandra S, Huls N A, Phan M H, Srinath S, Garcia M A, Lee Y, Wang C, Sun S, Iglesias O and Srikanth H 2014 *Nanotechnology.* **25** 055702.
- [16] Choi S H, Na H B, Park Y I, An K, Kwon S G, Jang Y, Park M H, Moon J, Son J S, Song I C, Moon W K and Hyeon T, S 2008 *J. Am. Chem. Soc.* **130** 15573–15580.
- [17] Hyeon T, Lee S S, Park J, Chung Y and Na H B 2001 *J. Am. Chem. Soc.* **123** 12798-12801.
- [18] Vreeland E C, Watt J, Schober G B, Hance B G, Austin M J, Price A D, Fellows D B, Monson T C, Hudak N S, Maldonado-Camargo L, Bohorquez A C, Rinaldi C and Huber D L 2015 *Chem. Mater.* **27** 6059-6066.

- 1
2
3 [19] Wang C, Xu C, Zeng H and Sun S 2009 *Adv. Mater.* **21** 3045-3052.
4
5 [20] Huang J, Sun Y, Huang S, Yu K, Zhao Q, Peng F, Yu H, Wang H and Yang J, 2011 *J.*
6 *Mater. Chem.* **21** 17930.
7
8 [21] Wei Y, Klajn R, Pinchuk A O, Grzybowski B A 2008 *Small.* **4** 1635-1639.
9
10 [22] Gallo J, García I, Padro D, Arnáiza B and Penadés S 2010 *J. Mater. Chem.* **20** 10010-
11 10020.
12
13 [23] Hao D, Cheng-Min S, Chao H, Zhi-Chuan X, Chen L, Yuan T, Xue-Zhao S and Hong-Jun
14 G 2010 *Chin. Phys. B.* **19** 066102.
15
16 [24] Zhang L, Dou Y H and Gu H C 2006 *J. Colloid. Interface. Sci.* **297** 660-664.
17
18 [25] Fantechi E, Roca A G, Sepúlveda B, Torruella P, Estradé S, Peiró F, Coy E, Jurga S, Bastús
19 N G, Nogués J and Puntès V 2017 *Chem. Mater.* **29** 4022-4035.
20
21 [26] Schladt T D, Graf T, Köhler O, Bauer H, Dietzsch M, Mertins J, Branscheid R, Kolb U and
22 Tremel W 2012 *Chem. Mater.* **24** 525-535.
23
24 [27] Muraca D, Sharma S K, Socolovsky L M, de Siervo A, Lopes G and Pirota K R 2012 *J.*
25 *Nanosci. Nanotechnol.* **12** 6961-6967.
26
27 [28] Knobel M, Nunes W C, Socolovsky L M, Biasi E D, Vargas J M and Denardin J C 2008, *J.*
28 *Nanosci. Nanotechnol.* **8** 2836-2857.
29
30 [29] Socolovsky L M and Moscoso-Londoño O. 2017 Consequences of Magnetic Interactions
31 Phenomena in Granular Systems, in: S.K. Sharma (Ed.) *Complex Magnetic Nanostructures -*
32 *Synthesis, Assembly and Applications*, Springer.
33
34 [30] Orozco-Henao J M, Coral D F, Muraca D, Moscoso-Londoño O, Mendoza Zélis P,
35 Fernandez van Raap M B, Sharma S K, Pirota K R and Knobel M 2016 *J. Phys. Chem. C.* **120**
36 12796-12809.
37
38 [31] Moscoso-Londoño O, Tancredi P, Muraca D, Mendoza Zélis P, Coral D, Fernández van
39 Raap M B, Wolff U, Neu V, Damm C, de Oliveira C L P, Pirota K R, Knobel K R and
40 Socolovsky L M 2017 *J. Magn. Magn. Mater.* **428** 105-118.
41
42 [32] Tancredi P, Rivas Rojas P C, Moscoso-Londoño O, Wolff U, Neu V, Damm C, Rellinghaus
43 B, Knobel M and Socolovsky L M 2017 *New J. Chem.* **41** 15033-15041.
44
45 [33] Jiang J, Gu H, Shao H, Devlin E, Papaefthymiou G C and Ying J Y 2008 *Adv. Mater.* **20**
46 4403-4407.
47
48
49
50
51
52
53
54
55
56
57
58
59
60

- 1
2
3 [34] Umut E, Pineider F, Arosio P, Sangregorio C, Corti M, Tabak F, Lascialfari A and Ghigna P
4 2012 *J. Magn. Magn. Mater.* **324** 2373-2379.
5
6 [35] Muraca D, de Siervo A and Pirola K R 2012 *J. Nanopart. Res.* **15** 1375
7
8 [36] Hansen M F and Mørup S 1999 *J. Magn. Magn. Mater.* **203** 214-216.
9
10 [37] Rondinone A J, Liu C and Zhang Z J 2001 *J. Phys. Chem. B.* **105** 7967-7971.
11
12 [38] Pineider F, Fernández C D J, Videtta V, Carlino E, Hourani A A, Wilhelm F, Rogalev A,
13 Cozzoli P D, Ghigna P and Sangregorio C 2013 *ACS Nano.* **7** 857-866.
14
15 [39] Schick I, Gehrig D, Montigny M, Balke B, Panthöfer M, Henkel A, Laquai F and Tremel W
16 2015 *Chem. Mater.* **27** 4877-4884.
17
18 [40] Sytnyk M, Kirchschrager R, Bodnarchuk M I, Primetzhofer D, Kriegner D, Enser H, Stangl
19 J, Bauer P, Voith M, Hassel A W, Krumeich F, Ludwig F, Meingast A, Kothleitner G,
20 Kovalenko M V and Heiss W 2013 *Nano. Lett.* **13** 586-593.
21
22 [41] Laha S S, Regmi R and Lawes G 2013 *J. Phys. D: Appl. Phys.* **46** 325004.
23
24 [42] León Félix L, Coaquira J A H, Martínez M A R, Goya G F, Mantilla J, Sousa M H,
25 Valladares L D L S, Barnes C H W, Morais P C 2017 *Sci. Rep.* **7** 41732.
26
27 [43] Morales M B, Phana M H, Pal S, Frey N A and Srikantha H 2009 *J. Appl. Phys.*, **105**
28 07B511.
29
30 [44] Masunaga S H, Jardim R F, Fichtner P F P and Rivas J 2009 *Phys. Rev. B.* **80** 184428.
31
32 [45] Landi G T, Arantes F R, Cornejo D R, Bakuzis A F, Andreu I, Natividad E, *J. Magn. Magn.*
33 *Mater.* **421** 138-151
34
35 [46] Lee J S, Tan R P, Wu J H and Kim Y K 2011 *Appl. Phys. Lett.* **99** 062506.
36
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